Geometry of phase space transport in a variety of mechanical systems

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University of Paderborn, May 2011

Chaotic phase space transport via lobe dynamics

- As our dynamical system, we consider a discrete map[1]

$$f : \mathcal{M} \longrightarrow \mathcal{M},$$

  e.g., $f = \phi_t^{t+T}$, where $\mathcal{M}$ is a differentiable, orientable, two-dimensional manifold e.g., $\mathbb{R}^2$, S^2

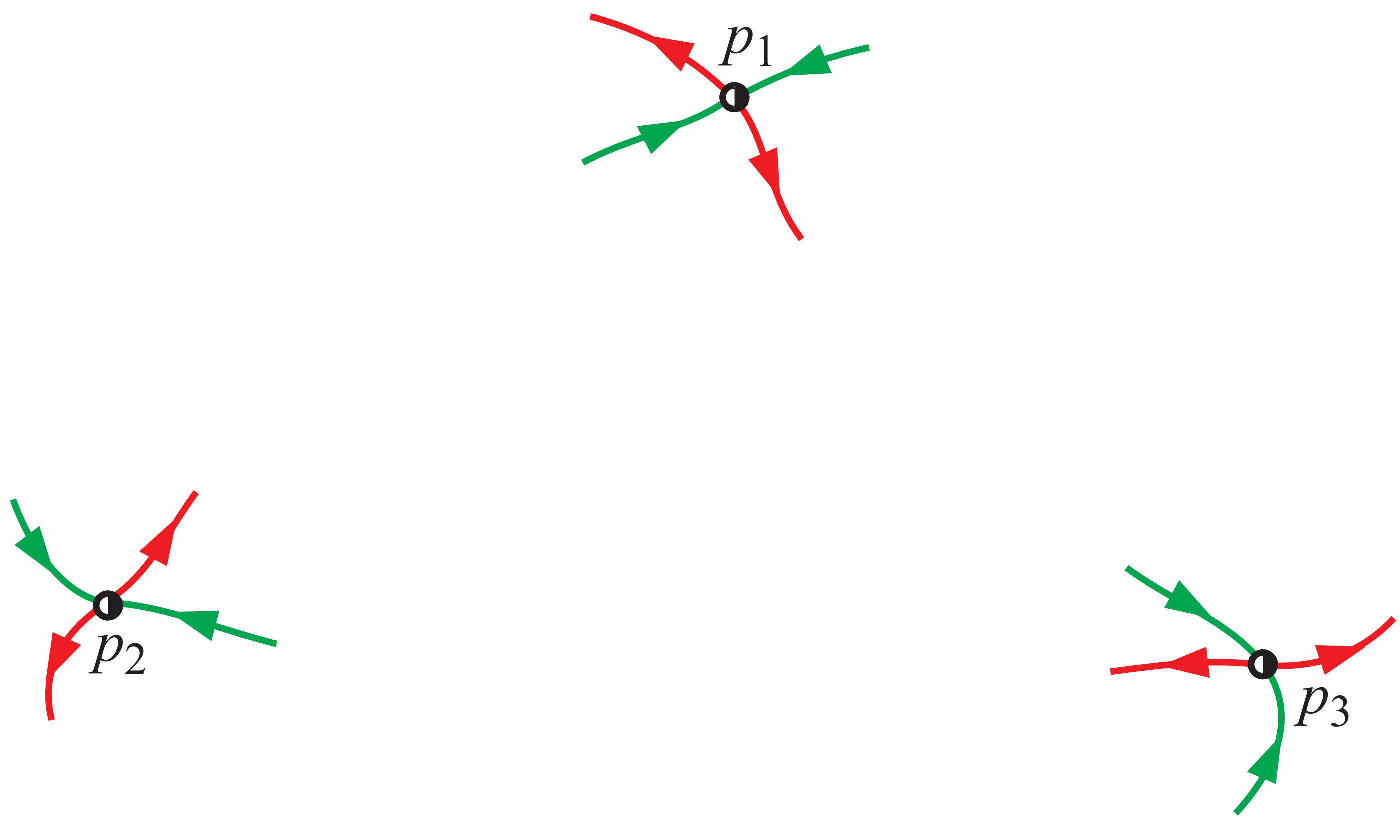
- To understand the transport of points under the map f, we consider the **invariant manifolds of unstable fixed points**

- Let $p_i, i = 1, \ldots, N_p$, denote a collection of saddle-type hyperbolic fixed points for f.

[1] Following Rom-Kedar and Wiggins [1990]

Partition phase space into regions

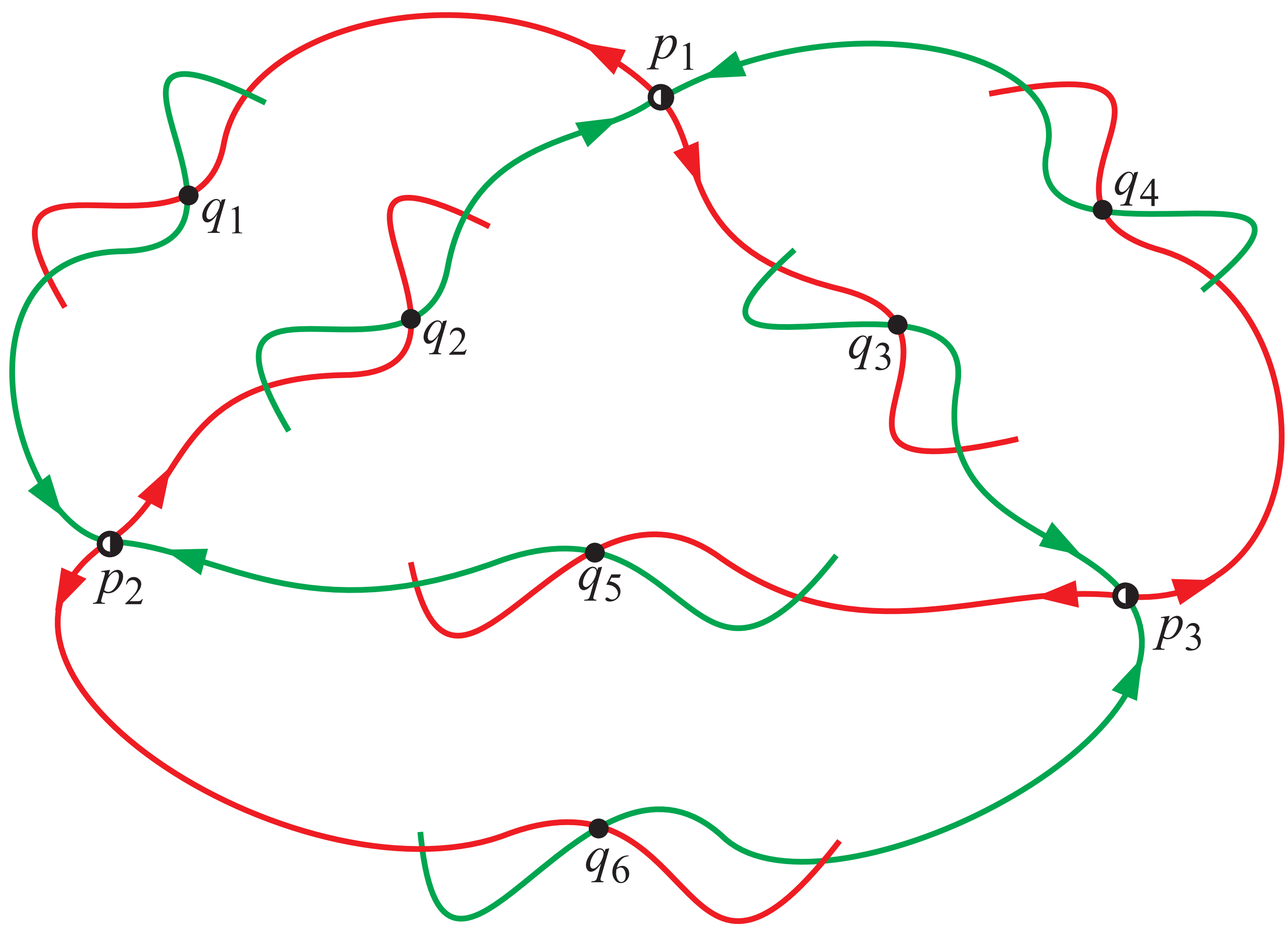
- Natural way to partition phase space
  - Pieces of $W^u(p_i)$ and $W^s(p_i)$ partition $\mathcal{M}$.



Unstable and stable manifolds in **red** and **green**, resp.

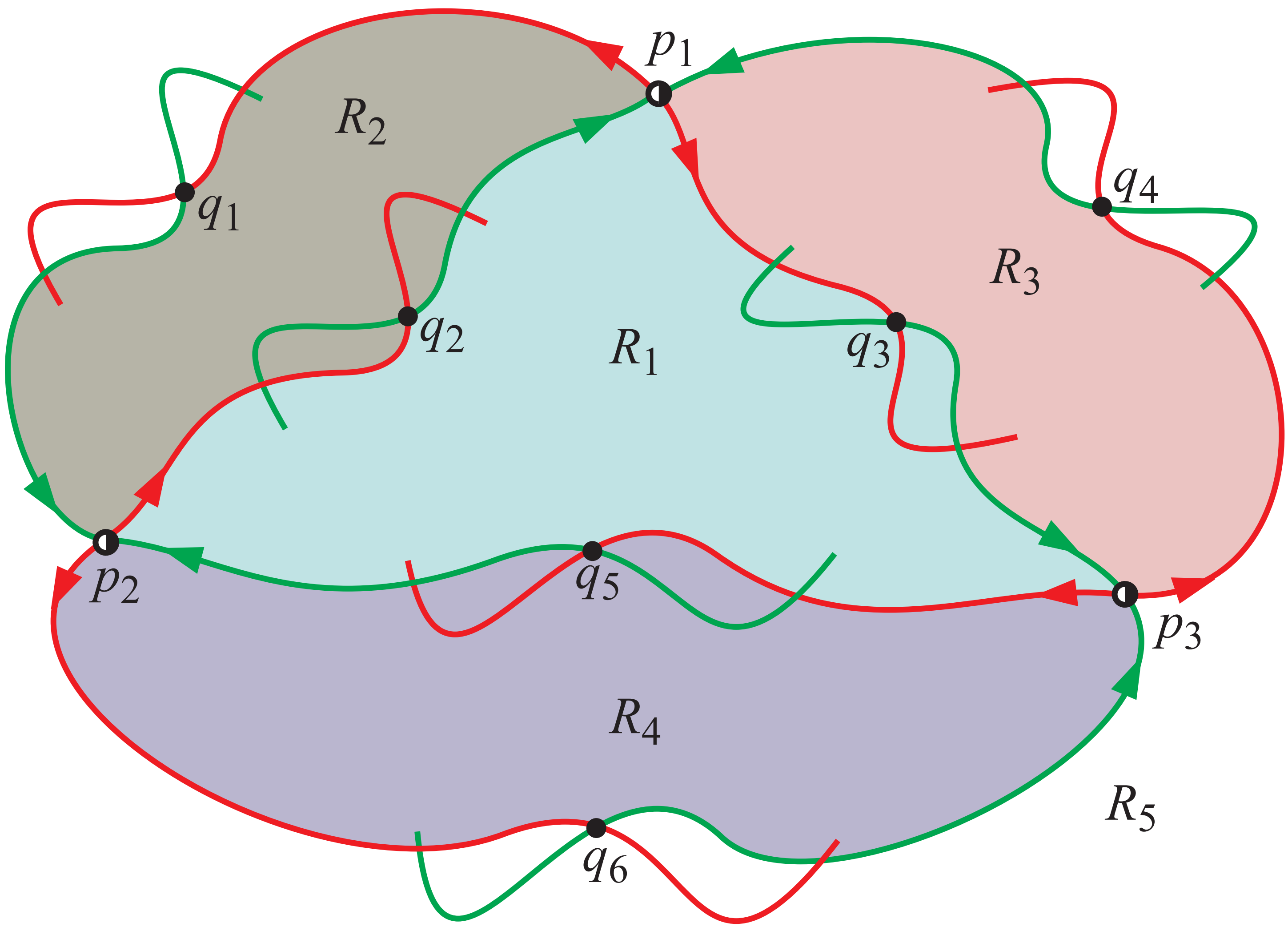
Partition phase space into regions

- Intersection of unstable and stable manifolds define **boundaries**.



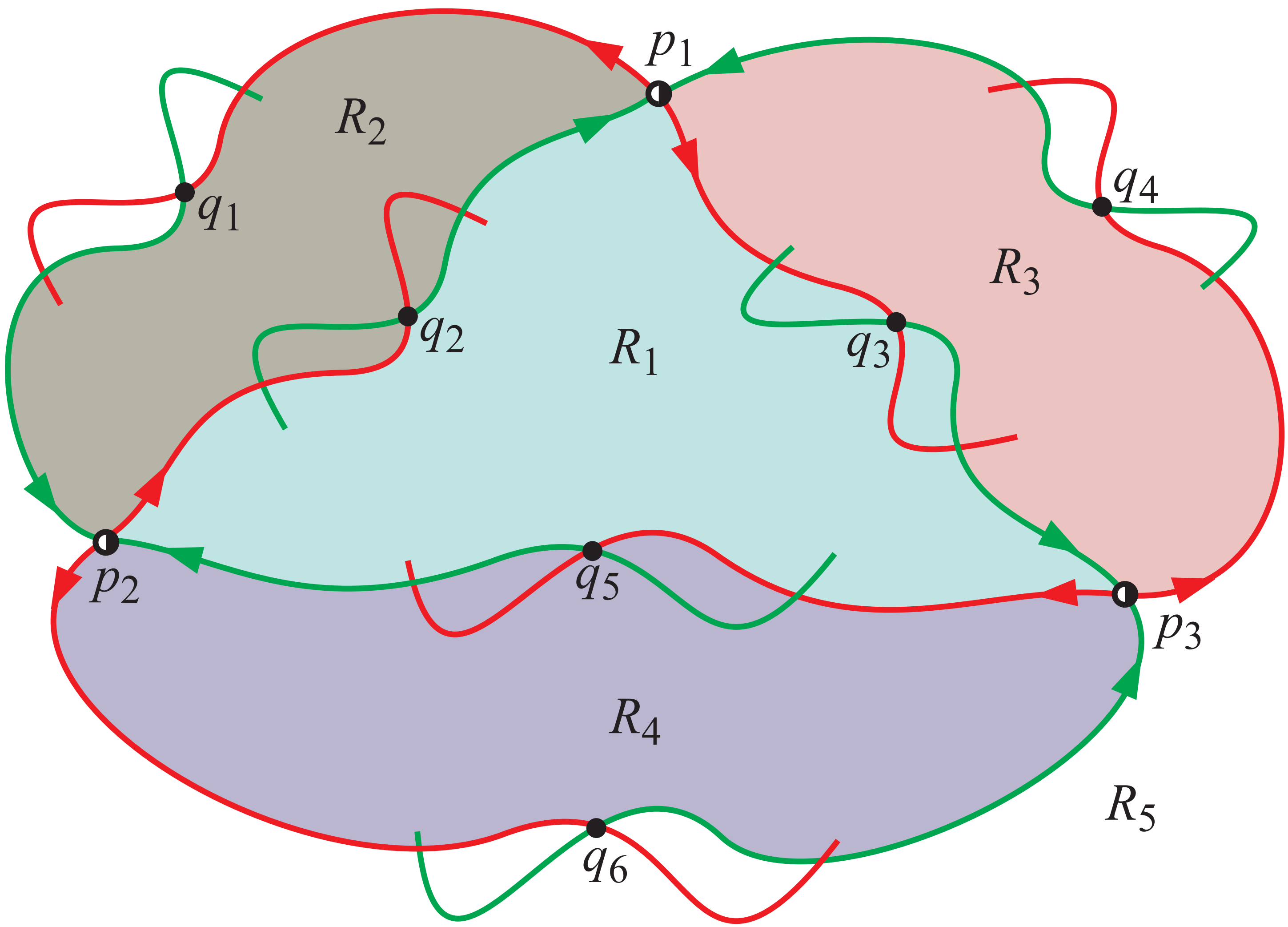
Partition phase space into regions

- These boundaries divide the phase space into **regions**.



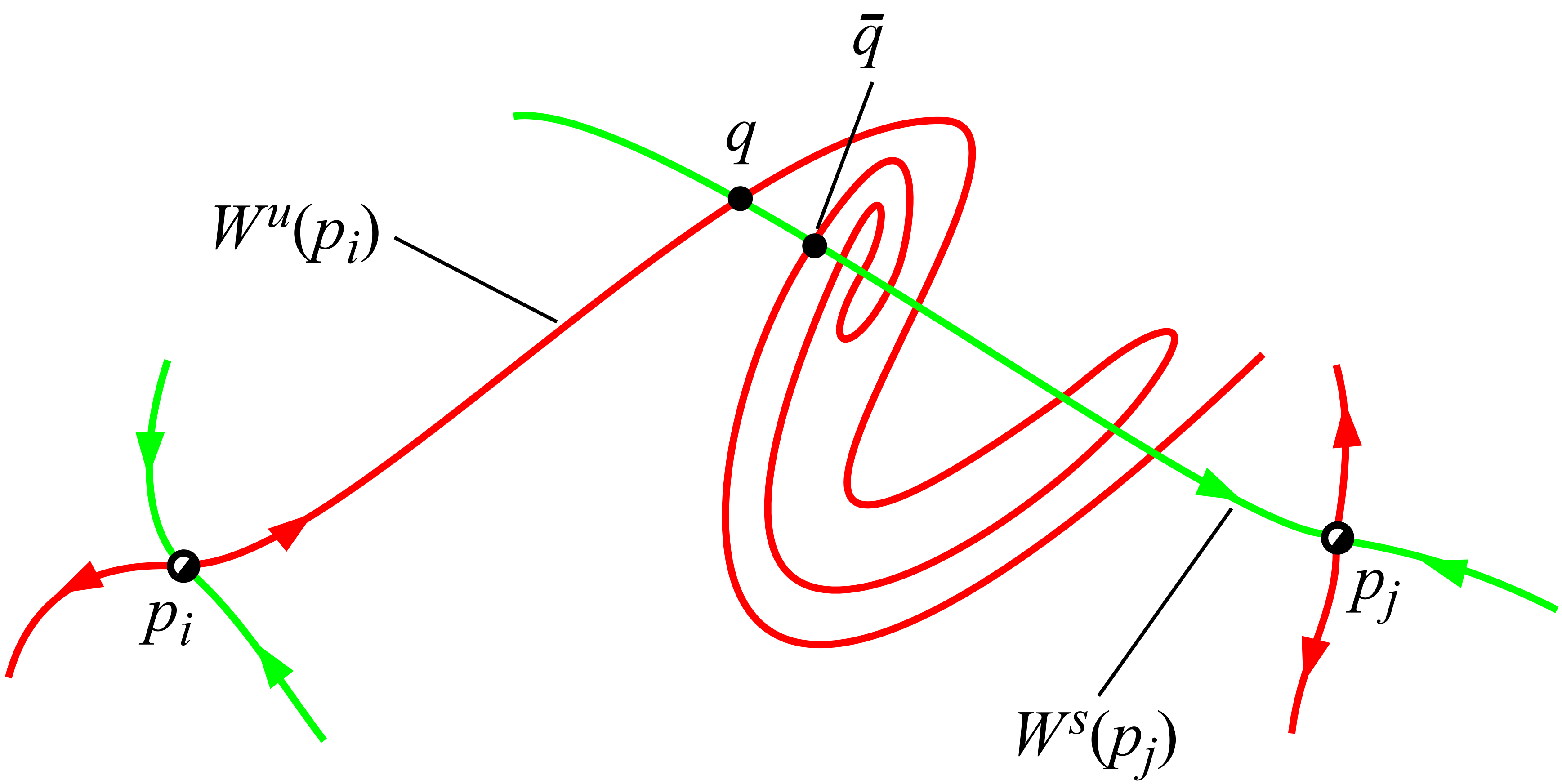
Label mobile subregions: 'atoms' of transport

- Can label mobile subregions based on their past and future whereabouts under one iterate of the map, e.g., $(\ldots, R_3, R_3, [R_1], R_1, R_2, \ldots)$



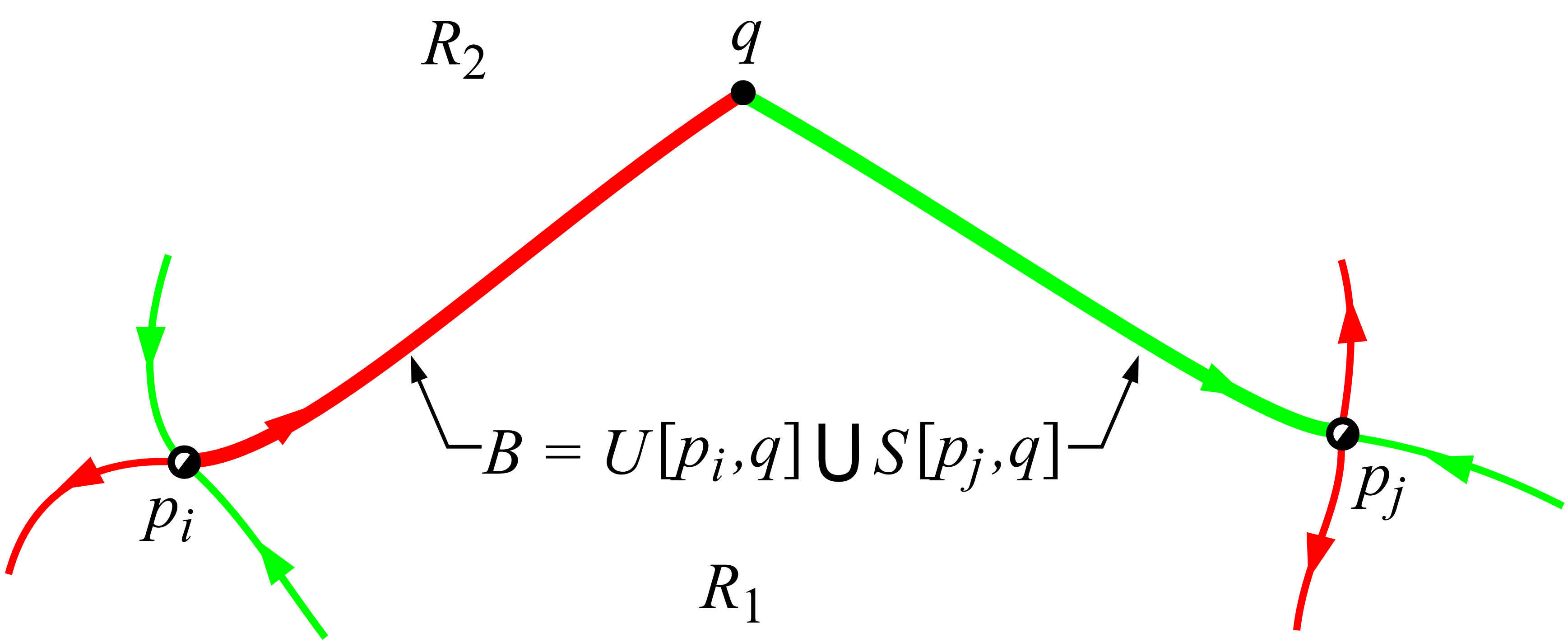
Primary intersection points (pips) and boundaries

- q is a primary intersection point (pip), $\bar{q}$ is not a pip.



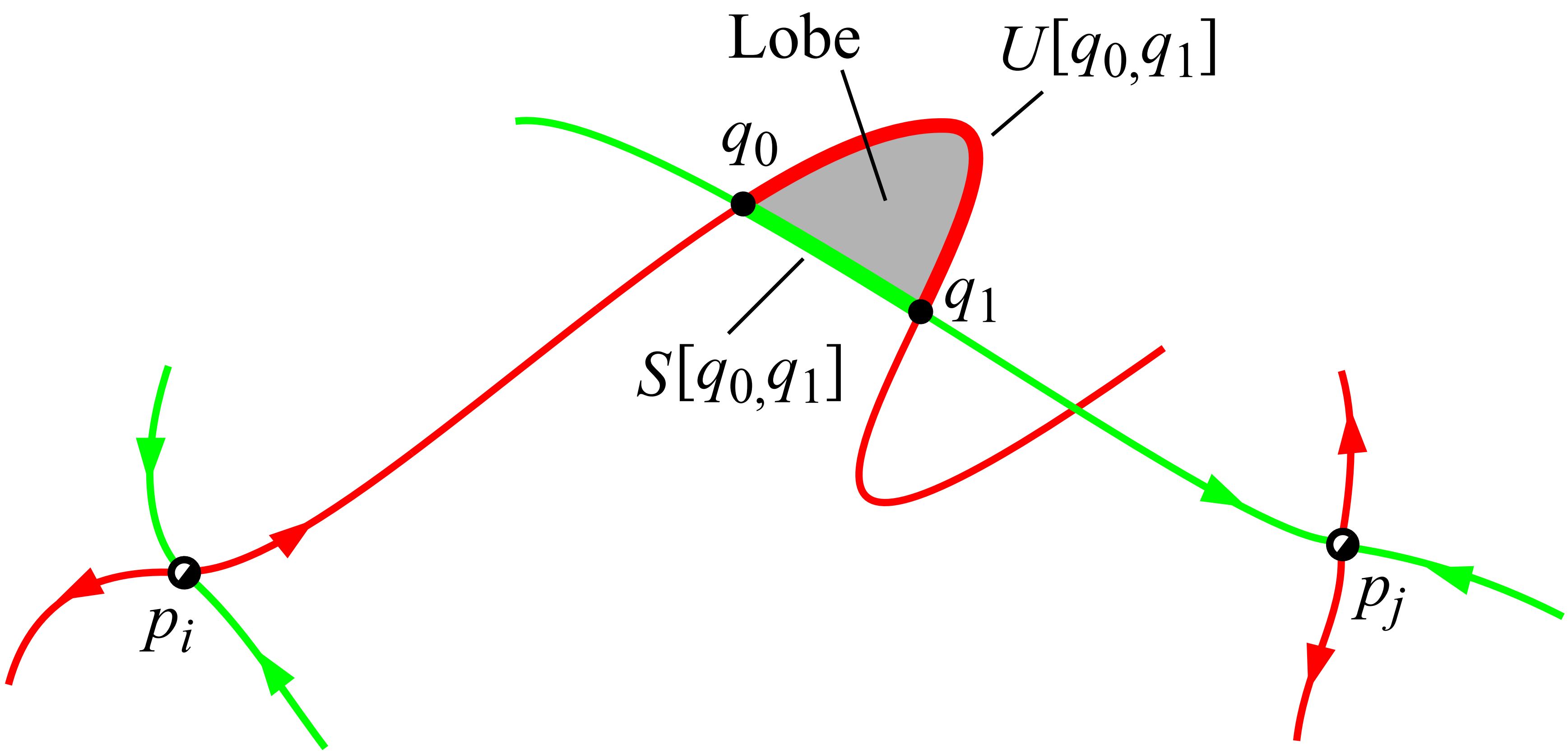
Primary intersection points (pips) and boundaries

- Suppose $W^u(p_i)$ and $W^s(p_j)$ intersect in the pip q. Define $B \equiv U[p_i, q] \bigcup S[p_j, q]$ as a **boundary** between "two sides," R_1 and R_2.



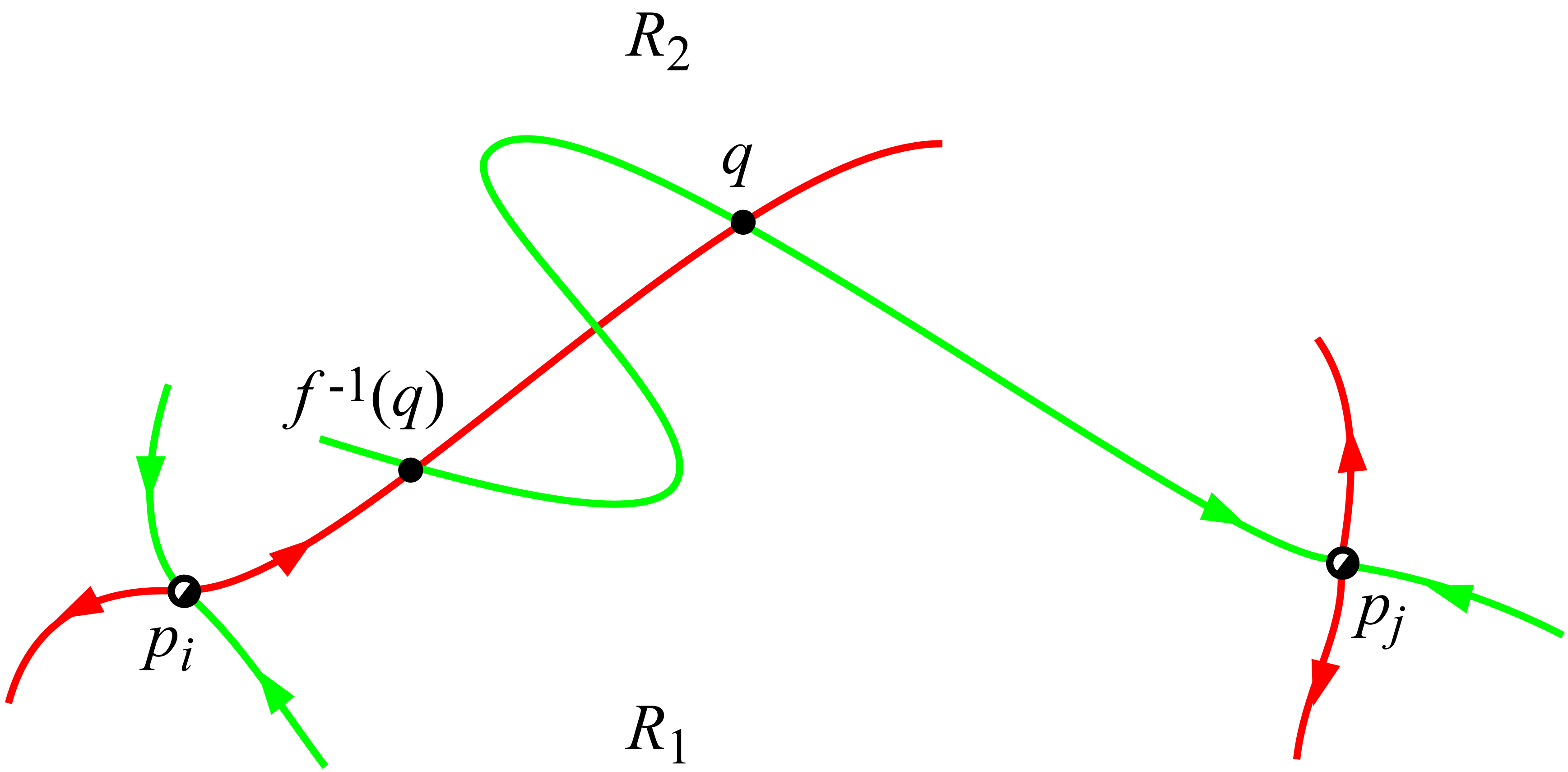
Lobes: the mobile subregions

- Let $q_0, q_1 \in W^u(p_i) \bigcap W^s(p_j)$ be two adjacent pips, i.e., there are no other pips on $U[q_0, q_1]$ and $S[q_0, q_1]$. The region interior to $U[q_0, q_1] \bigcup S[q_0, q_1]$ is a **lobe**.



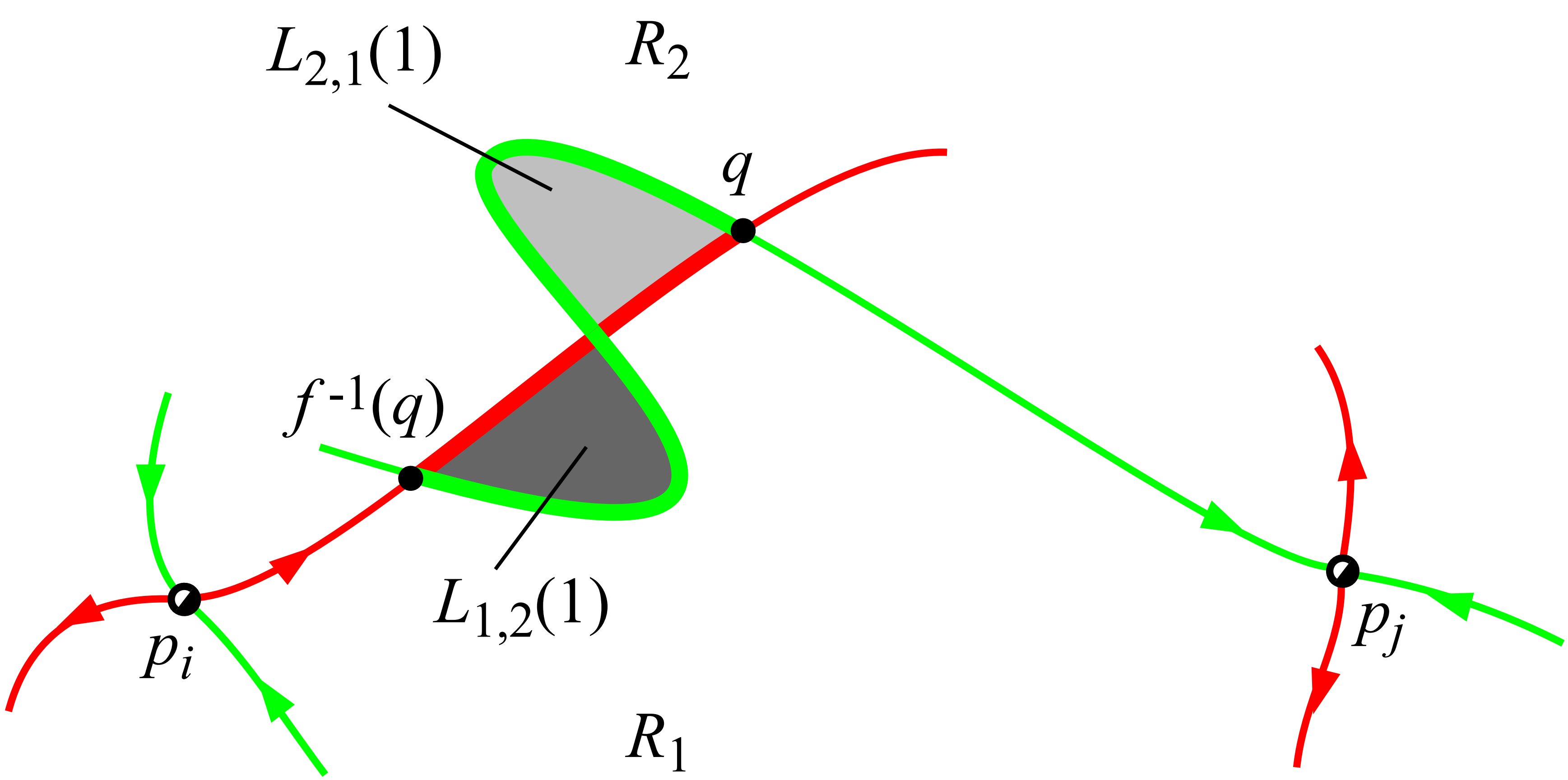
Lobe dynamics: transport across a boundary B

- $f^{-1}(q)$ is a pip. f is orientation-preserving $\Rightarrow$ there's *at least one* pip on $U[f^{-1}(q), q]$ where the $W^u(p_i), W^s(p_j)$ intersection is topologically transverse.



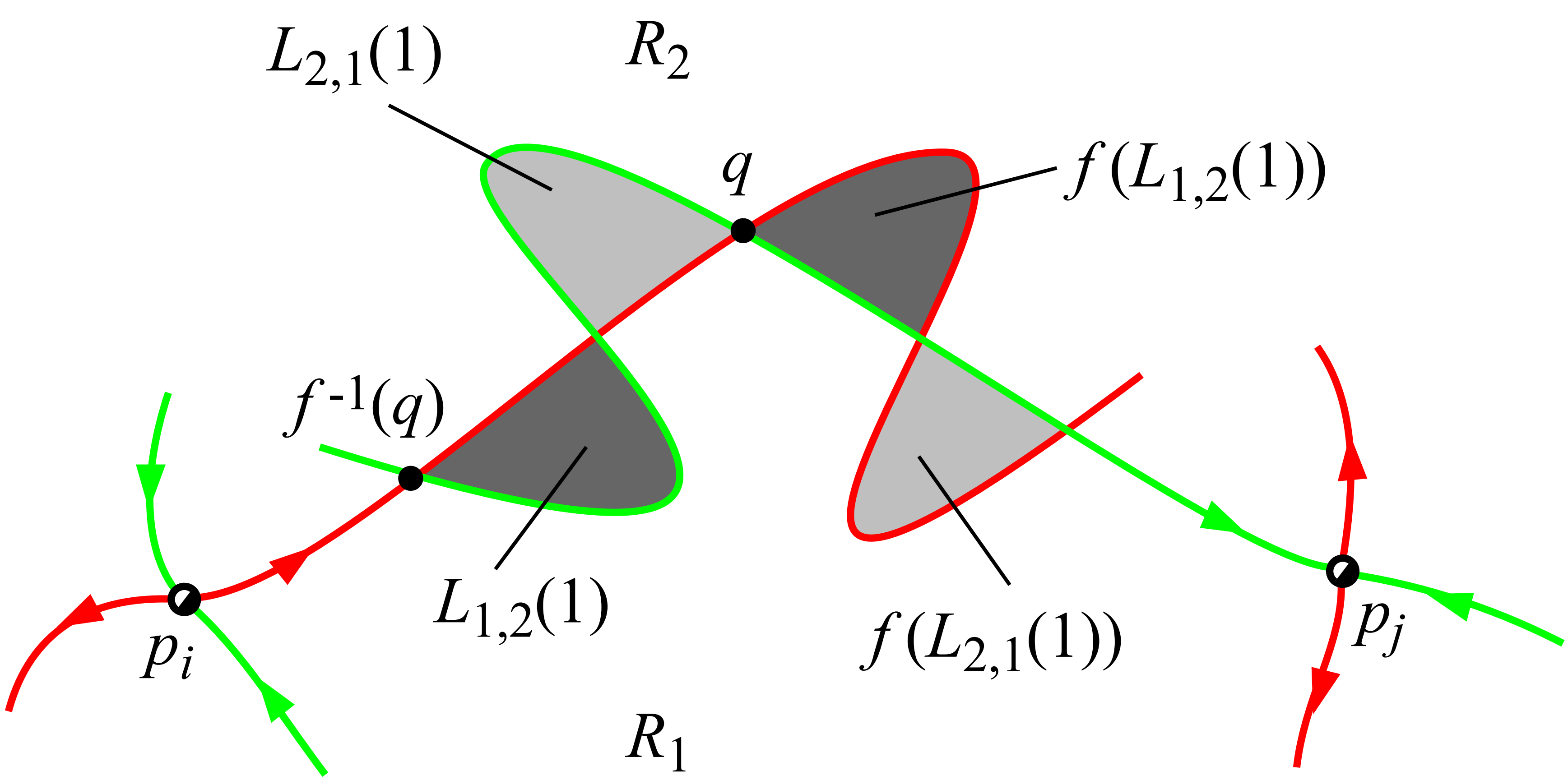
Lobe dynamics: transport across a boundary B

- $U[f^{-1}(q), q] \bigcup S[f^{-1}(q), q]$ forms boundary of two lobes; one in R_1, labeled $L_{1,2}(1)$, or equivalently $([R_1], R_2)$, where $f(([R_1], R_2)) = (R_1, [R_2])$, etc. for $L_{2,1}(1)$



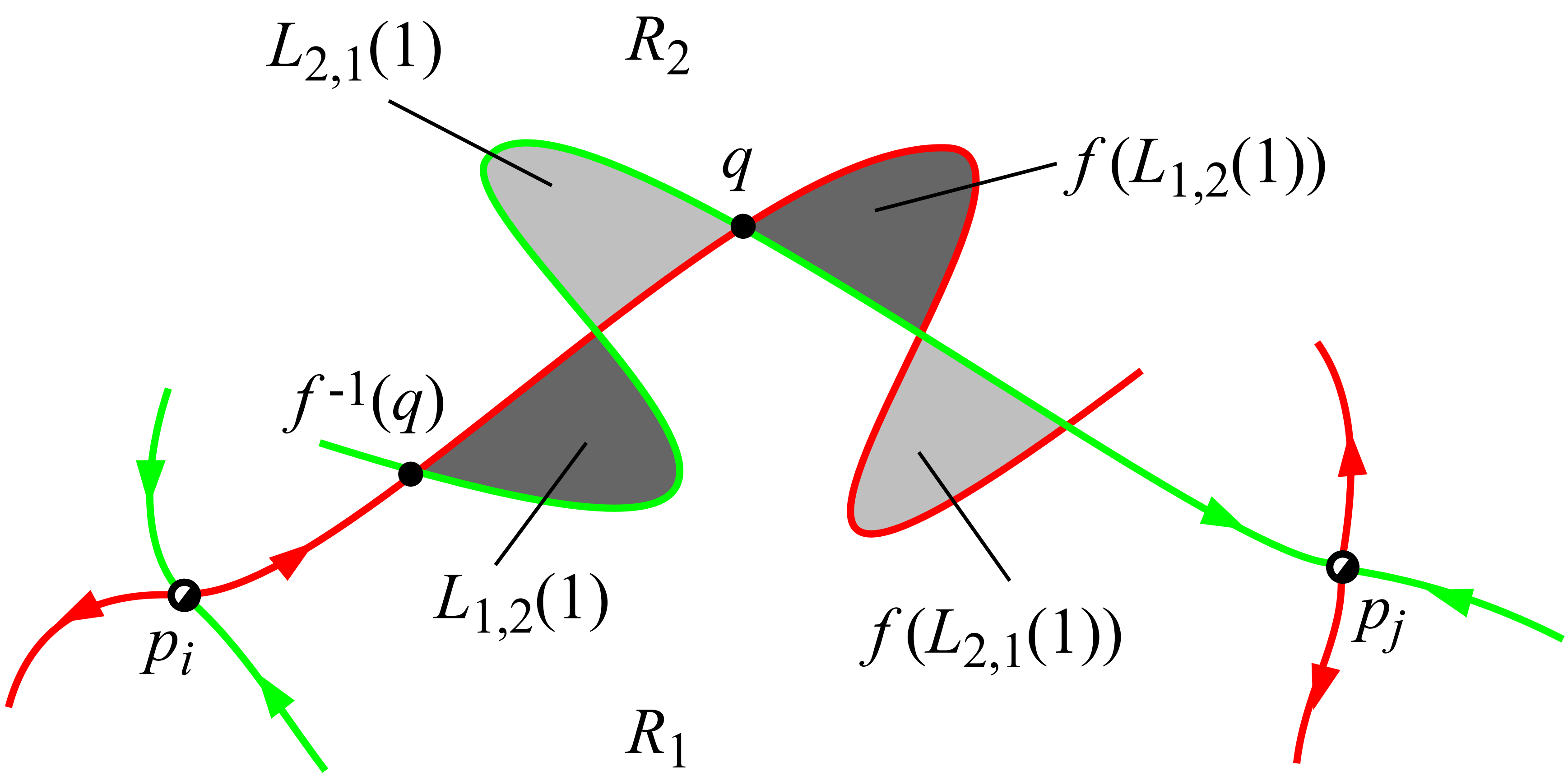
Lobe dynamics: transport across a boundary B

- Under one iteration of f, *only points in* $L_{1,2}(1)$ can move from R_1 into R_2 by crossing B, etc.
- The two lobes $L_{1,2}(1)$ and $L_{2,1}(1)$ are called a **turnstile**.



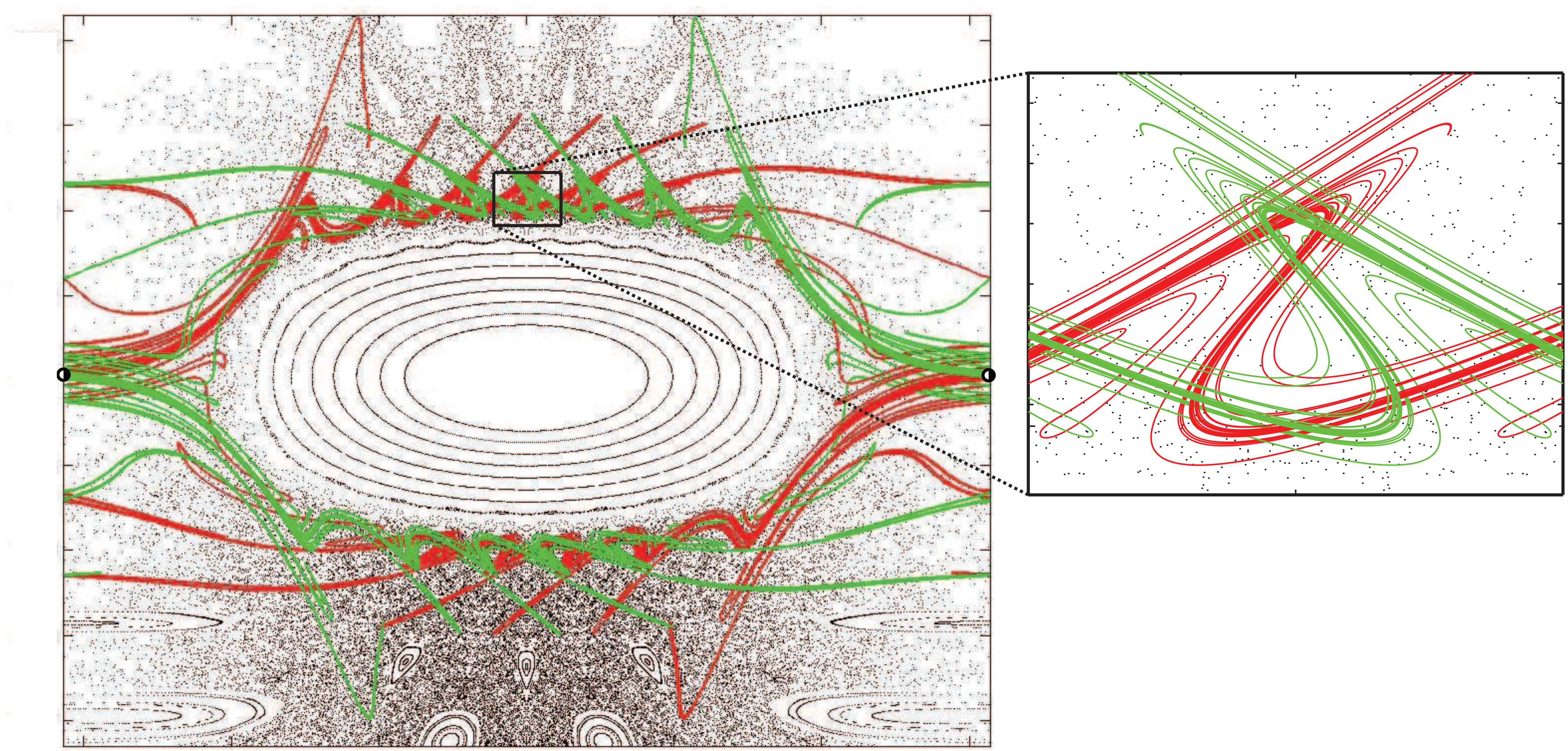
Lobe dynamics: transport across a boundary B

- Essence of lobe dynamics: **the dynamics associated with crossing B is reduced to the dynamics of the turnstile lobes associated with B.**



Identifying atoms of transport by itinerary

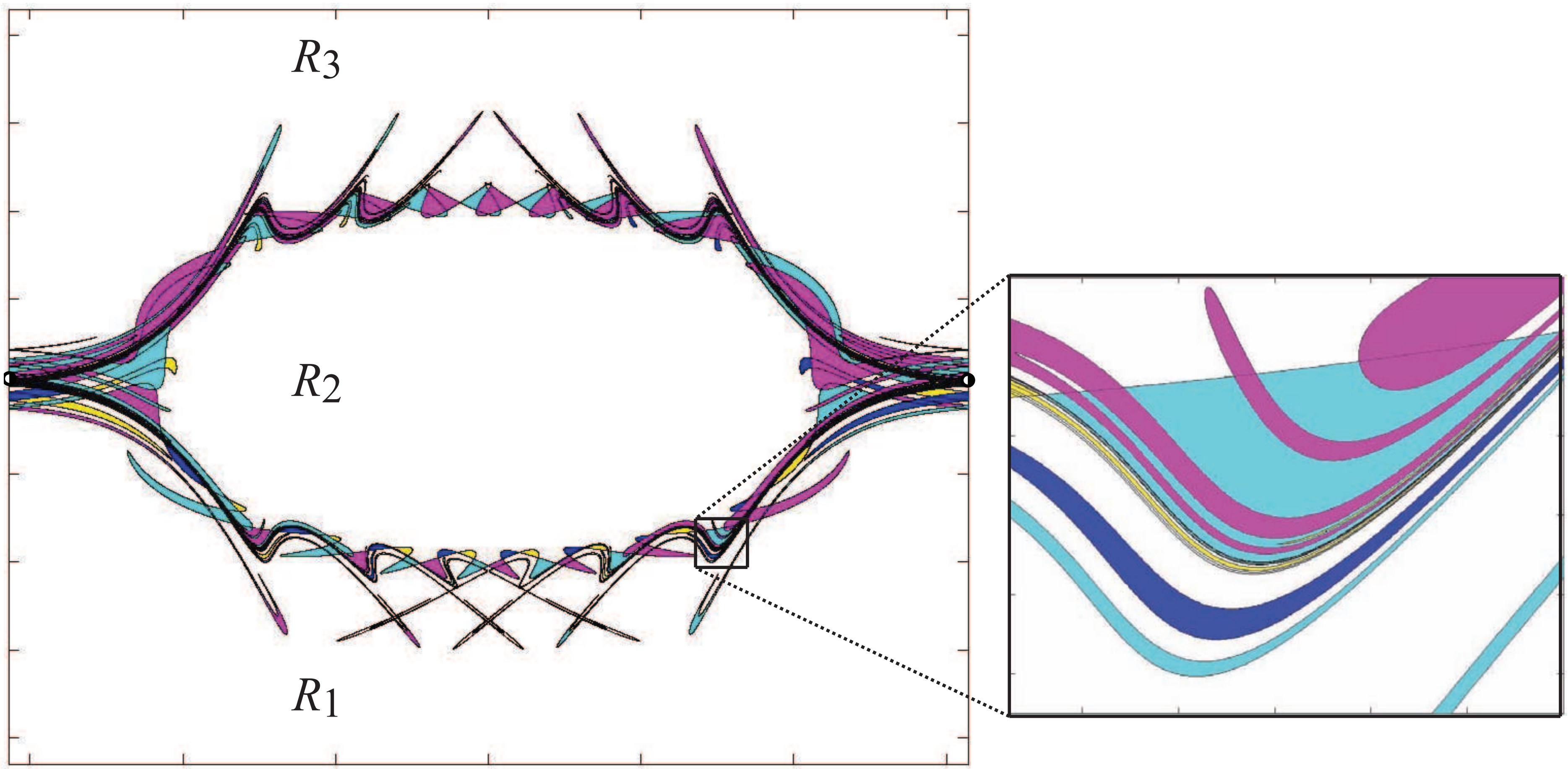
- In a complicated flow, can still identify manifolds ...



Unstable and stable manifolds in **red** and **green**, resp.

Identifying atoms of transport by itinerary

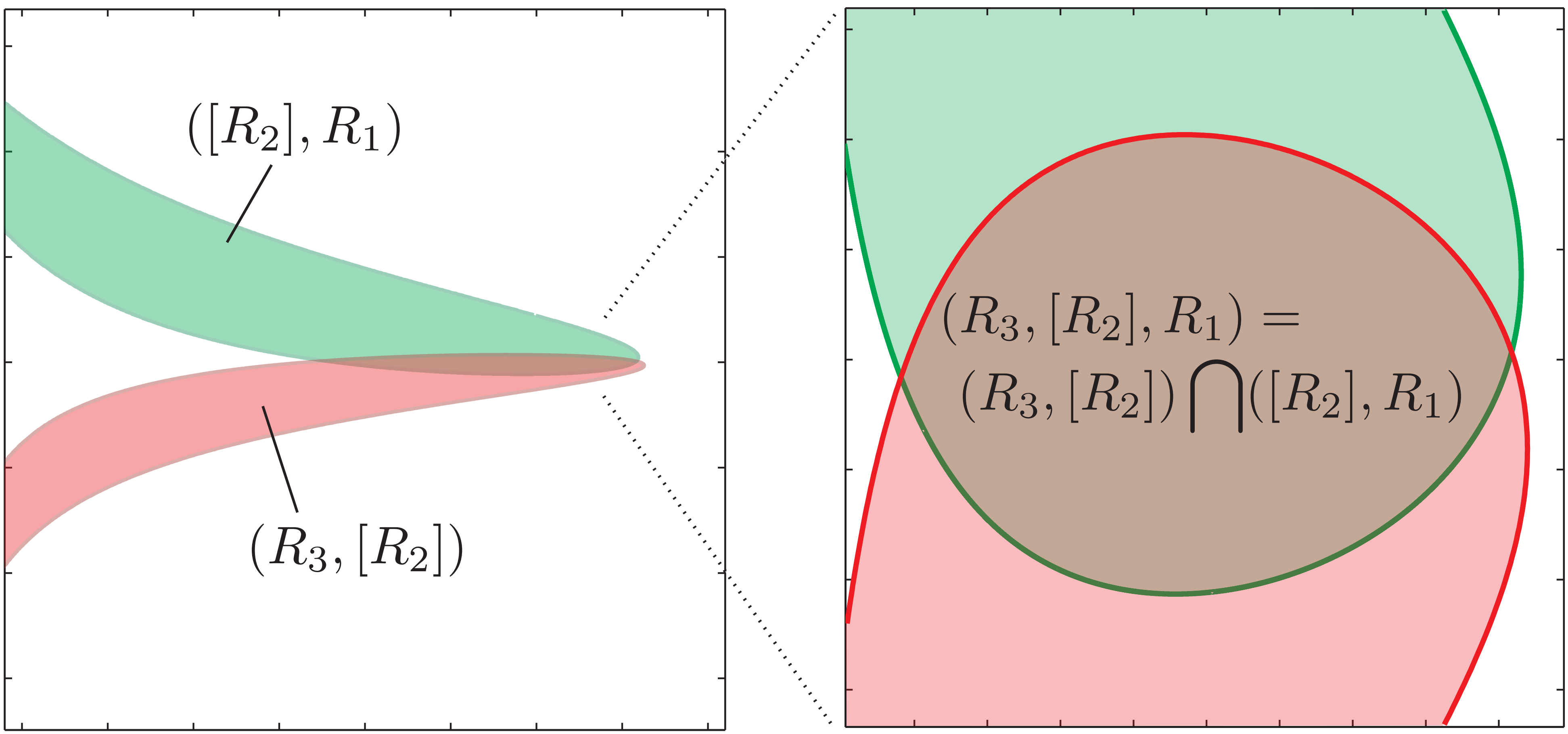
- … and lobes



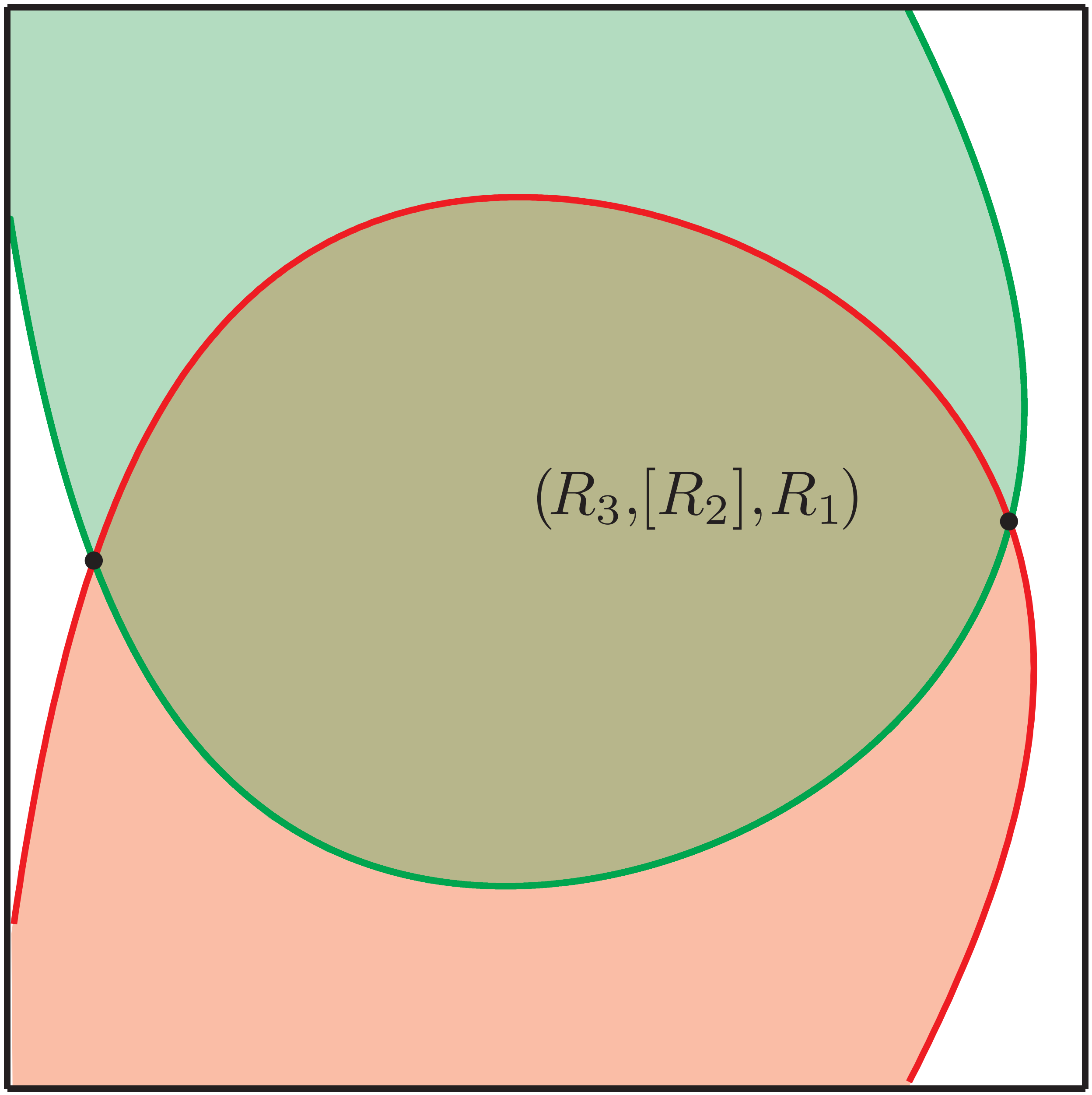
Significant amount of fine, filamentary structure.

Identifying atoms of transport by itinerary

- e.g., with three regions $\{R_1, R_2, R_3\}$, label lobe intersections accordingly.
  - Denote the intersection $(R_3, [R_2]) \bigcap ([R_2], R_1)$ by $(R_3, [R_2], R_1)$

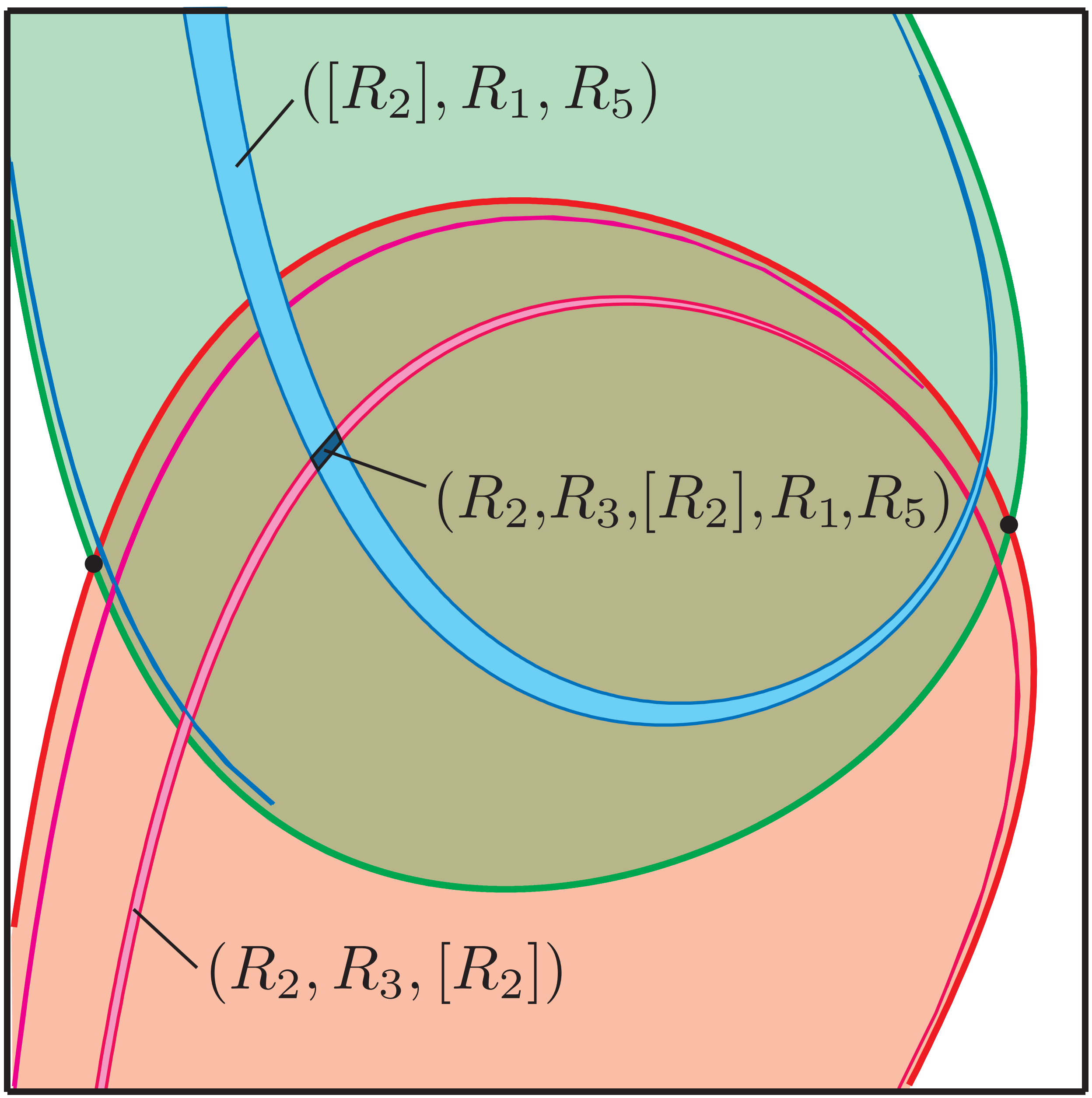


Identifying atoms of transport by itinerary



Longer itineraries...

Identifying atoms of transport by itinerary



... correspond to smaller pieces of phase space; horseshoe dynamics, etc

Aperiodic, finite-time setting

- Many systems defined from data or large-scale simulations

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  — no fixed points or periodic orbits (or their manifolds)

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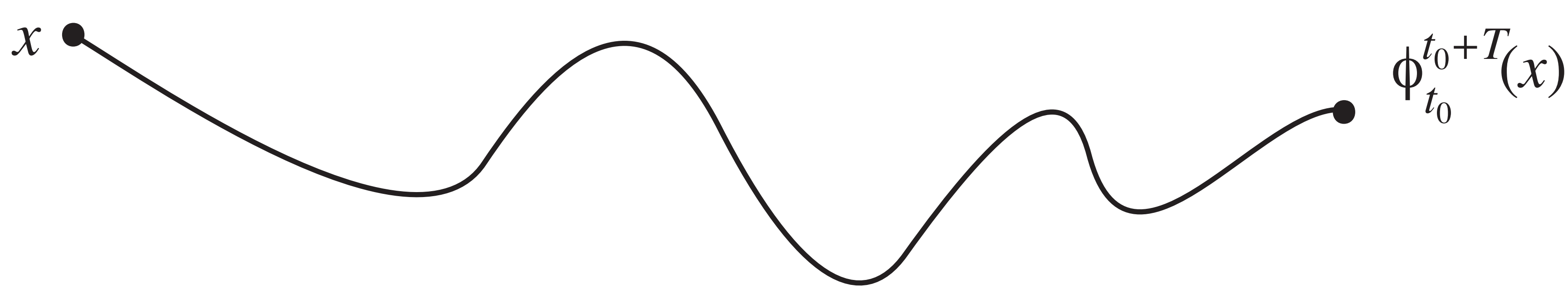
Aperiodic, finite-time setting

- Many systems defined from data or large-scale simulations
- e.g., atmospheric winds are a time-chaotic flow field
  — no fixed points or periodic orbits (or their manifolds)
- How do we get at transport?
- Recall the flow

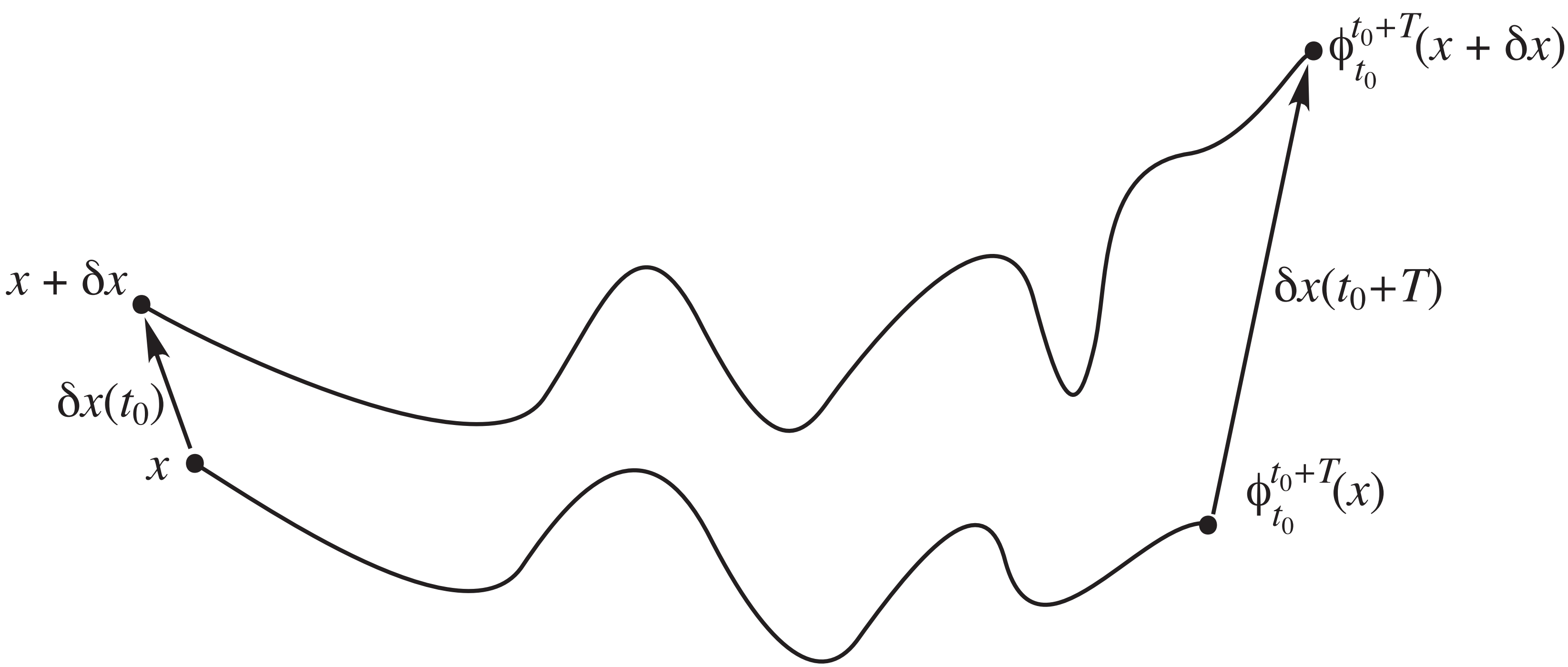
$$x \mapsto \phi_{t_0}^{t_0+T}(x)$$



Aperiodic, finite-time setting

- Small initial perturbations $\delta x(t_0)$ grow like

$$\begin{aligned}\delta x(t_0+T) &= \phi_{t_0}^{t_0+T}(x+\delta x(t_0)) - \phi_{t_0}^{t_0+T}(x) \\ &= \frac{d\phi_{t_0}^{t_0+T}(x)}{dx}\delta x(t_0) + O(||\delta x(t_0)||^2)\end{aligned}$$



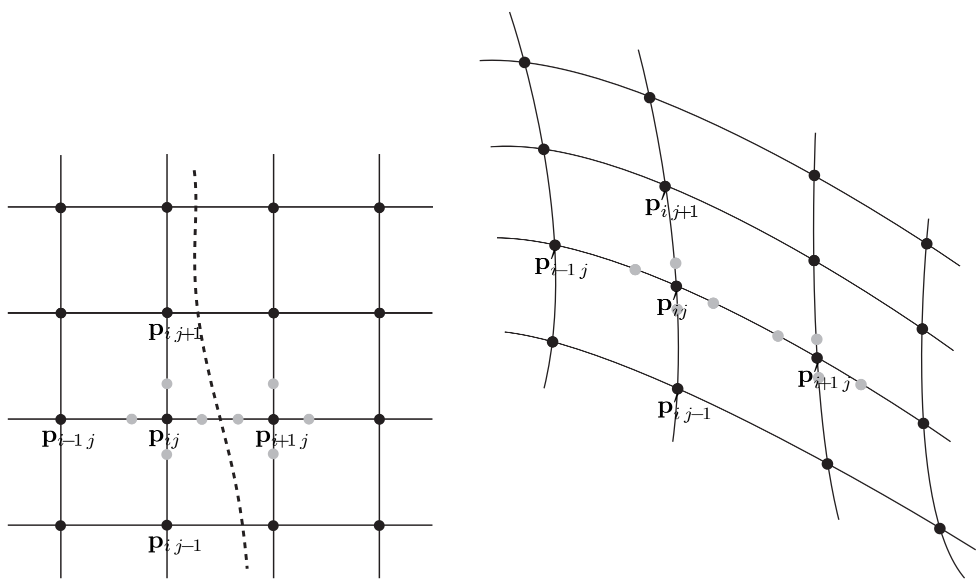
Invariant manifold analogs: FTLE-LCS approach

- The finite-time Lyapunov exponent (FTLE),

$$\sigma_t^T(x) = \frac{1}{|T|} \log \left\| \frac{d\phi_t^{t+T}(x)}{dx} \right\|$$

  measures the maximum stretching rate over the interval T of trajectories starting near the point x at time t

- Ridges of σ_t^T are candidate hyperbolic codim-1 surfaces; finite-time analogs of stable/unstable manifolds; Lagrangian coherent structures[2]



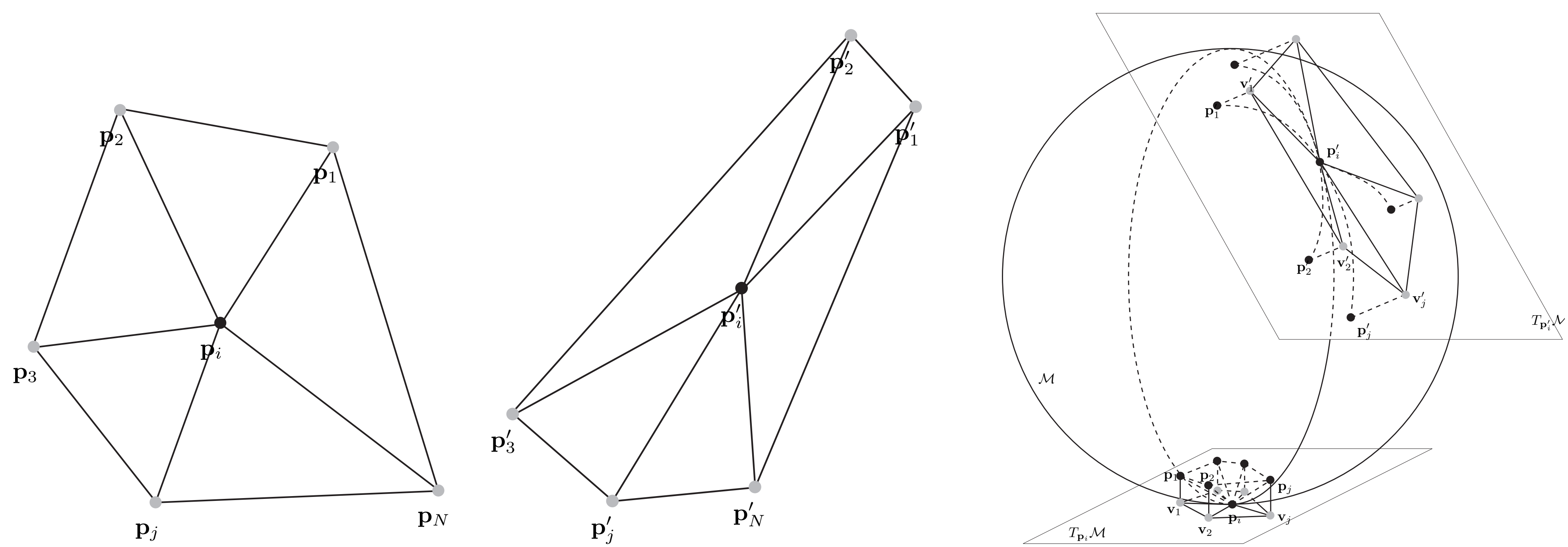
[2] cf. Bowman, 1999; Haller & Yuan, 2000; Haller, 2001; Shadden, Lekien, Marsden, 2005

Invariant manifold analogs: FTLE-LCS approach

- We can define the FTLE for Riemannian manifolds[3]

$$\sigma_t^T(x) = \frac{1}{|T|} \ln \left\| \mathrm{D}\phi_t^{t+T} \right\| \doteq \frac{1}{|T|} \log \left(\max_{\mathrm{y} \neq 0} \frac{\left\| \mathrm{D}\phi_t^{t+T}(\mathrm{y}) \right\|}{\|\mathrm{y}\|} \right)$$

with y a small perturbation in the tangent space at x.



[3]Lekien & Ross [2010] Chaos

Transport barriers: LCS

- Ridges correspond to dynamical barriers[3] or Lagrangian coherent structures (LCS): repelling surfaces for $T > 0$, attracting for $T < 0$

cylinder

Moebius strip

Each frame has a different initial time t

[3] Lekien & Ross [2010] Chaos

Atmospheric flows: Antarctic polar vortex

ozone data

Atmospheric flows: Antarctic polar vortex

ozone data + LCSs (red = repelling, blue = attracting)

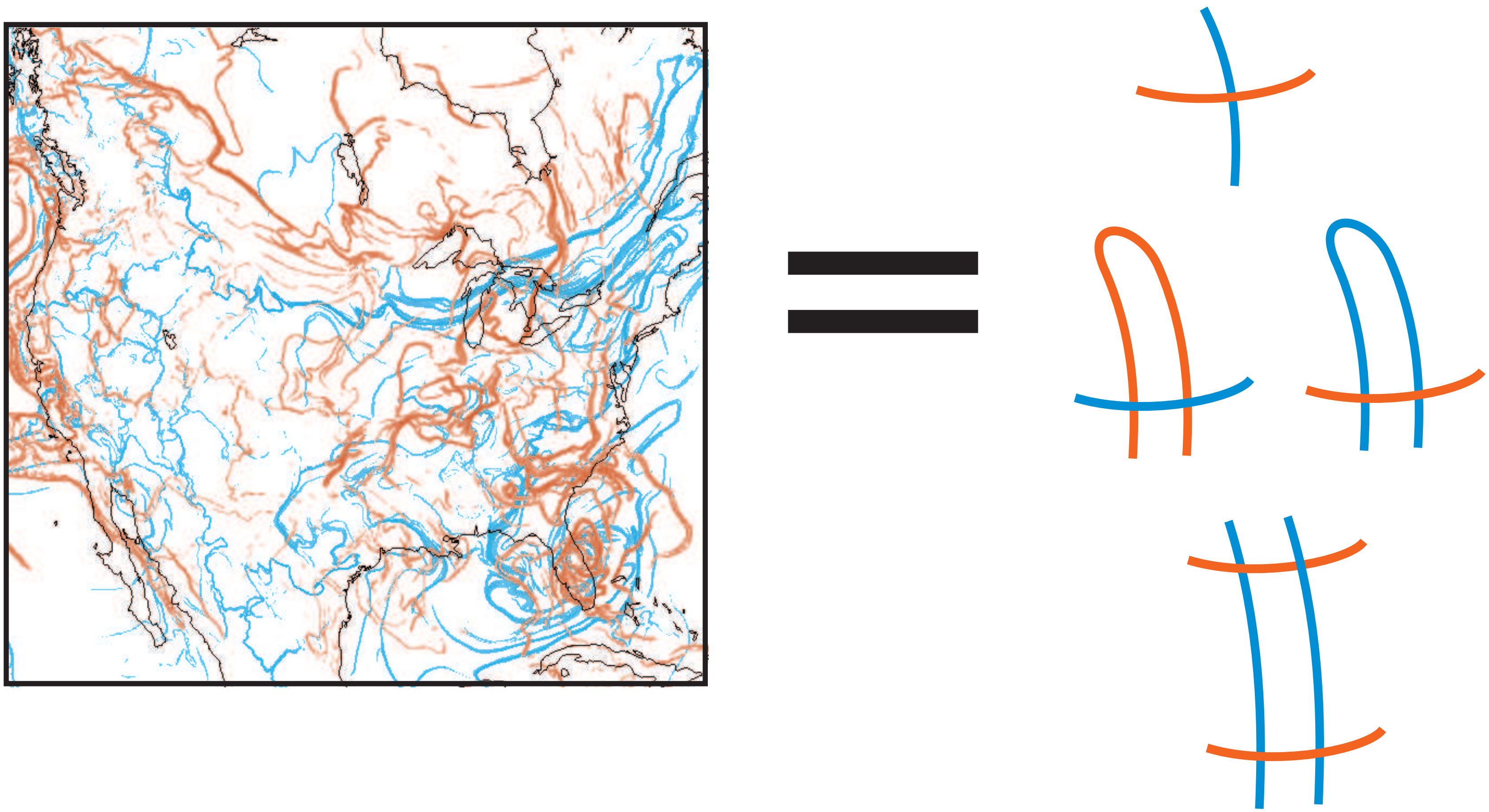
Atmospheric flows: Antarctic polar vortex

air masses on either side of a repelling LCS

Atmospheric flows: continental U.S.

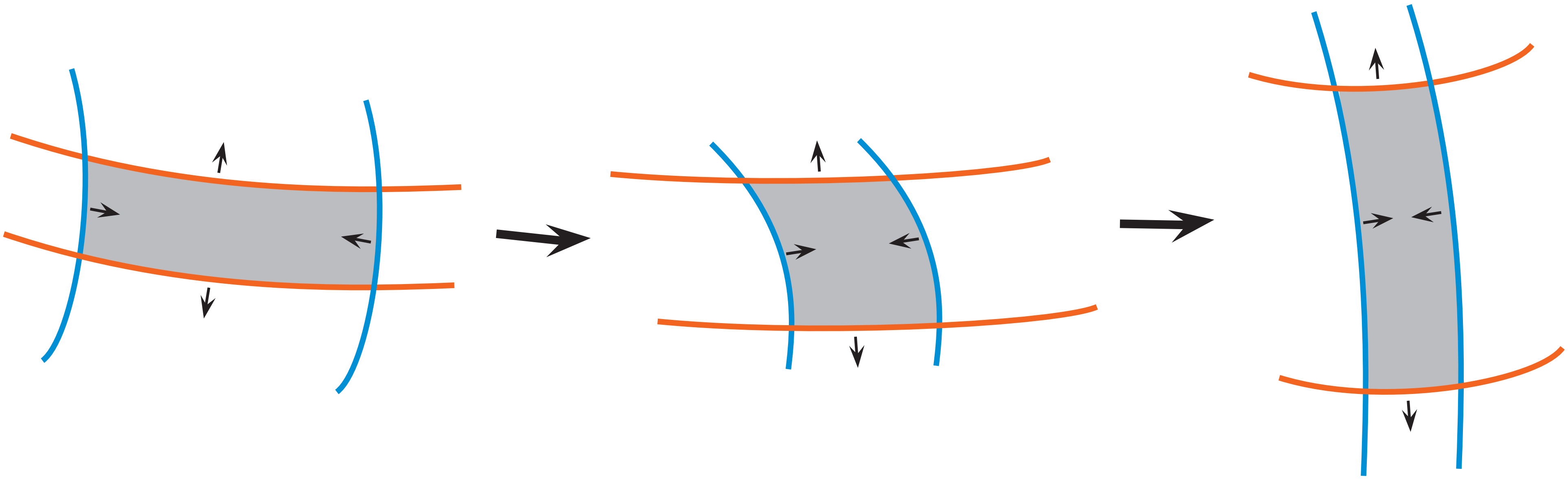
LCSs: orange = repelling, blue = attracting

Classification of motifs



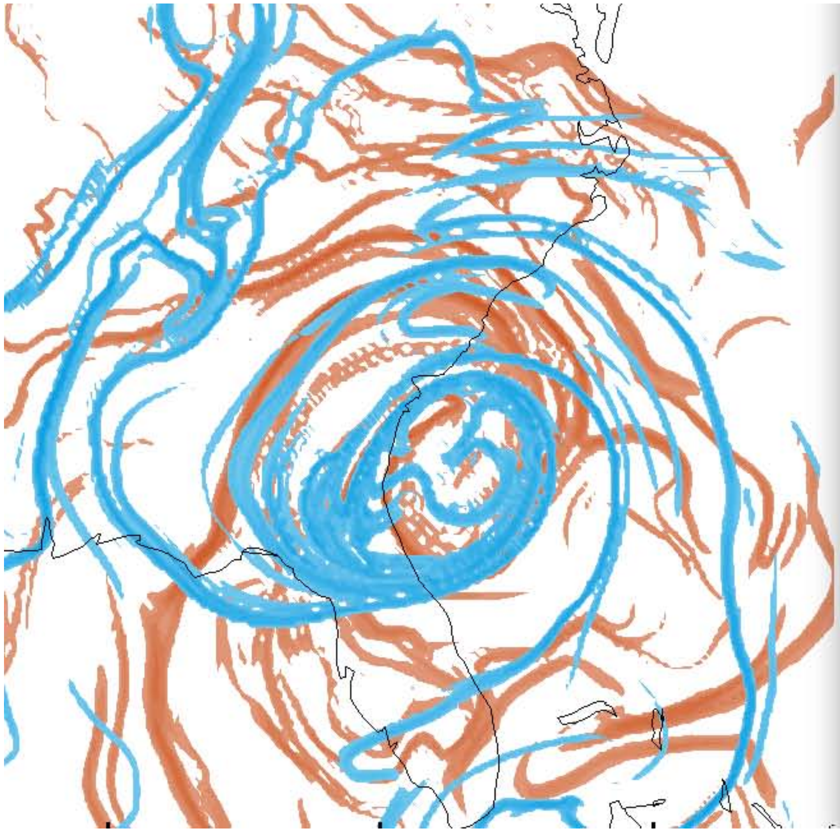
- **Regions bounded by attracting and repelling curves**
- **Atmosphere is naturally parsed into discrete 'cells'**

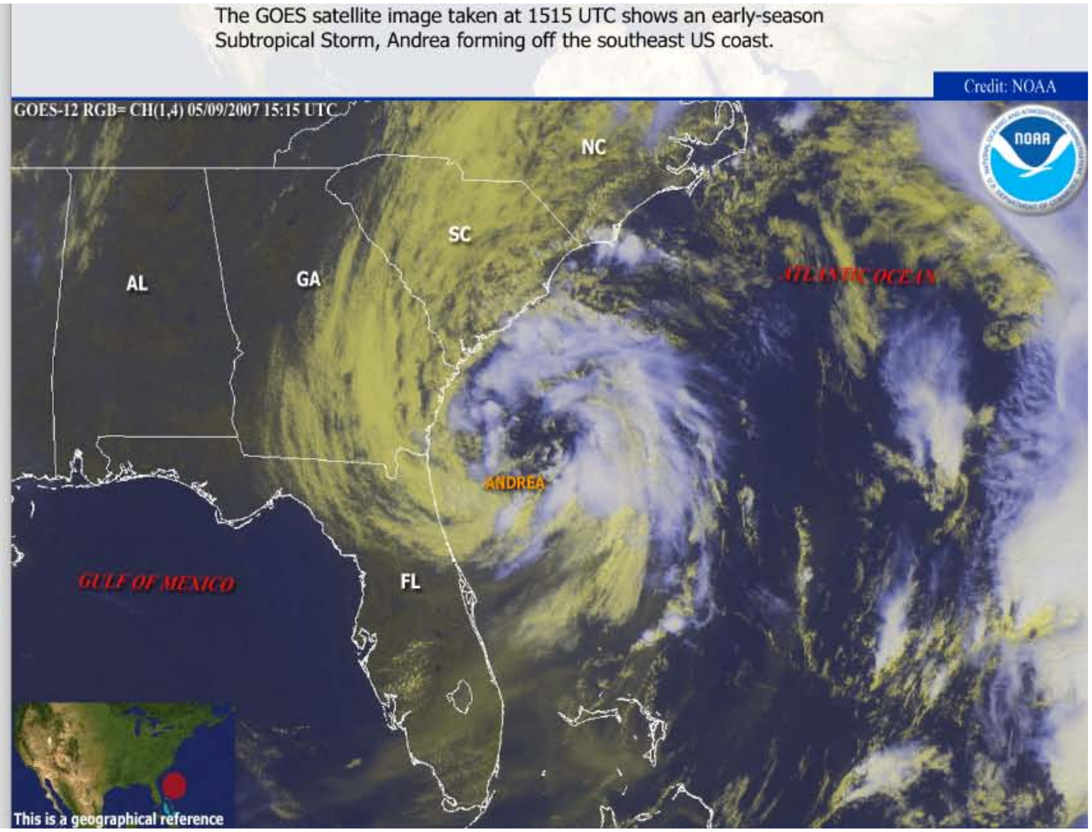
Motion of ‘cells’



- **Packets have their own dynamics as consequence of repelling and attracting natures of boundaries**

Atmospheric flows and lobe dynamics





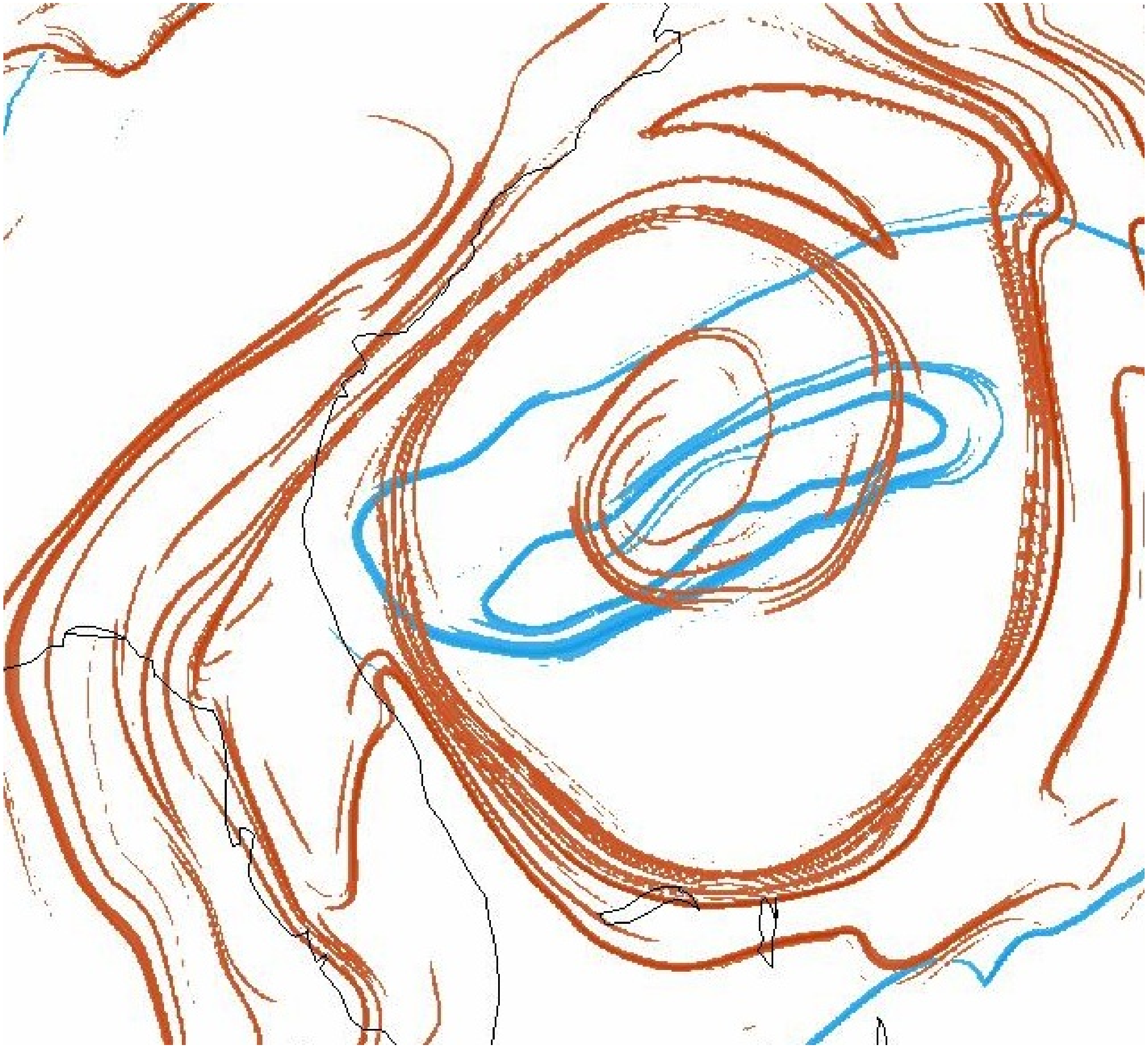
orange = repelling LCSs, blue = attracting LCSs

satellite

Hurricane Andrea, 2007

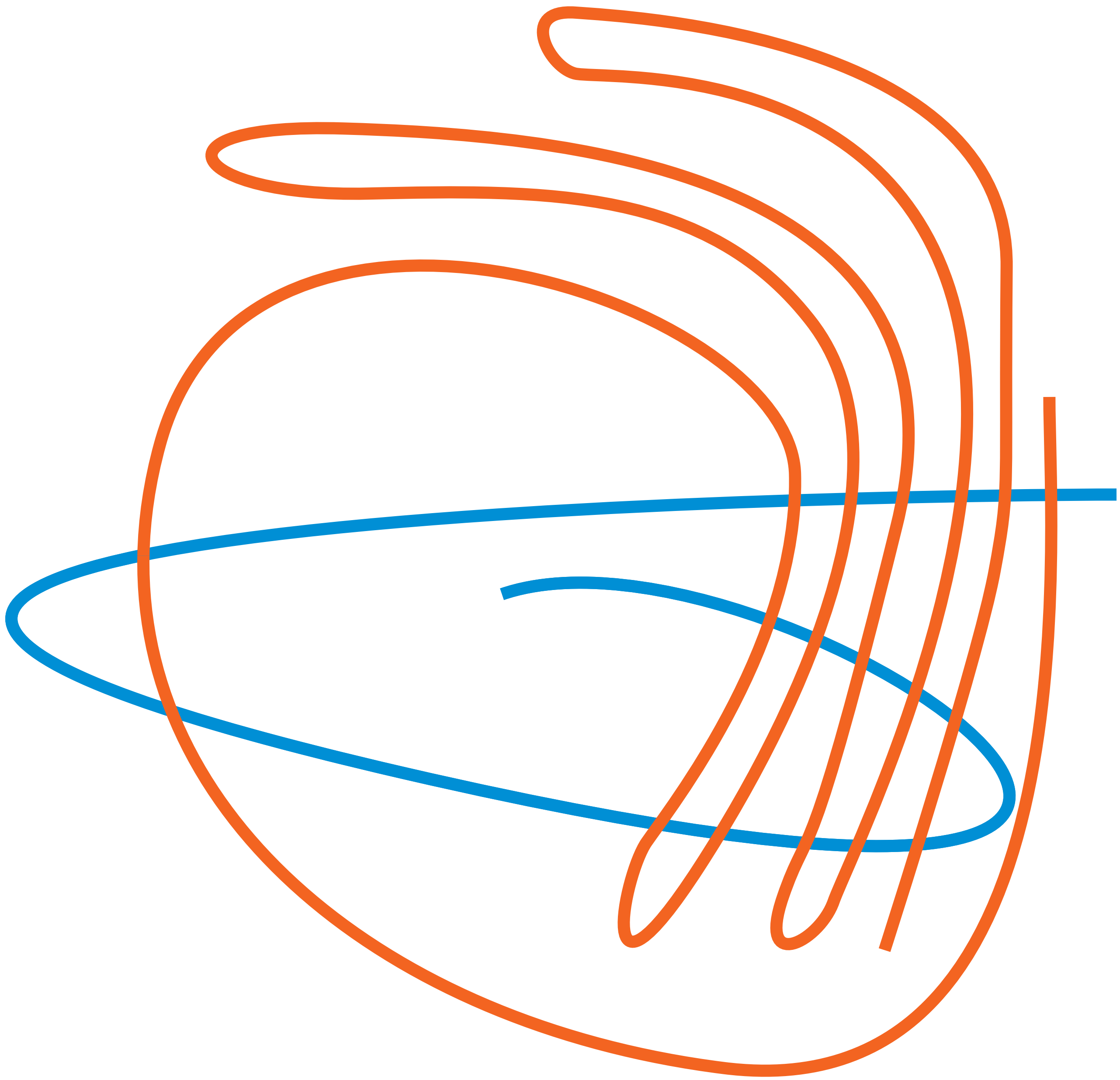
cf. Sapsis & Haller [2009], Du Toit & Marsden [2010], Lekien & Ross [2010]

Atmospheric flows and lobe dynamics



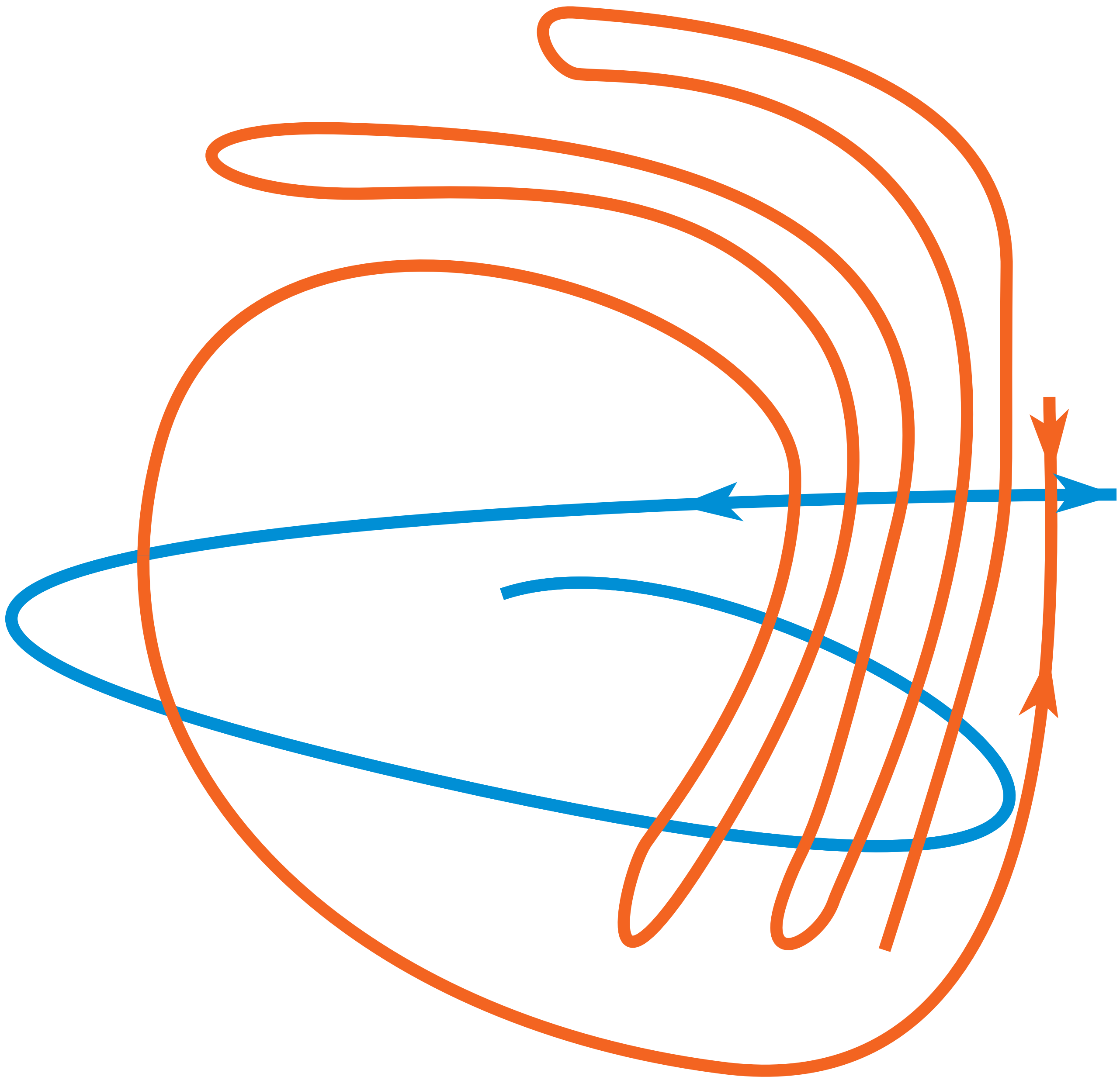
Hurricane Andrea at one snapshot; LCS shown (orange = repelling, blue = attracting)

Atmospheric flows and lobe dynamics



orange = repelling (stable manifold), blue = attracting (unstable manifold)

Atmospheric flows and lobe dynamics



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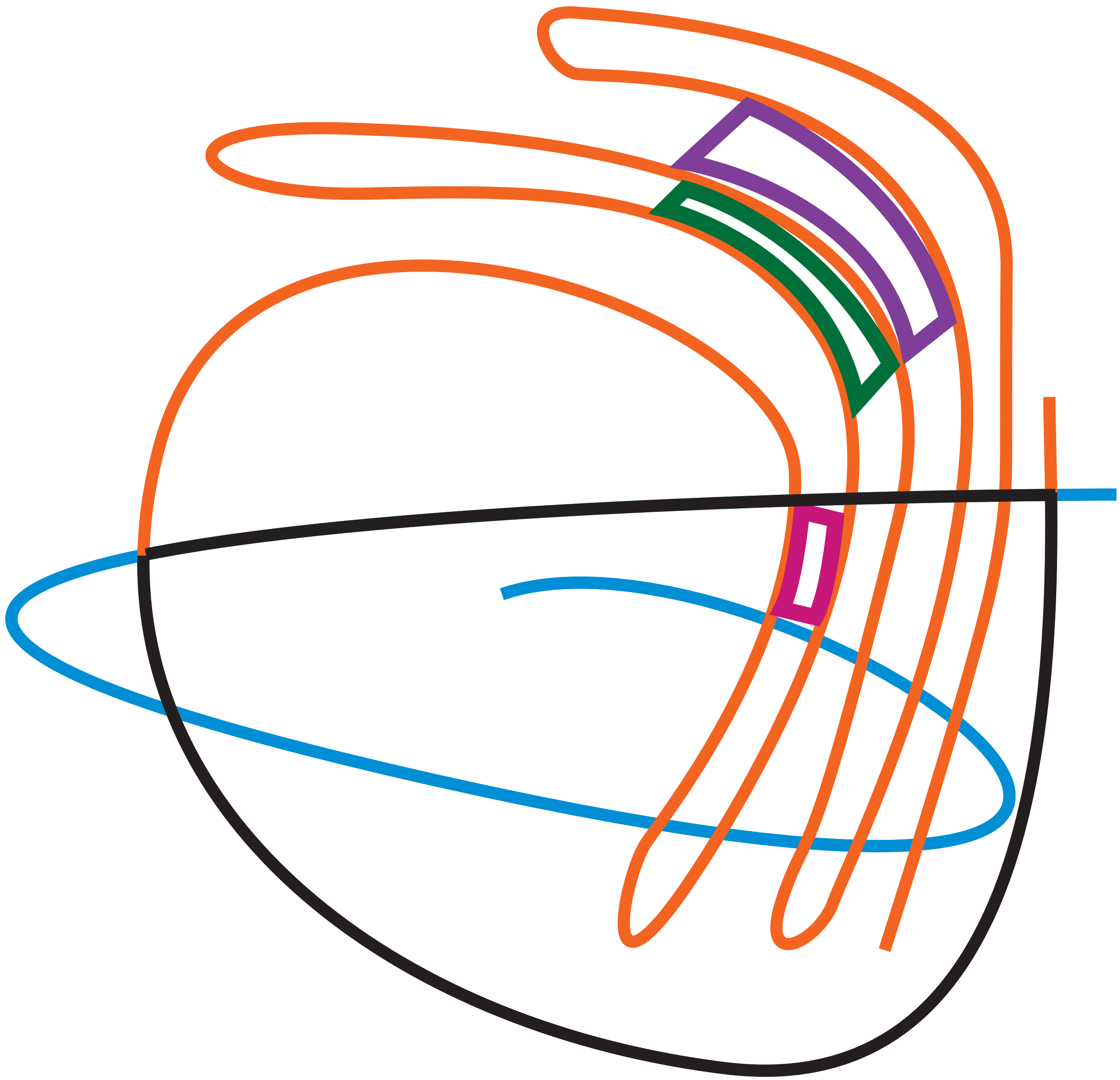
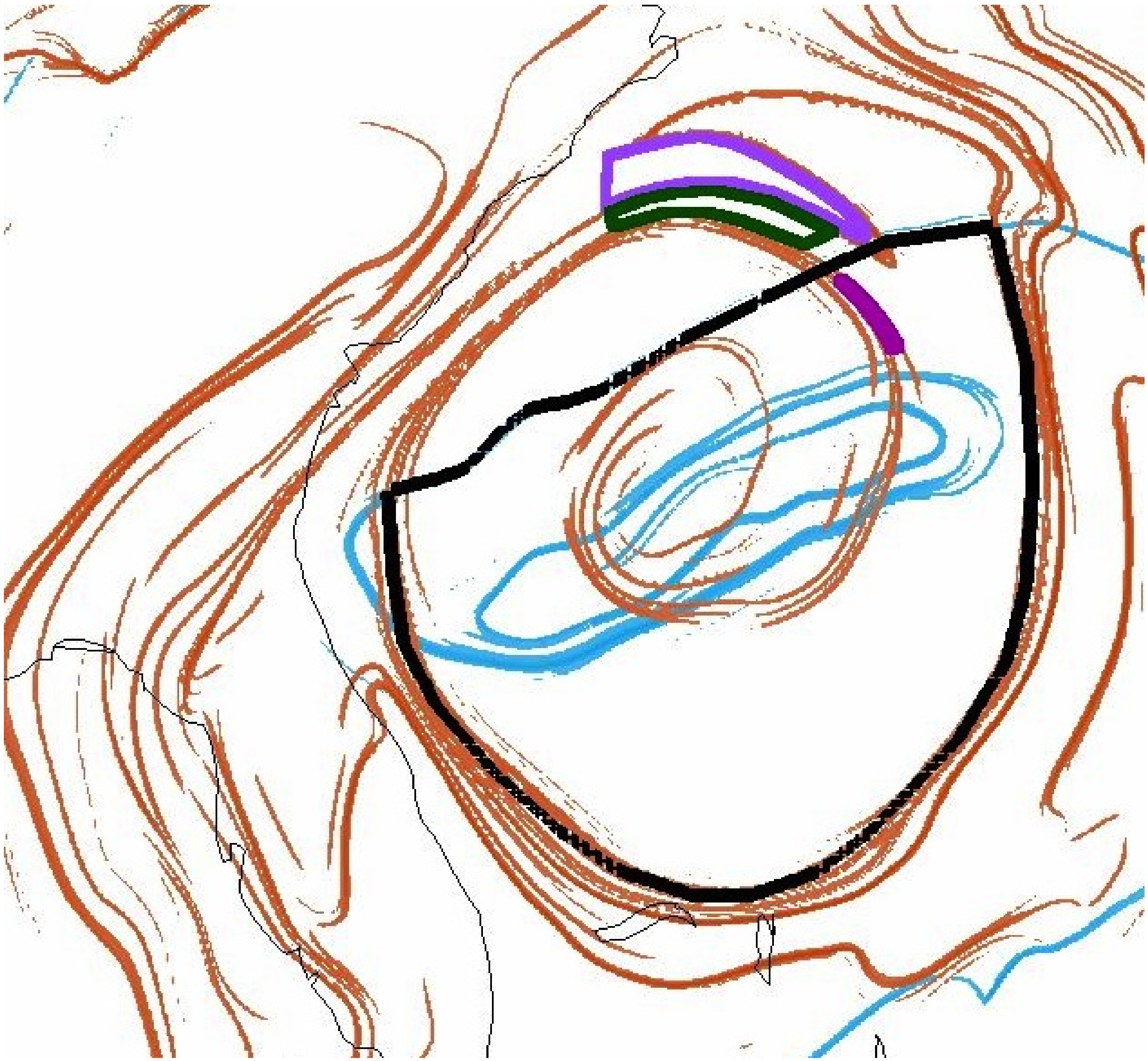
Atmospheric flows and lobe dynamics



Portions of lobes colored; magenta = outgoing, green = incoming, purple = stays out

Atmospheric flows and lobe dynamics



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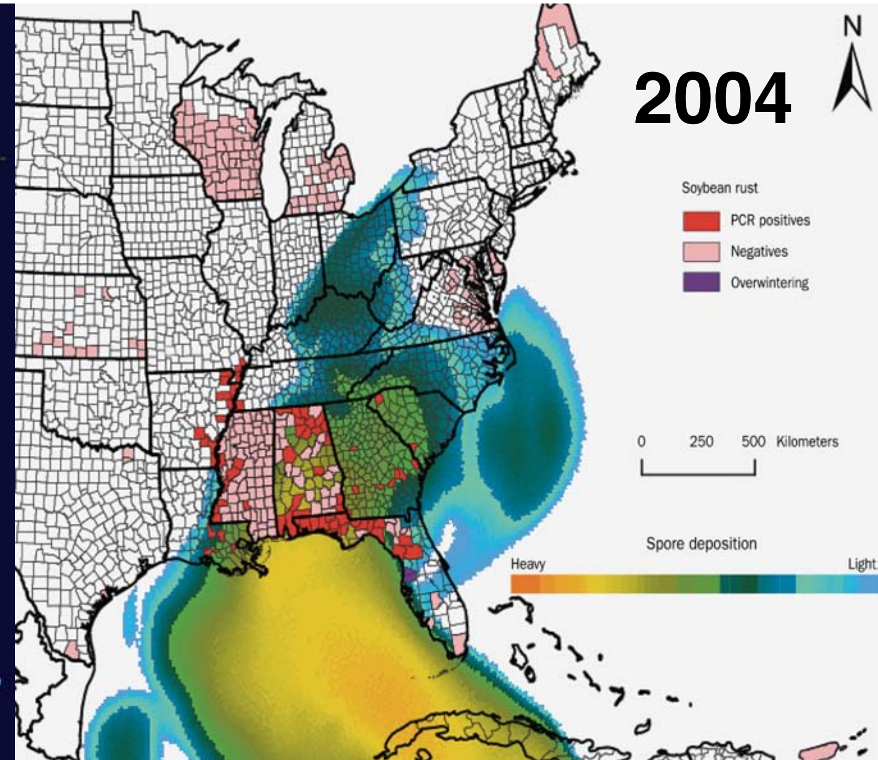
Atmospheric flows and lobe dynamics

Sets behave as lobe dynamics dictates

Invasive species riding the atmosphere

Hurricane Ivan (2004) brought new crop disease (soybean rust) to U.S.

From Rio Cauca region of Colombia



Disease extent

Invasive species riding the atmosphere

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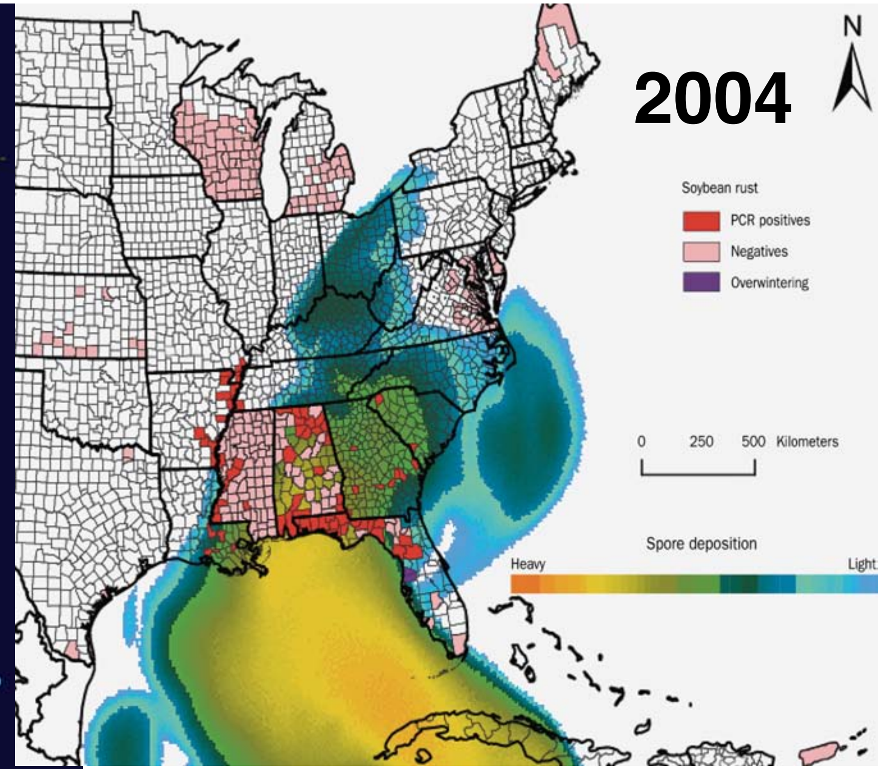
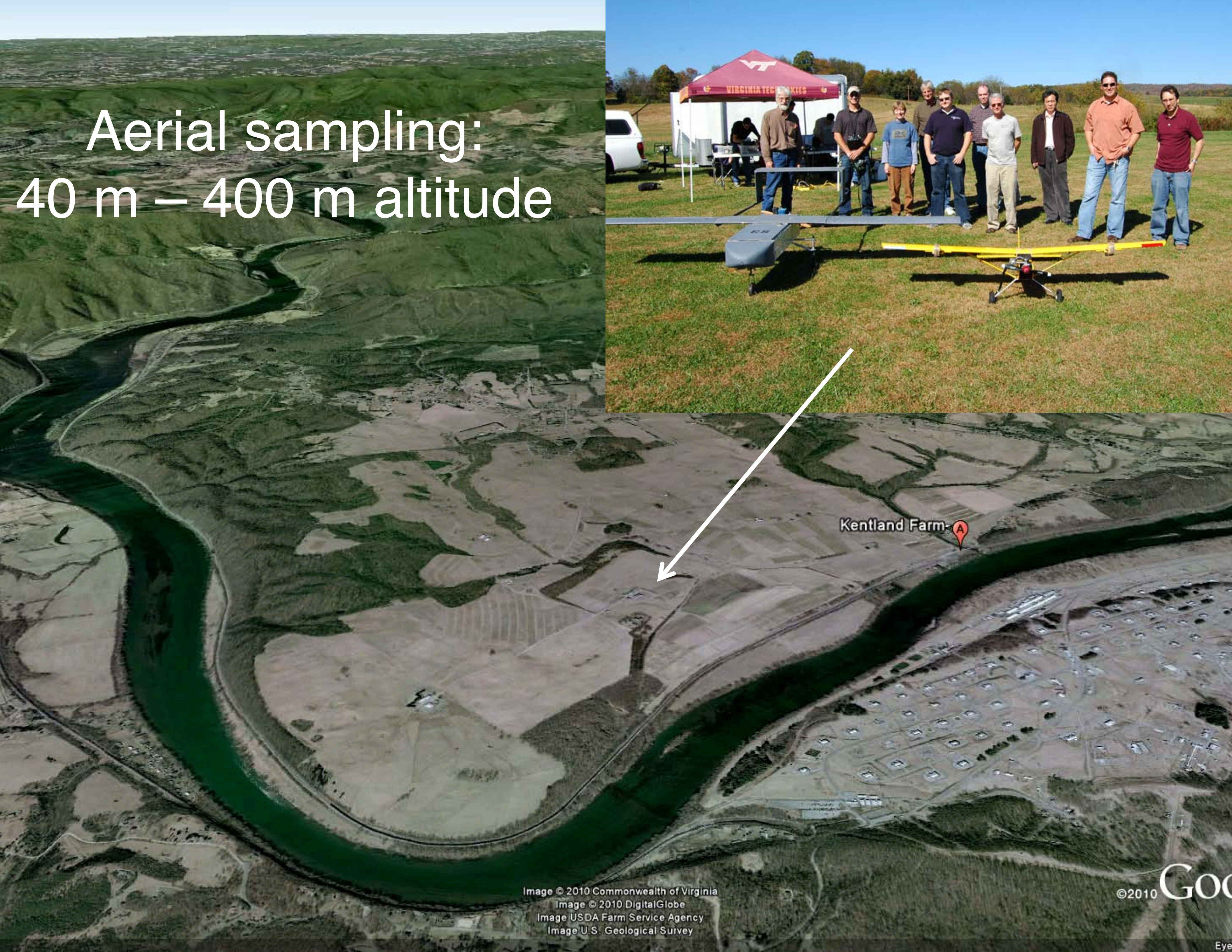
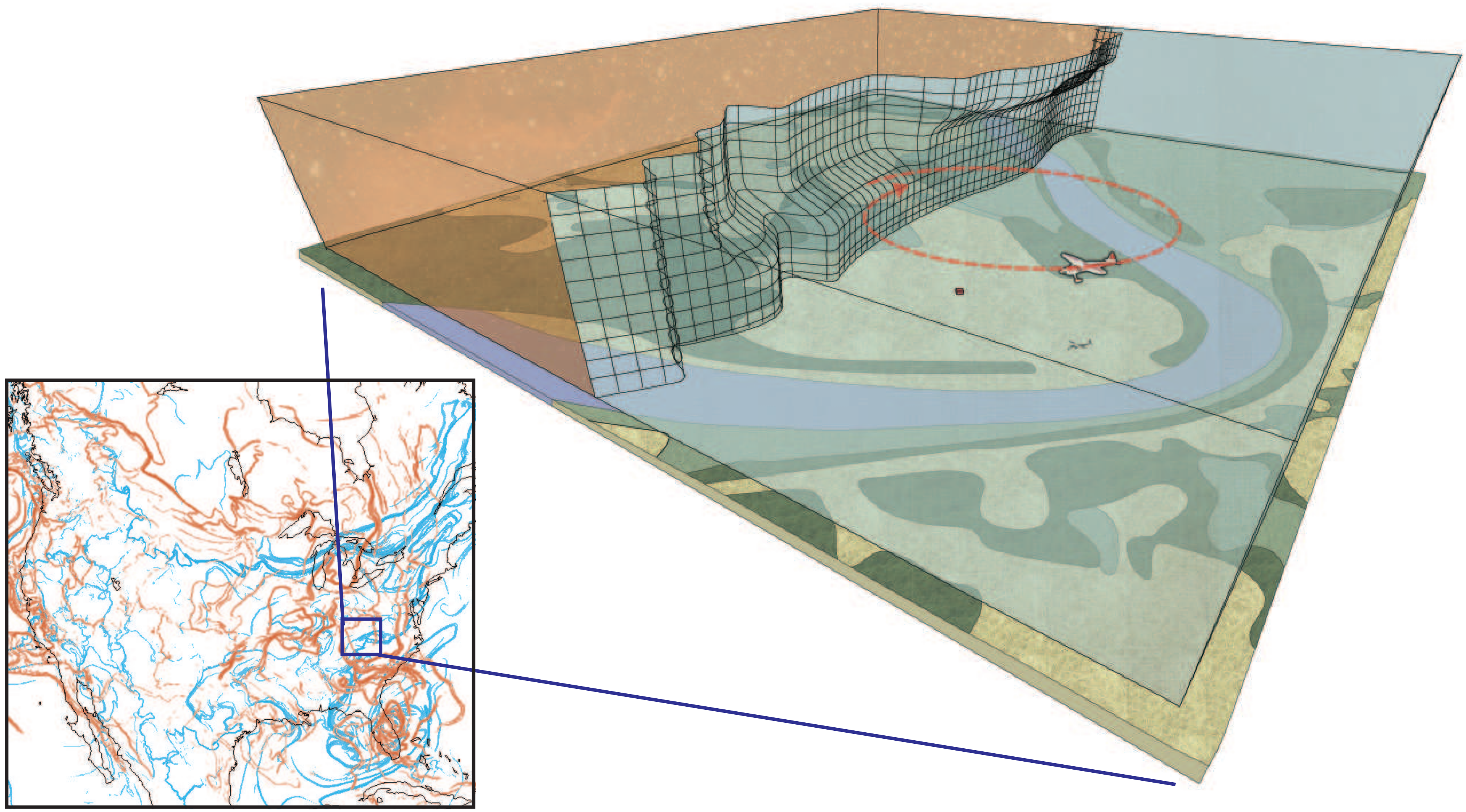


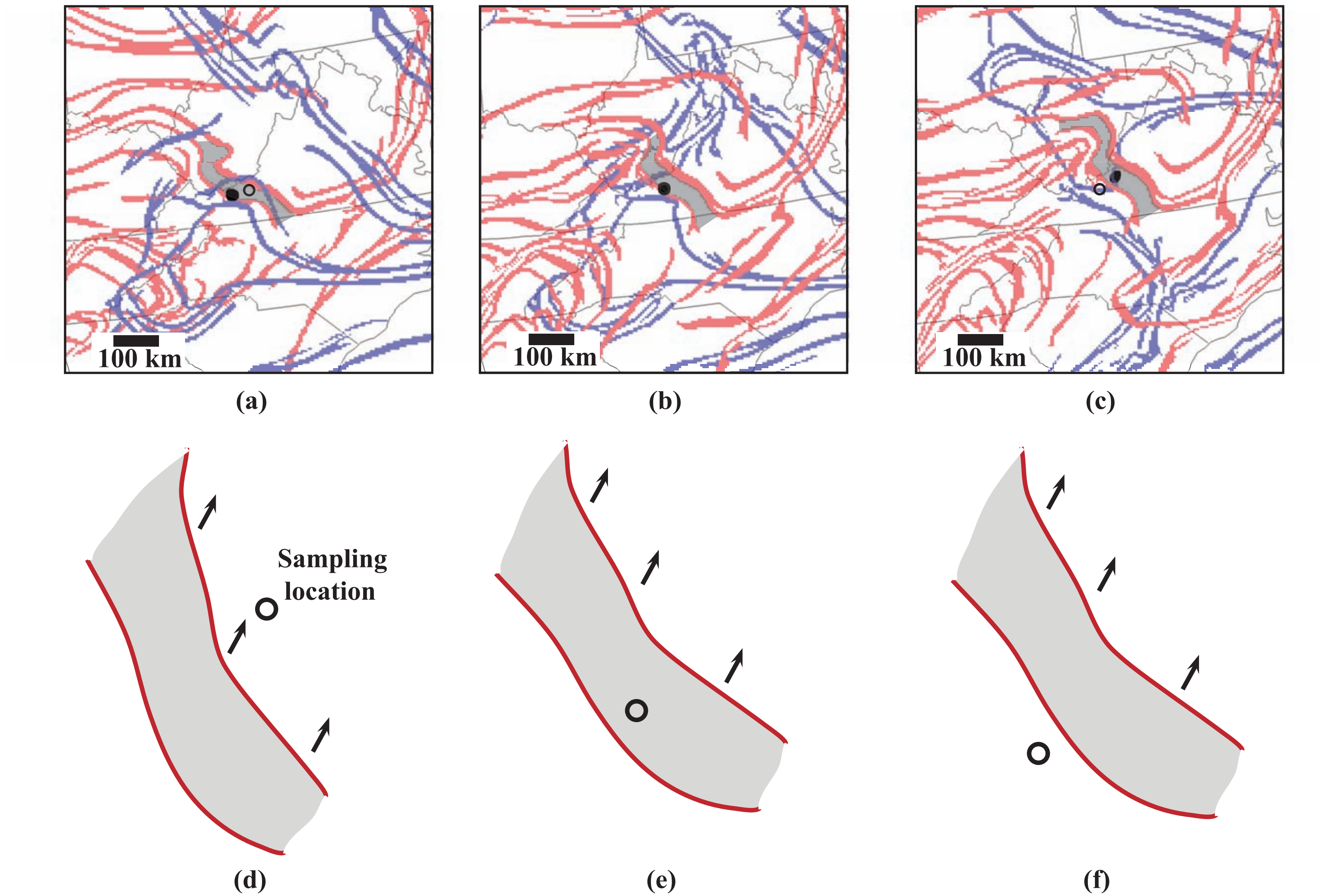
Disease extent

Cost of invasive organisms is ***$137 billion*** **per year in U.S.**


Aerial sampling:
40 m – 400 m altitude
Kentland Farm
Image © 2010 Commonwealth of Virginia
Image © 2010 DigitalGlobe
Image USDA Farm Service Agency
Image U.S. Geological Survey

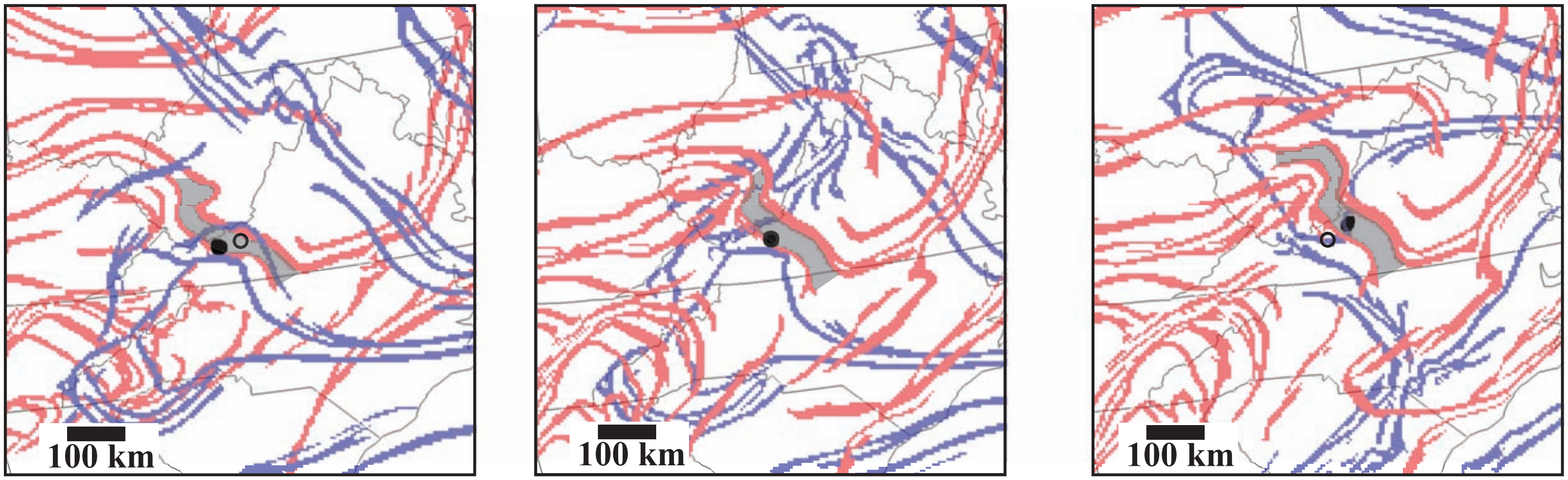


Pathogen transport: filament bounded by LCS

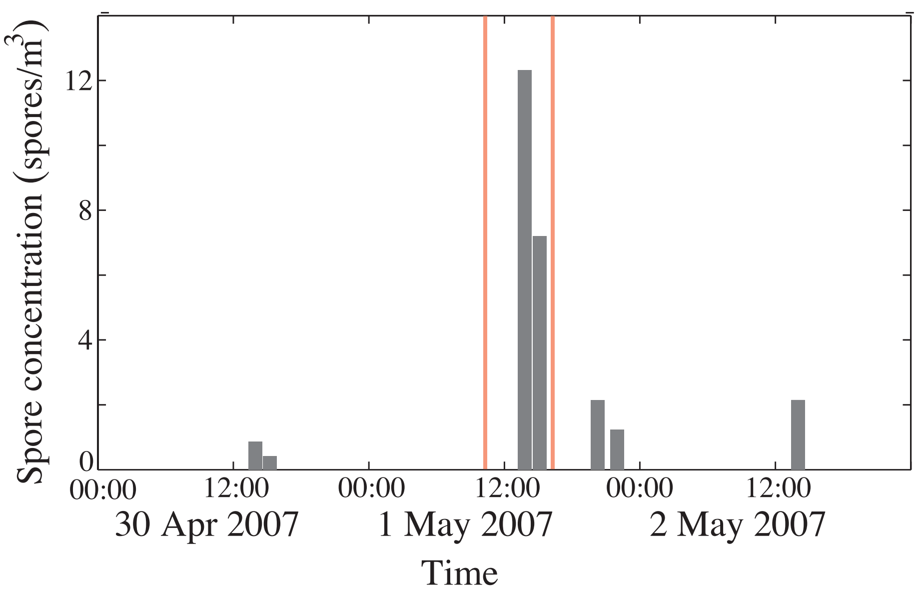


12:00 UTC 1 May 2007 15:00 UTC 1 May 2007 18:00 UTC 1 May 2007

Pathogen transport: filament bounded by LCS



(a) (b) (c)



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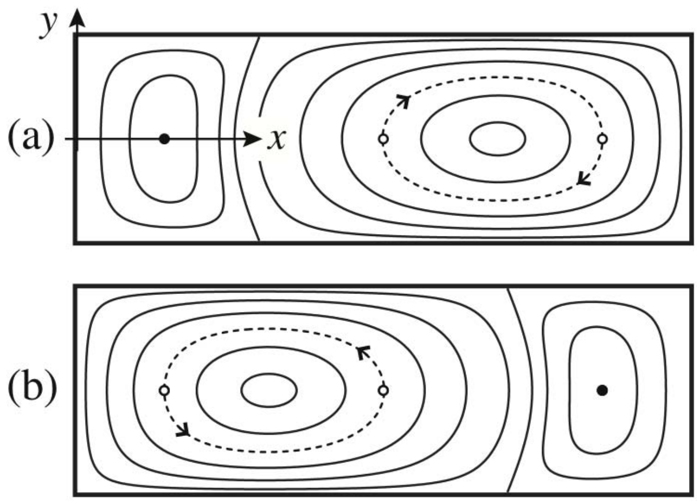
Lobe dynamics: another fluid example

- Fluid example: time-periodic Stokes flow[2]



streamlines

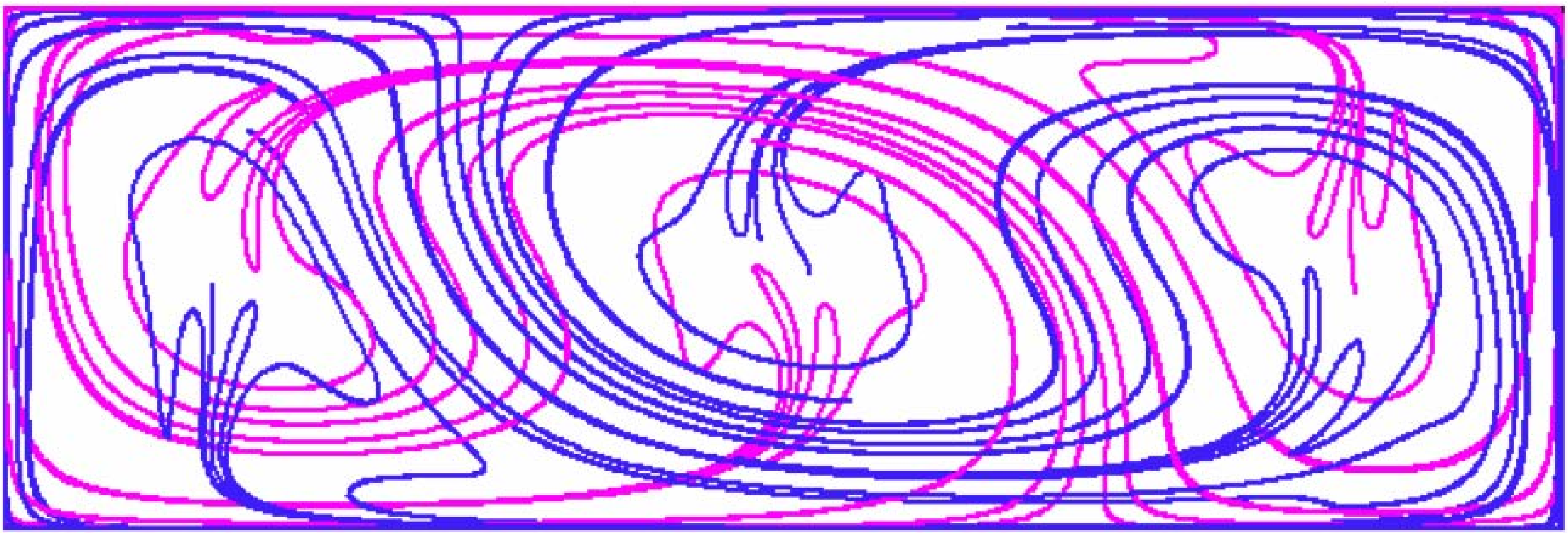
tracer blob

Lid-driven cavity flow

[2]Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Lobe dynamics: another fluid example

- Fluid example: time-periodic Stokes flow[2]



some invariant manifolds of saddles

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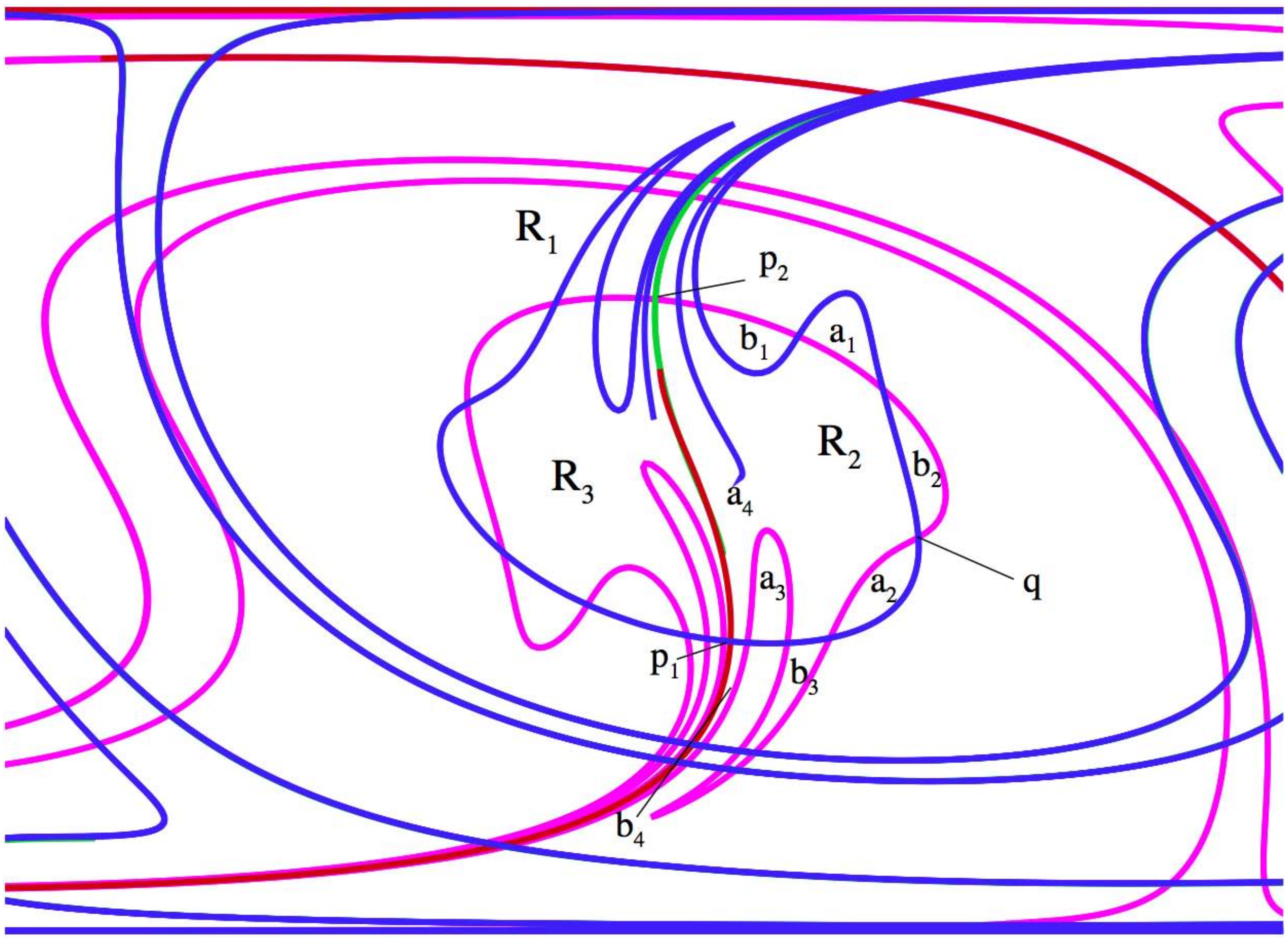

regions and lobes labeled

[2]Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Stable/unstable manifolds and lobes in fluids

- Fluid example: time-periodic Stokes flow[2]

material blob at $t = 0$

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Stable/unstable manifolds and lobes in fluids

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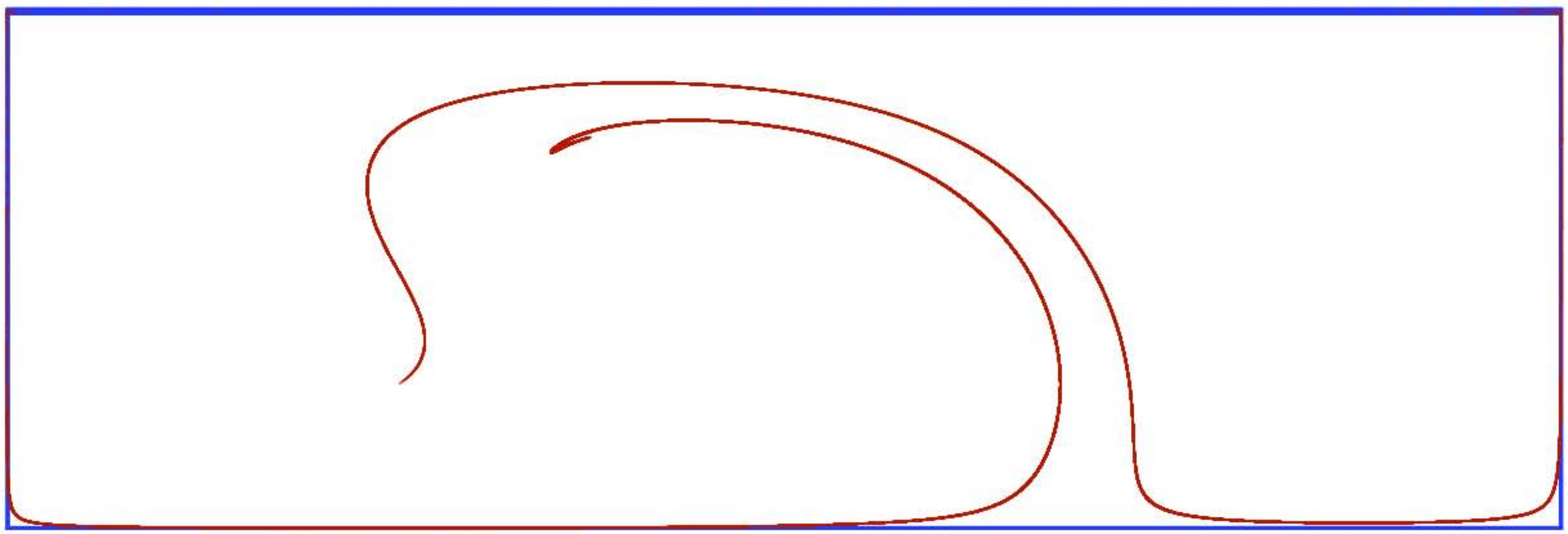
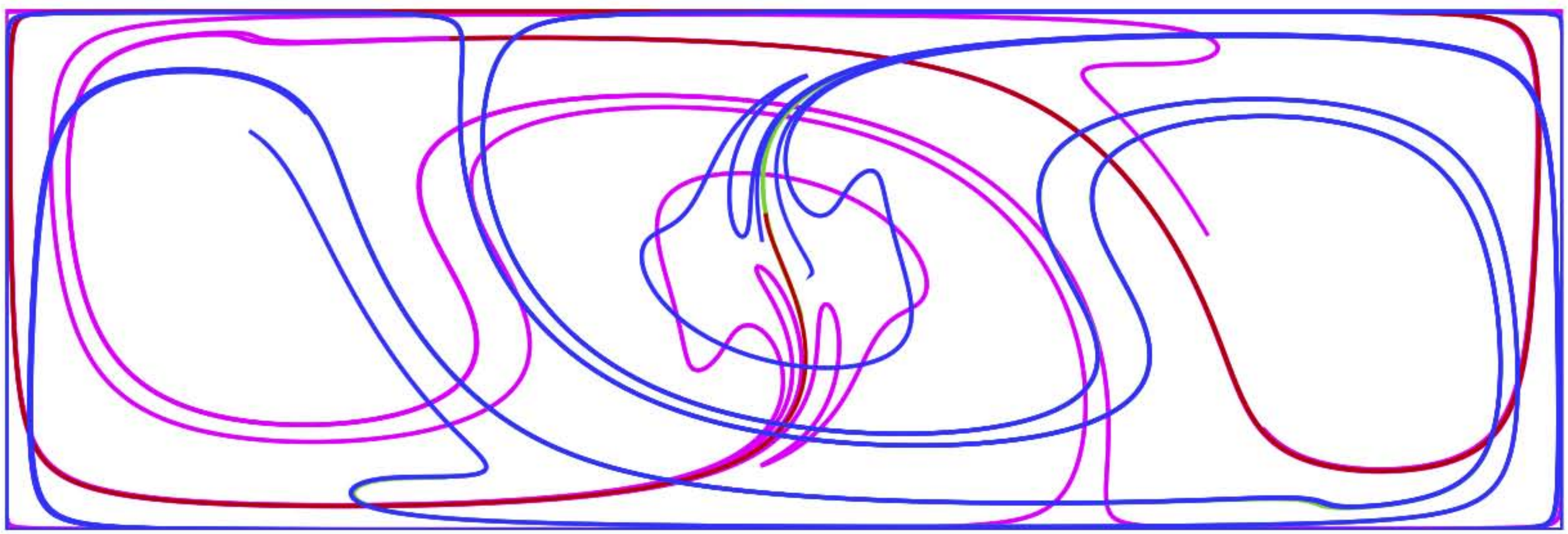


material blob at $t = 5$

[2]Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Stable/unstable manifolds and lobes in fluids

- Fluid example: time-periodic Stokes flow[2]



some invariant manifolds of saddles

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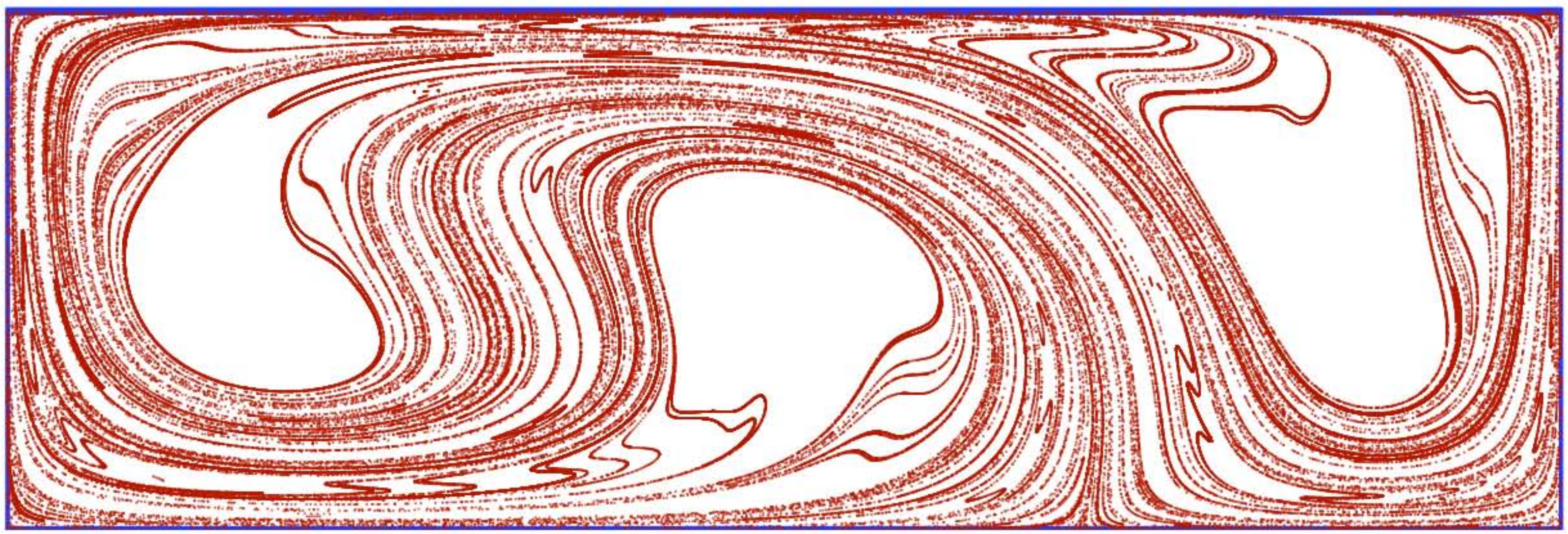
- Fluid example: time-periodic Stokes flow[2]

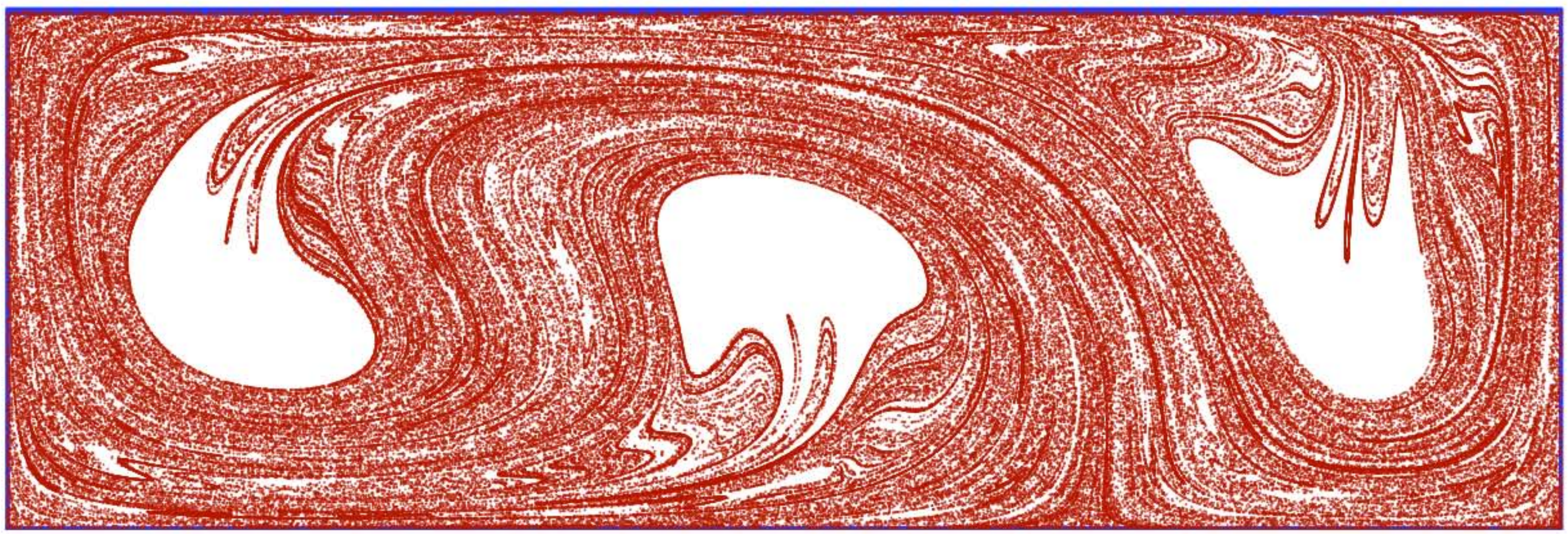


material blob at $t = 10$

[2]Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Stable/unstable manifolds and lobes in fluids

- Fluid example: time-periodic Stokes flow[2]

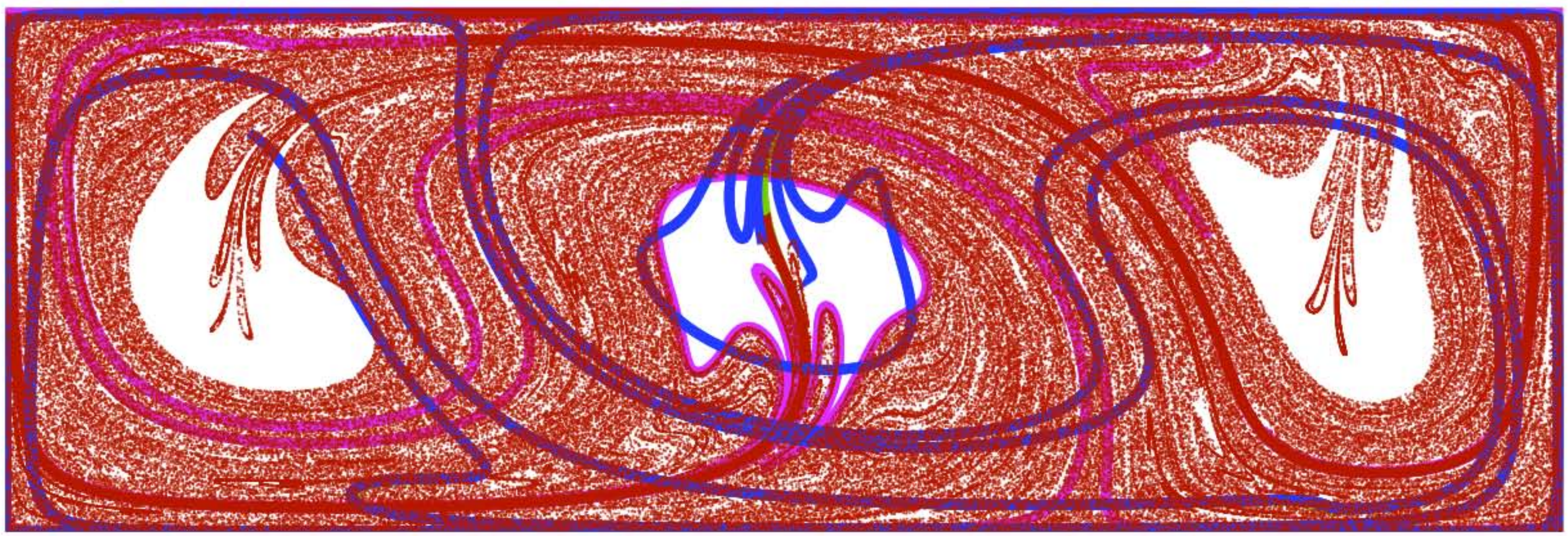


material blob at $t = 15$

[2] Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Stable/unstable manifolds and lobes in fluids

- Fluid example: time-periodic Stokes flow[2]



material blob and manifolds

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Stable/unstable manifolds and lobes in fluids

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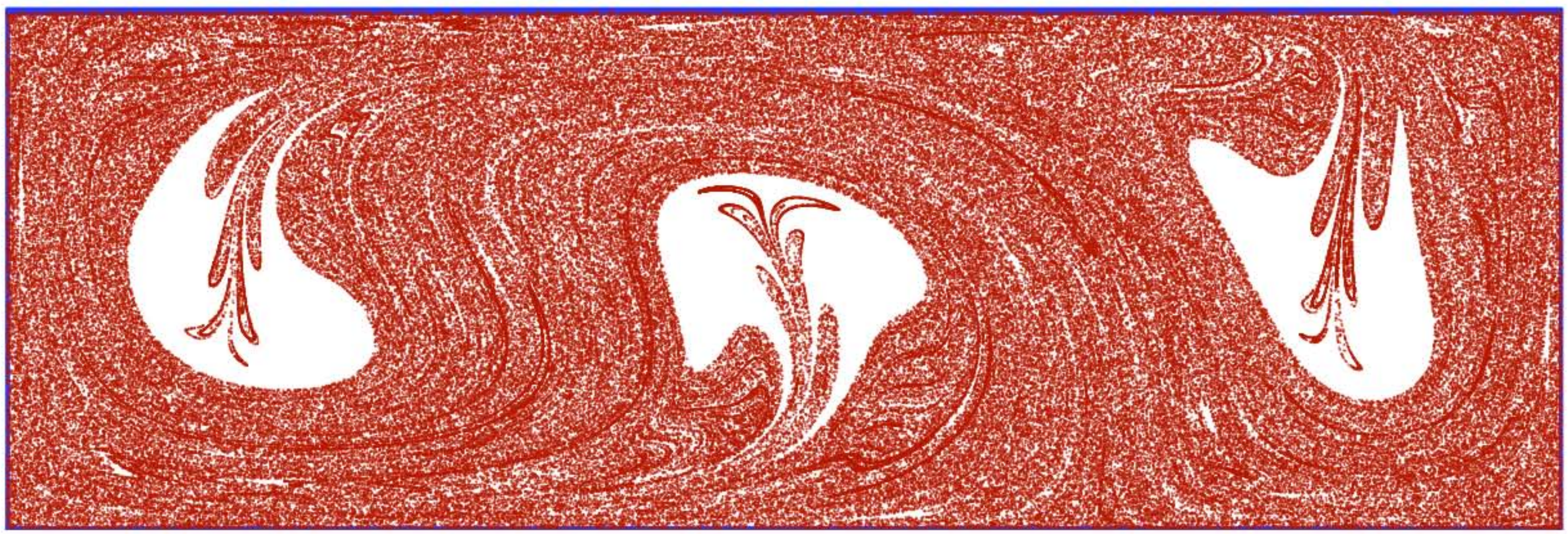

material blob at $t = 20$

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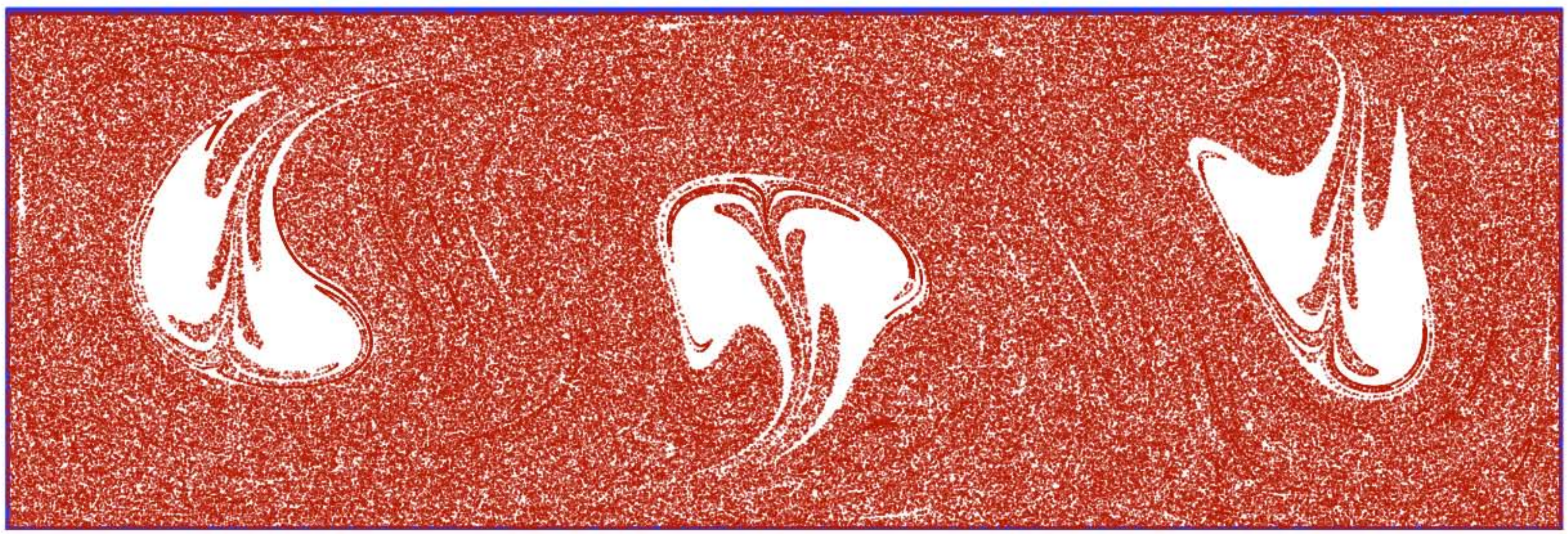
Stable/unstable manifolds and lobes in fluids

- Fluid example: time-periodic Stokes flow[2]



material blob at $t = 25$

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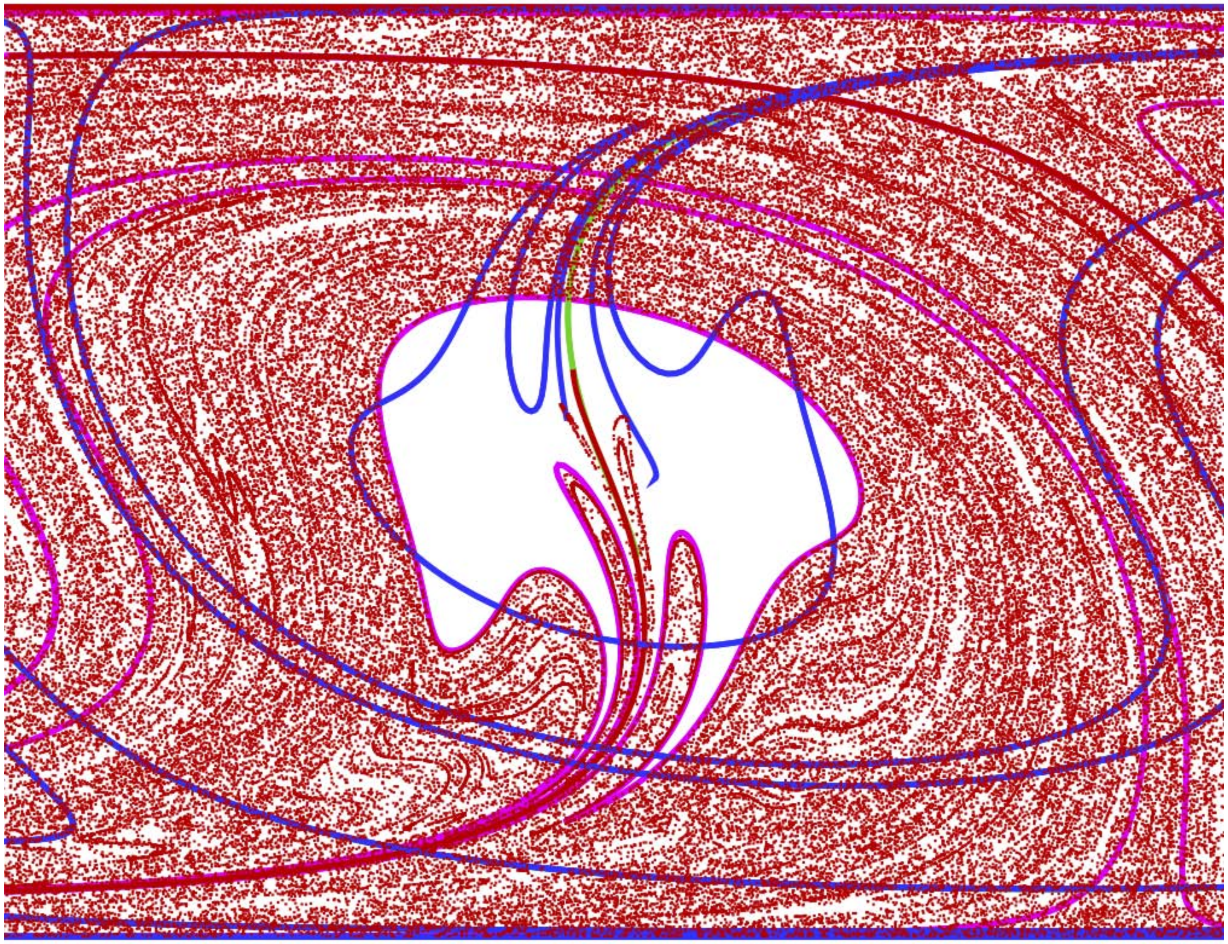
Stable/unstable manifolds and lobes in fluids

- Fluid example: time-periodic Stokes flow[2]

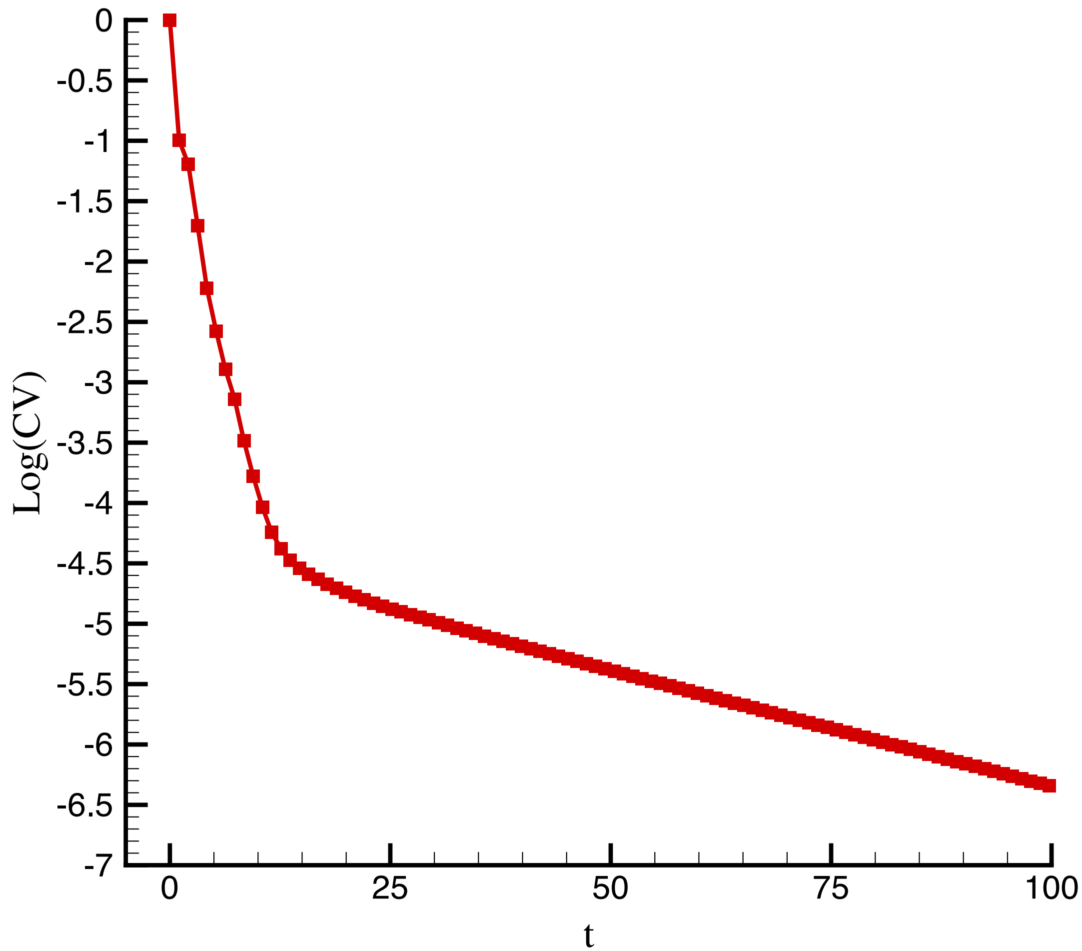


- Saddle manifolds and lobe dynamics provide template for motion

[2]Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Stable/unstable manifolds and lobes in fluids

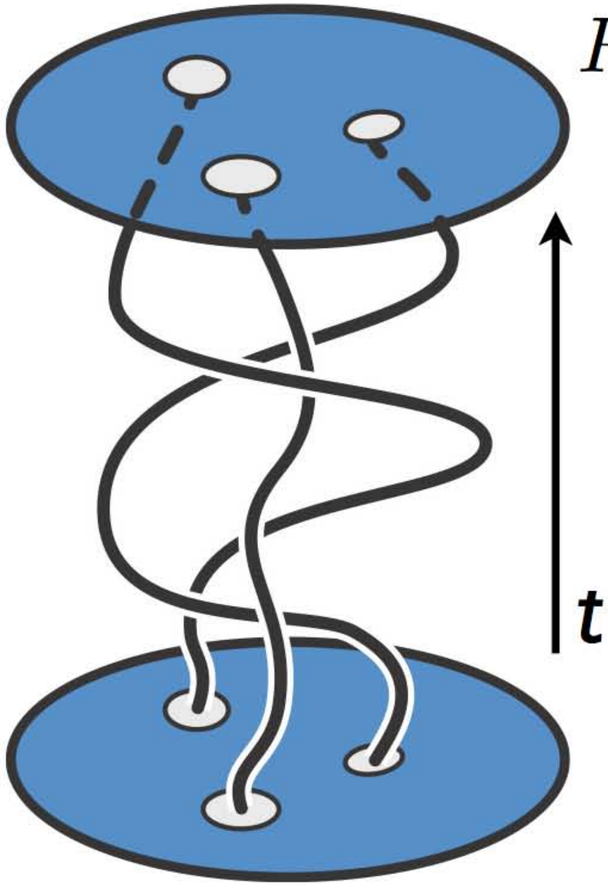
- Fluid example: time-periodic Stokes flow[2]



- Homogenization has two exponential rates: slower one related to lobes

[2]Computations of Mohsen Gheisarieha and Mark Stremler (Virginia Tech)

Braiding of stirrers



R_N : 2D fluid region with *N* stirring ‘rods’

- stirrers move on periodic orbits
- stirrers = solid objects or *fluid particles*
- stirrer motions generate diffeomorphism
  $f : R_N \to R_N$
- stirrer trajectories generate braids
  in 2+1 dimensional space-time

Thurston-Nielsen classification theorem

- Thurston (1988) Bull. Am. Math. Soc.
- A stirrer motion f is isotopic to a stirrer motion g of one of three types

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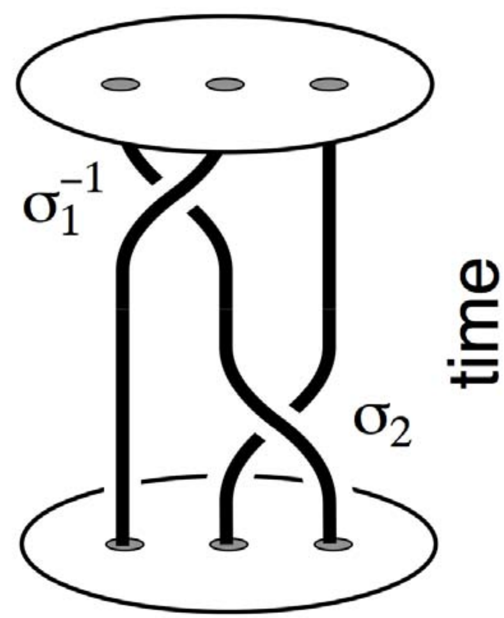
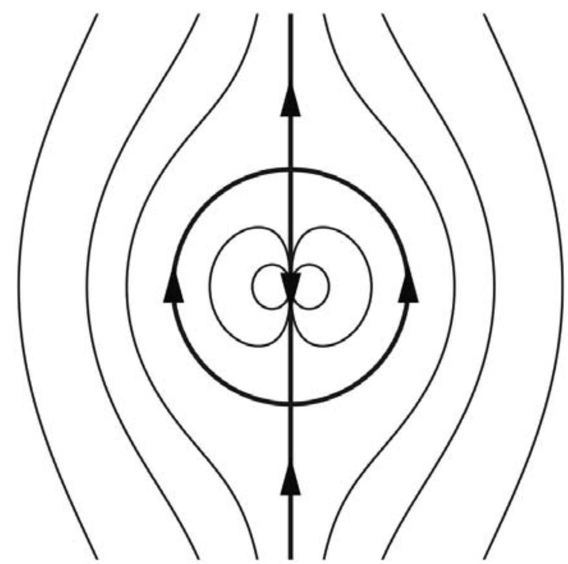
  (iii) reducible: g contains both f.o. and pA regions
- h_{TN} computed from 'braid word', e.g., $\sigma_{-1}\sigma_2$
- $\log(\lambda_{PF}(A))$ provides a **lower bound** on the true topological entropy
- i.e., non-trivial material lines grow like $\ell \sim \ell_0 \lambda^n$, where $\lambda \geq \lambda_{\mathrm{TN}}$



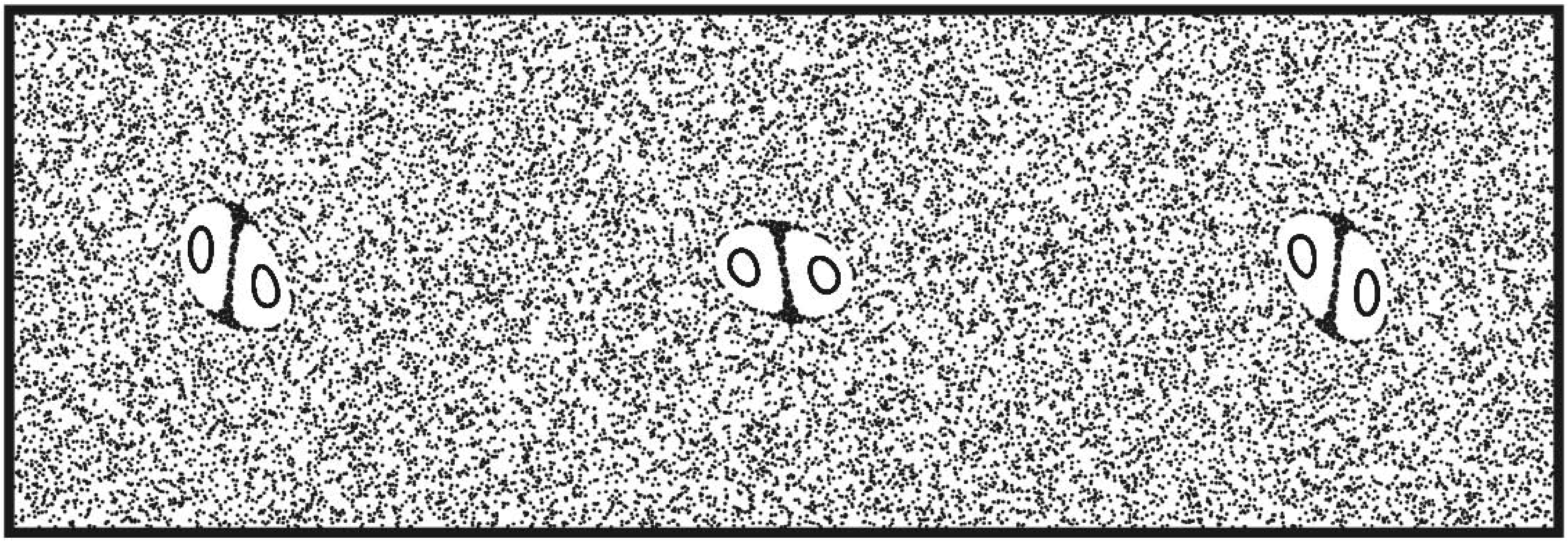
Identifying 'ghost rods': periodic points

tracer blob for $\tau_f > 1$

- Bifurcation parameter τ_f to this system
- Critical point $\tau_f^* = 1$
- For $\tau_f > 1$, pairs of elliptic and saddle points
- Below $\tau_f < 1$, pairs vanish



Identifying ‘ghost rods’: periodic points



Poincaré section for $\tau_f > 1$

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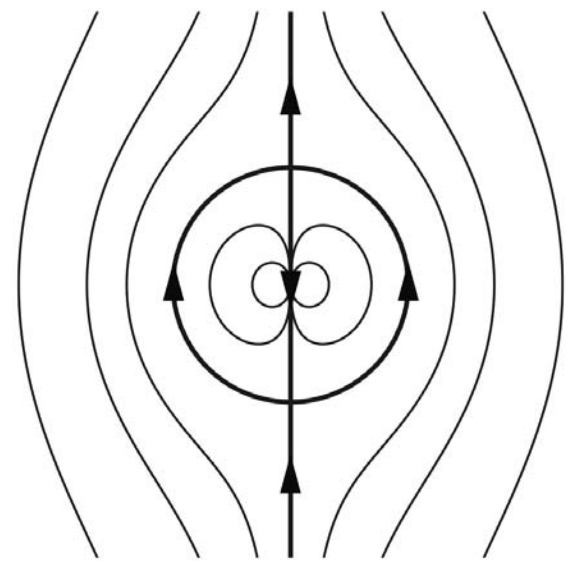
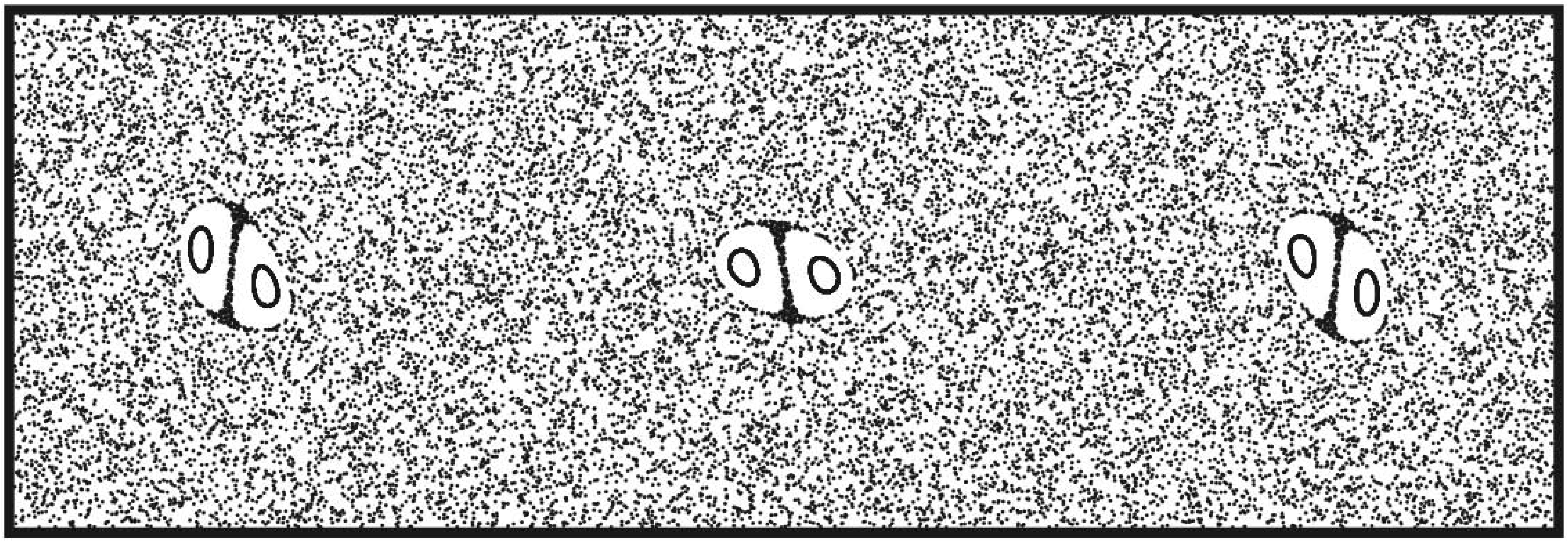
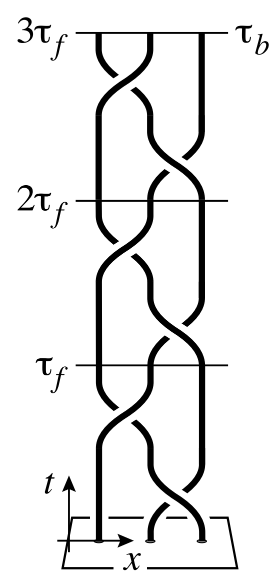


Identifying 'ghost rods': periodic points

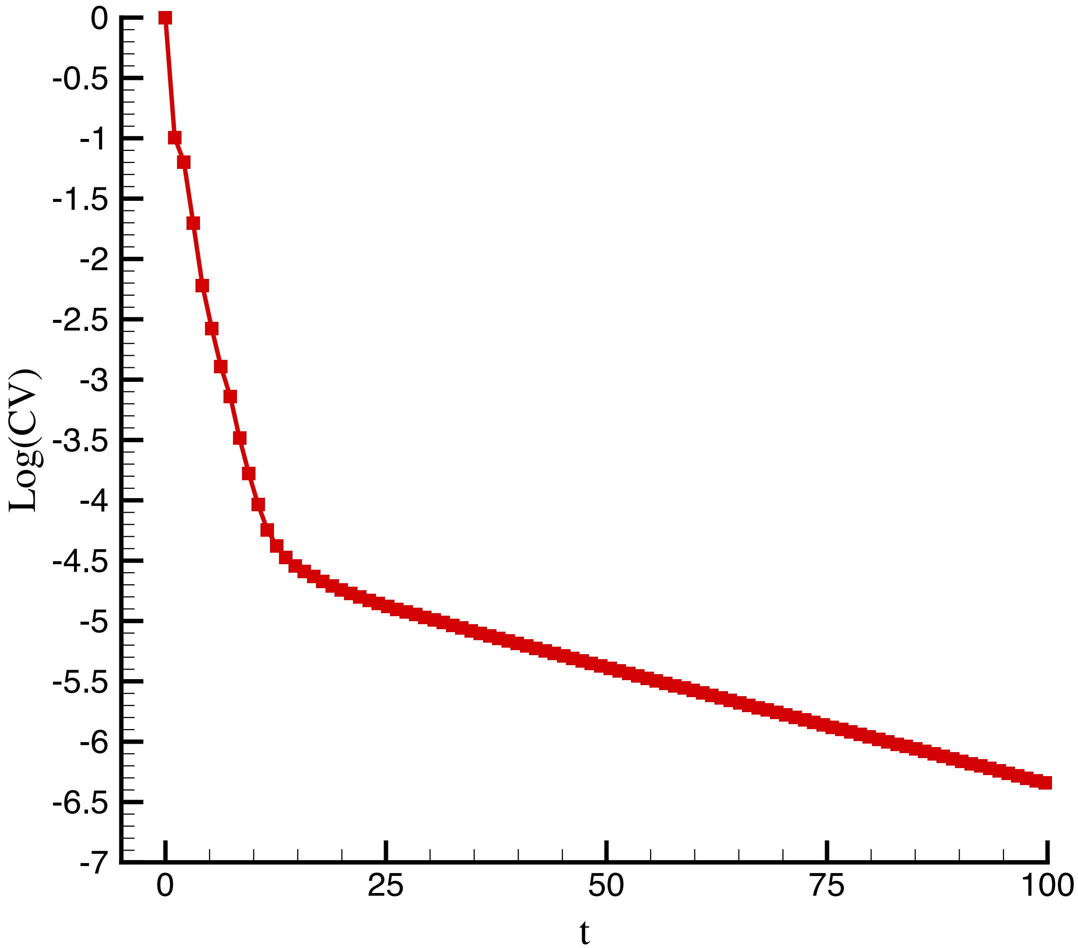


Poincaré section for $\tau_f > 1$

- Periodic points of period 3 $\Rightarrow$ act as 'ghost rods'
- Their braid $\Rightarrow h_{\rm TN} = 0.96242$
- Actual $h_{\rm flow} \approx 0.964$
- $h_{\rm TN}$ is an excellent lower bound

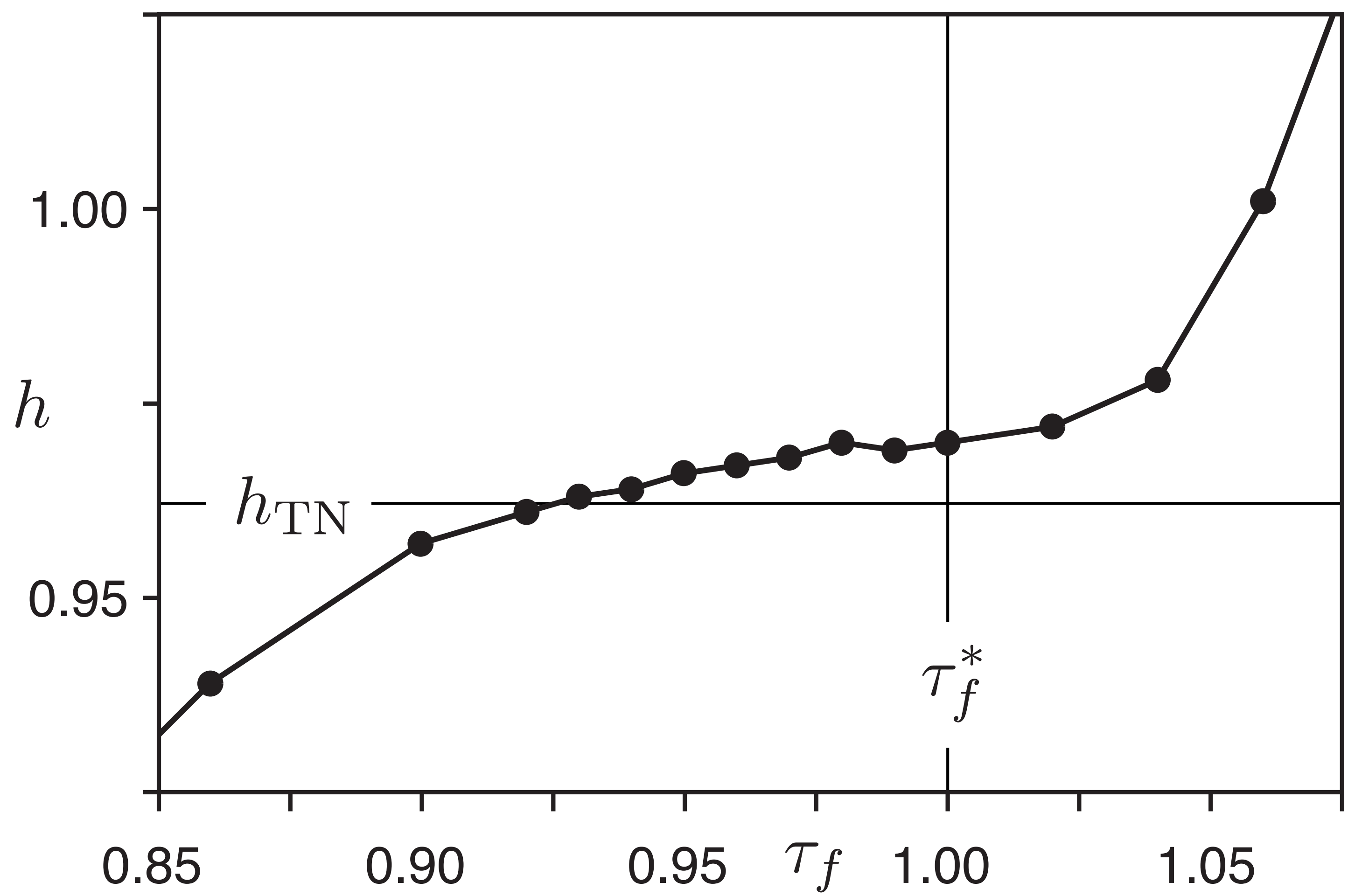


Identifying 'ghost rods': periodic points



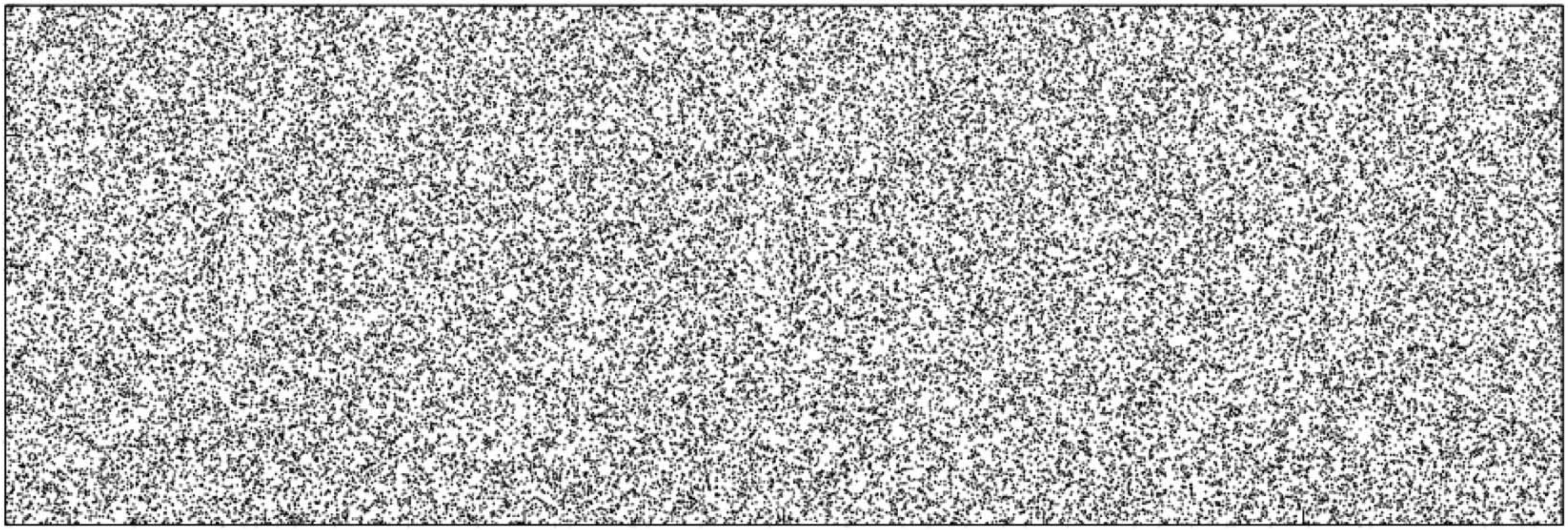
- Homogenization has two exponential rates: slower one related to lobes
- **Fast rate due to braiding of ghost rods!**

Topological entropy continuity across critical point



topological entropy as a function of τ_f

Identifying 'ghost rods'?



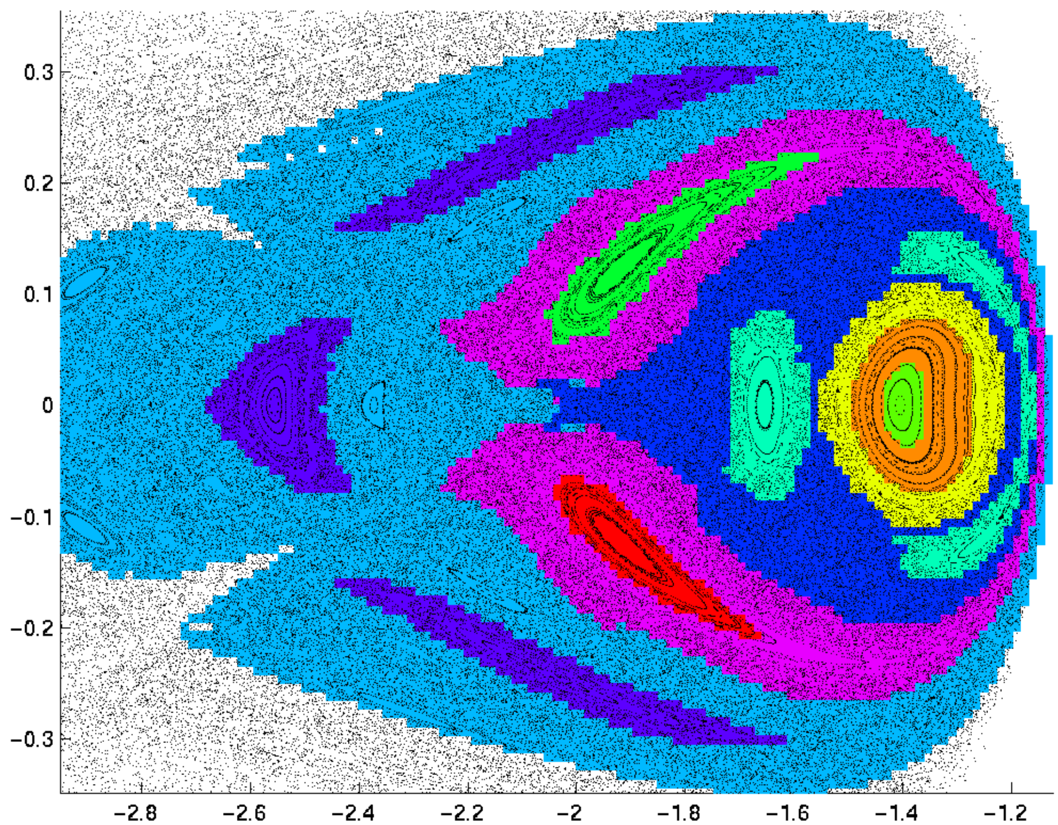
Poincaré section for $\tau_f < 1 \Rightarrow$ no obvious structure!

- Note the absence of any elliptical islands
- No periodic orbits of low period were found
- Is the phase space featureless?

Almost-invariant set (AIS) approach

- Partition phase space into **loosely coupled regions**

"Leaky" regions with a long residence time[3]



3-body problem phase space is divided into several invariant and almost-invariant sets.

---

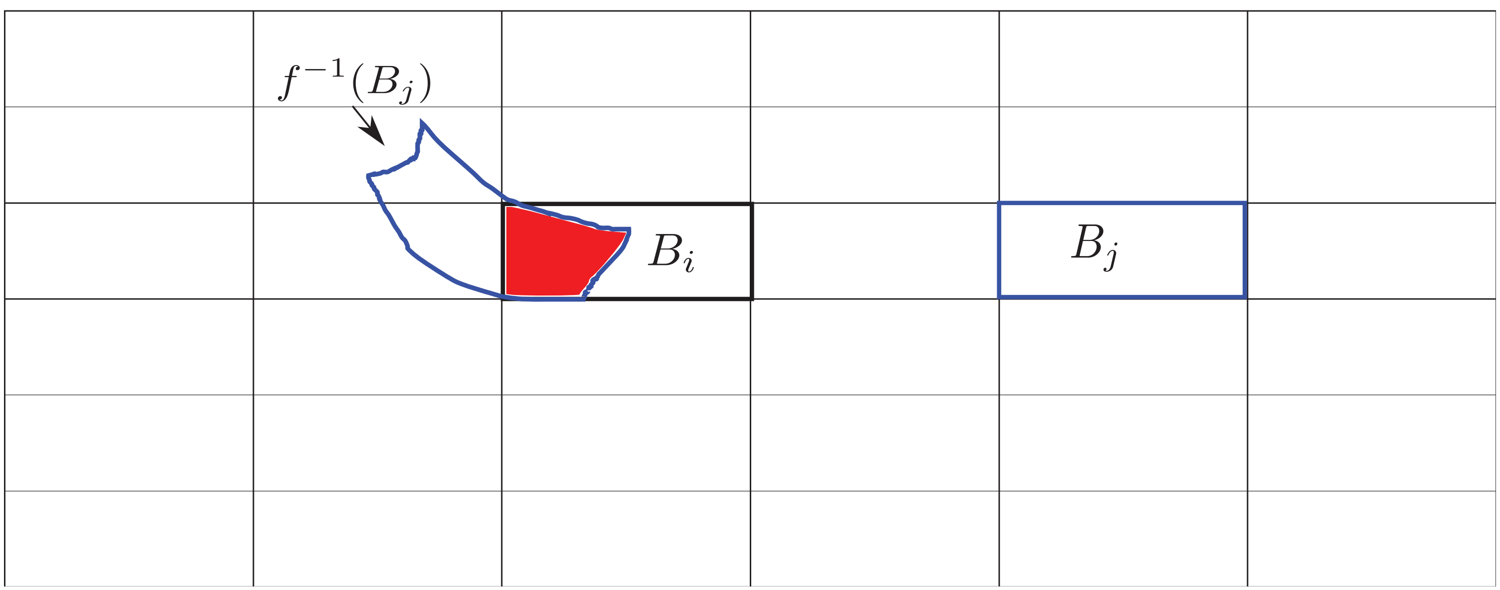
[3] work of Dellnitz, Junge, Froyland, et al

Almost-invariant set (AIS) approach

- Create box partition of phase space $\mathcal{B} = \{B_1, \ldots B_q\}$, with q large
- Consider a q-by-q **transition (Ulam) matrix**, P, for our dynamical system, where

$$P_{ij} = \frac{\mu(B_i \cap f^{-1}(B_j))}{\mu(B_i)},$$

the *transition probability* from B_i to B_j using, e.g., $f = \phi_t^{t+T}$

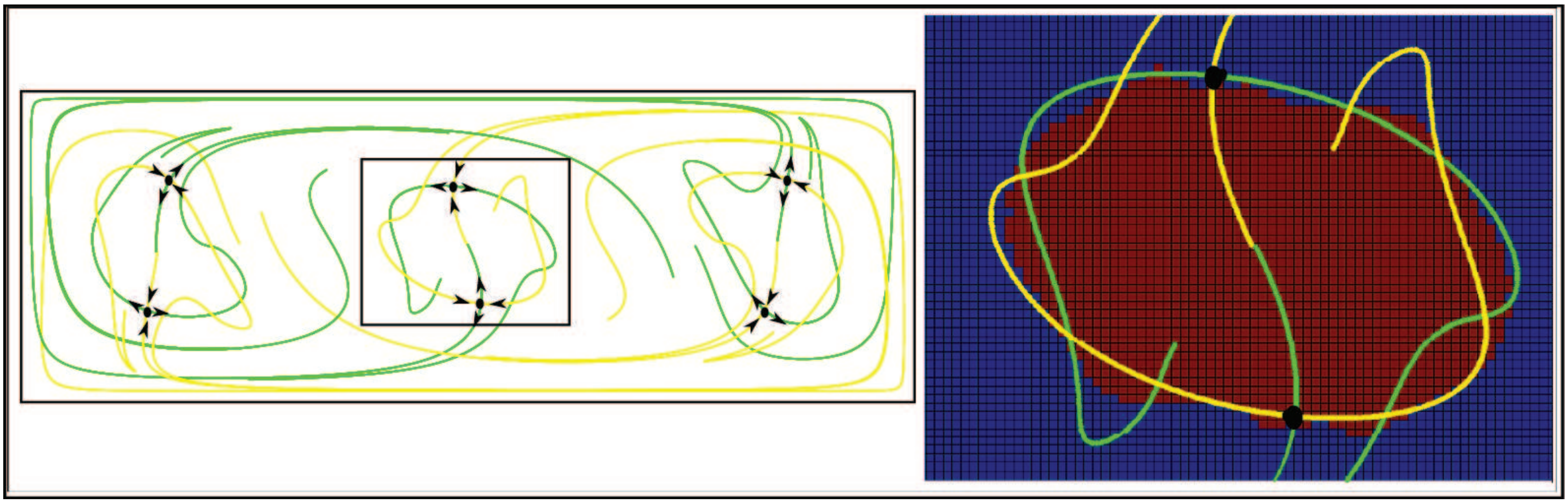


- P approximates our dynamical system via a finite state Markov chain.

Almost-invariant set (AIS) approach

- A set B is called almost invariant over the interval $[t, t+T]$ if
$$\rho_\mu(B) = \frac{\mu(B \cap \phi^{-1}(B))}{\mu(B)} \approx 1.$$
Can maximize value of ρ_μ over all possible combinations of sets $B \in \mathcal{B}$.
- In practice, AIS or relatedly, almost-cyclic sets (ACS), identified via **eigenvectors** of P or graph-partitioning
- Appropriate for non-autonomous, aperiodic, finite-time settings

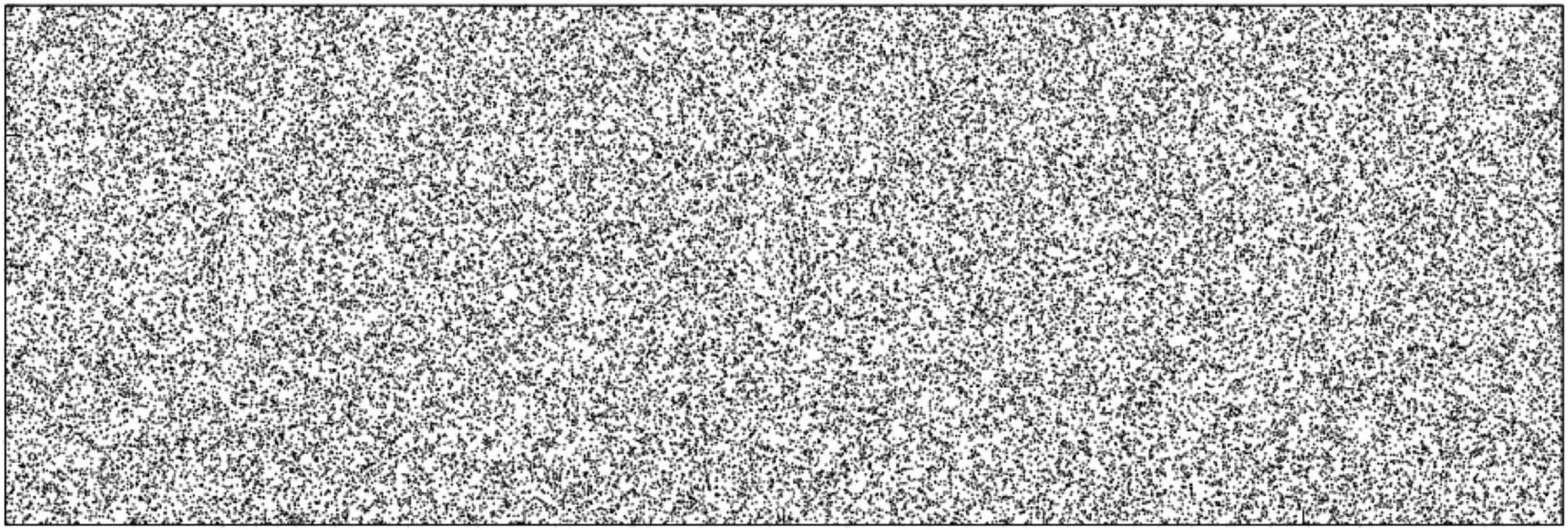
Identifying 'ghost rods': almost-cyclic sets



- Return to $\tau_f > 1$ case, where periodic points and manifolds exist
- Good agreement between AIS boundaries and manifolds of fixed points
- Known previously[4] and applies to more general objects than fixed points, i.e. normally hyperbolic invariant manifolds (NHIMs)

[4] Dellnitz, Junge, Lo, Marsden, Padberg, Preis, Ross, Thiere [2005] Phys. Rev. Lett.; Dellnitz, Junge, Koon, Lekien, Lo, Marsden, Padberg, Preis, Ross, Thiere [2005] Int. J. Bif. Chaos

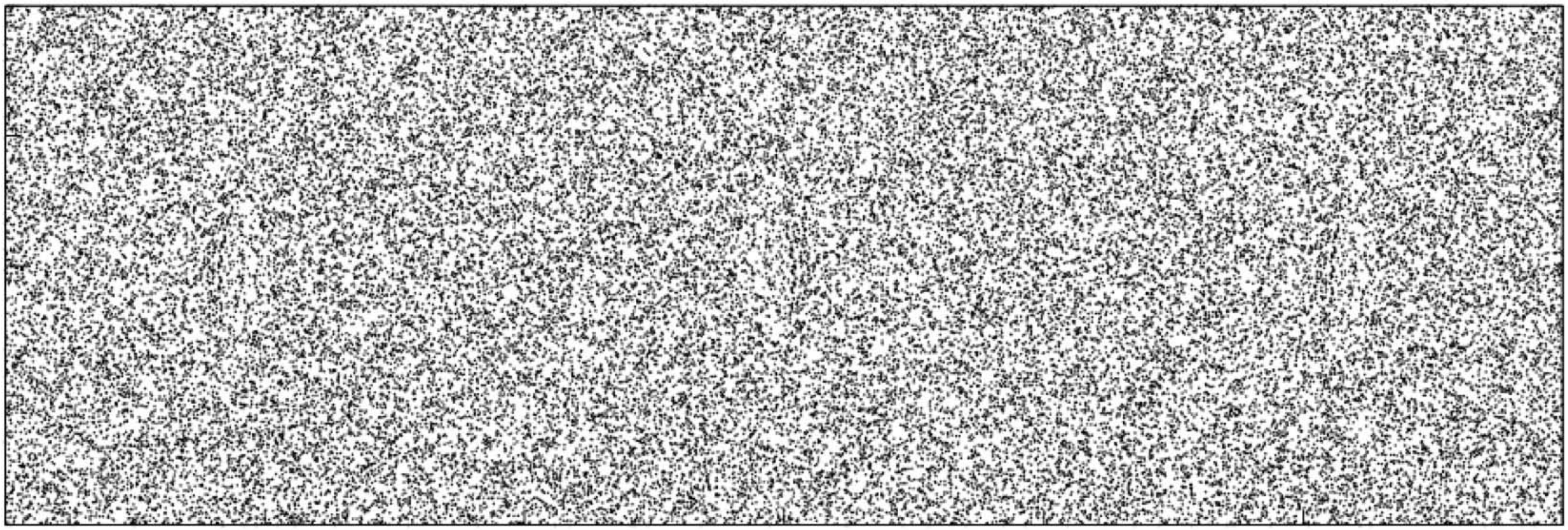
Identifying ‘ghost rods’: almost-cyclic sets



Poincaré section for $\tau_f < 1 \Rightarrow$ no obvious structure!

- Return to $\tau_f < 1$ case, where no periodic orbits of low period known
- Is the phase space featureless?

Identifying ‘ghost rods’: almost-cyclic sets

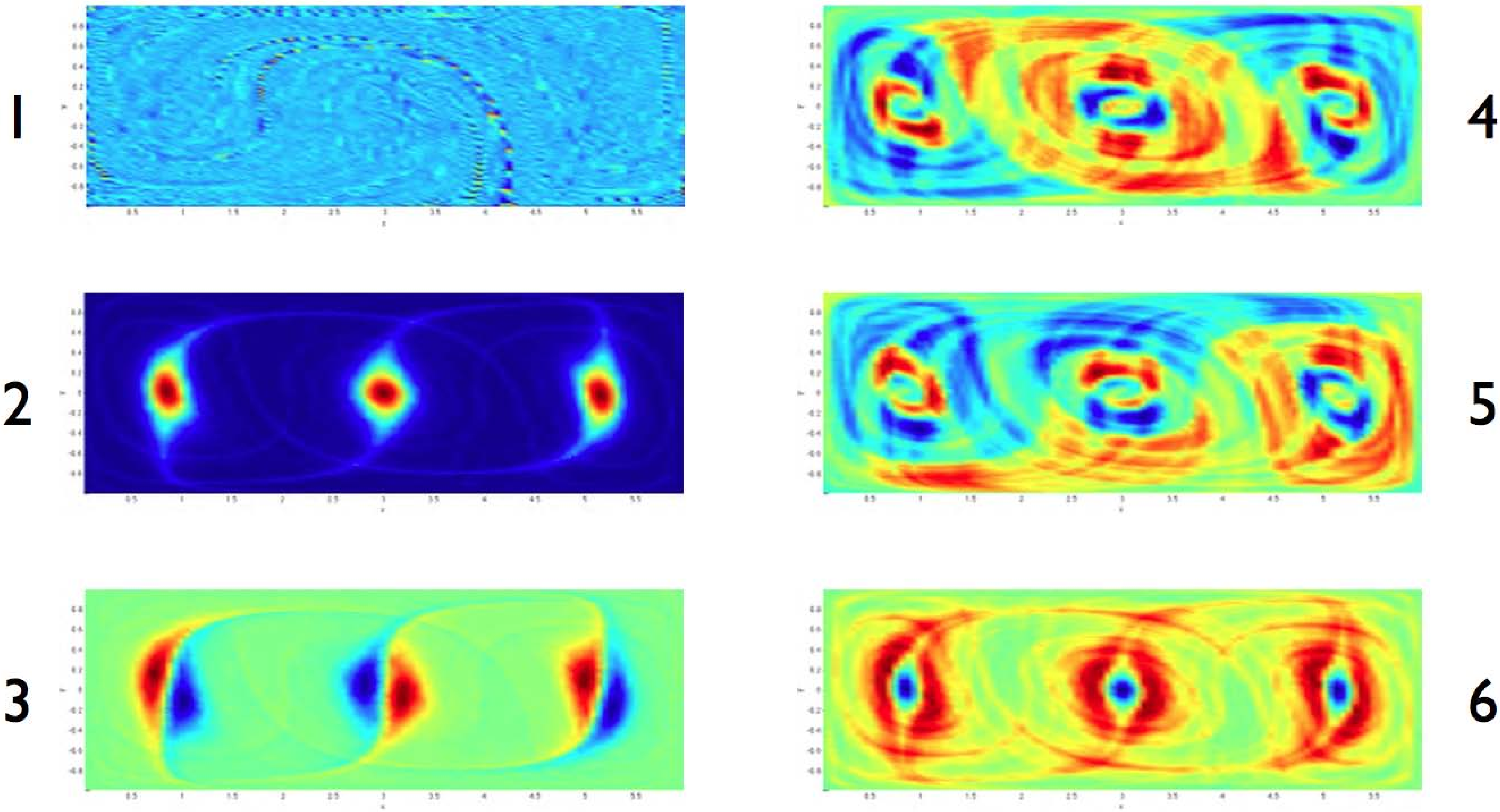


Poincaré section for $\tau_f < 1 \Rightarrow$ no obvious structure!

- Return to $\tau_f < 1$ case, where no periodic orbits of low period known
- Is the phase space featureless?

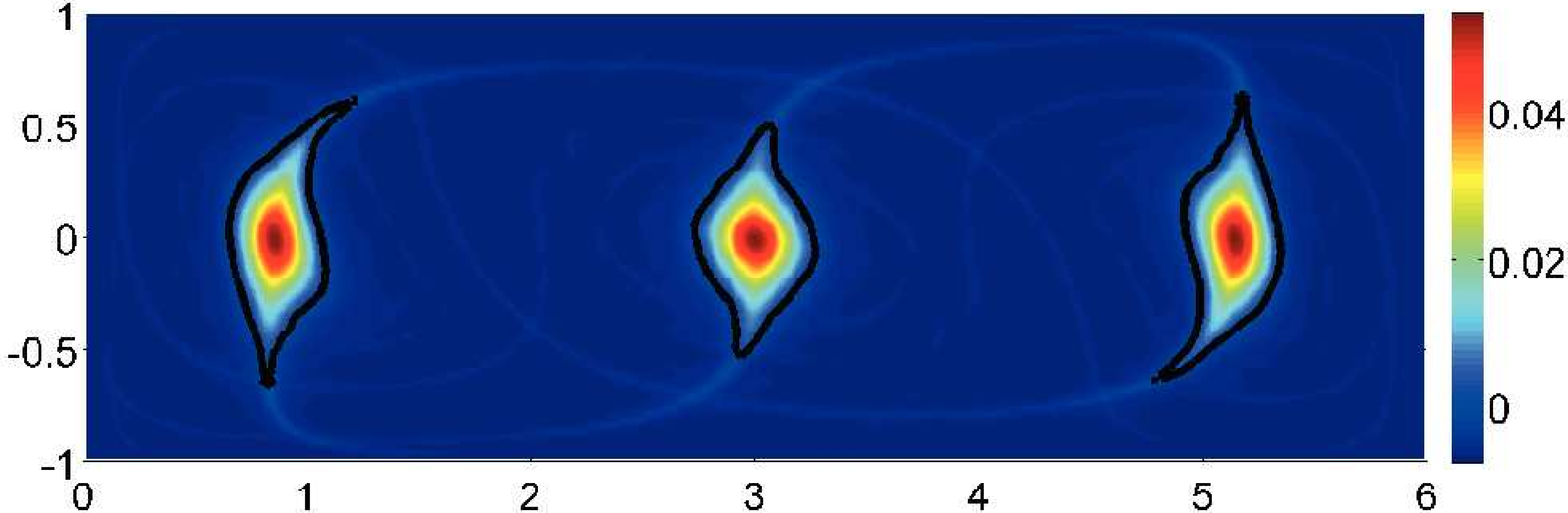
- Consider transition matrix $P_t^{t+\tau_f}$ induced by Poincaré map $\phi_t^{t+\tau_f}$

Identifying 'ghost rods': almost-cyclic sets



Top six eigenvalues for $\tau_f = 0.99$

Identifying ‘ghost rods’: almost-cyclic sets



The zero contour (black) is the boundary between the two almost-invariant sets.

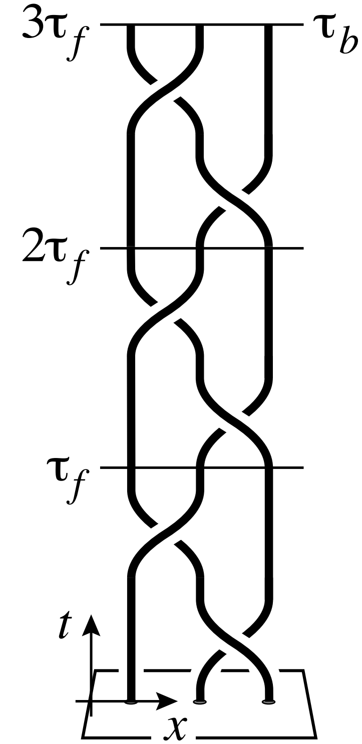
- The disconnected AIS is made of three almost-cyclic sets, with period 3

Identifying 'ghost rods': almost-cyclic sets

Almost-cyclic sets stirring the surrounding fluid like 'ghost rods'
— **works even when periodic orbits are absent!**

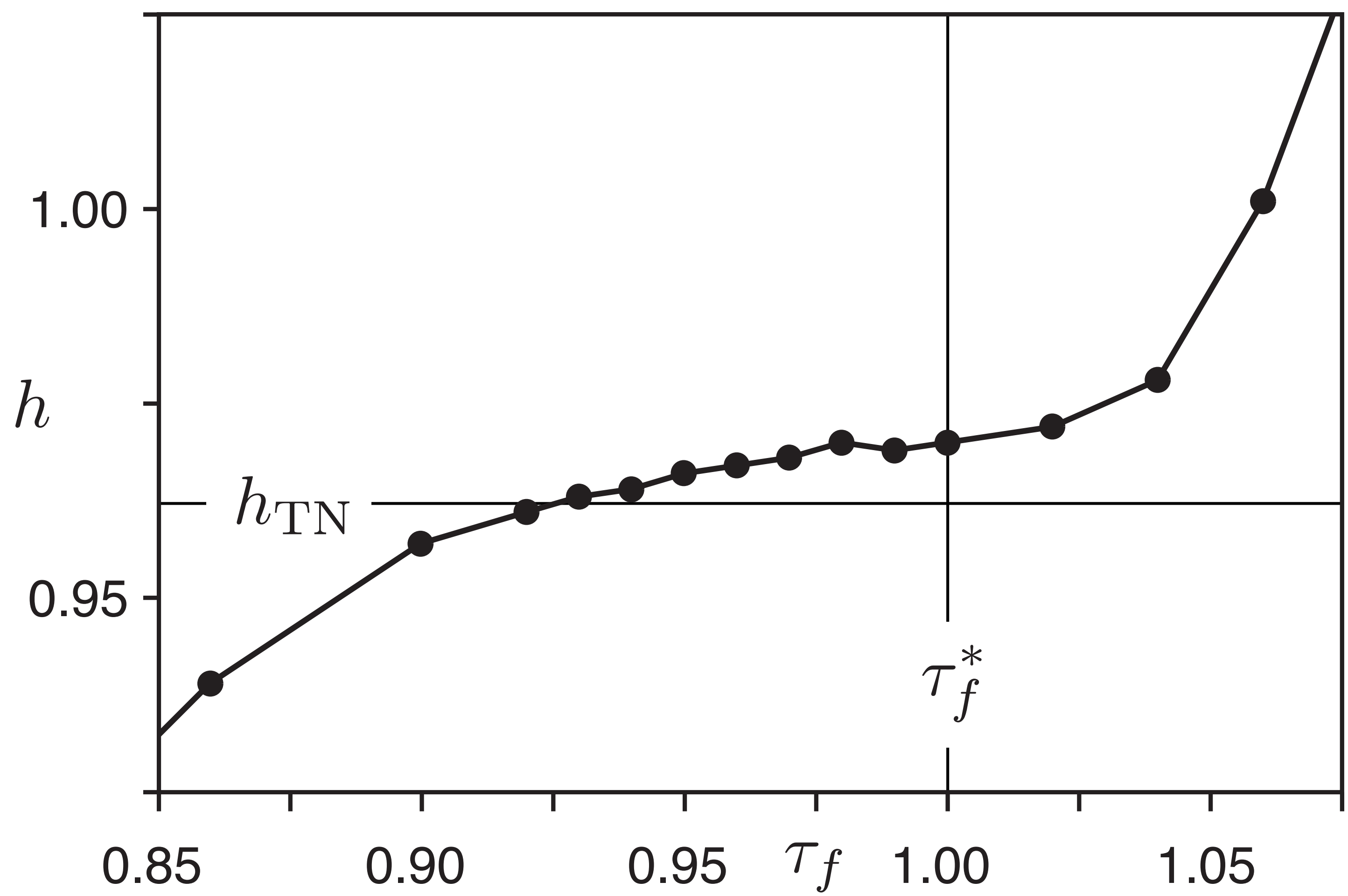
Movie shown is second eigenvector for $P_t^{t+\tau_f}$ for $t \in [0, \tau_f)$

Identifying 'ghost rods': almost-cyclic sets



Braid of ACSs gives lower bound of entropy via Thurston-Nielsen
— One only needs approximately cyclic blobs of fluid
— Even though the theorems require exactly periodic points!
— Stremler, Ross, Grover, Kumar [2011] Phys. Rev. Lett.

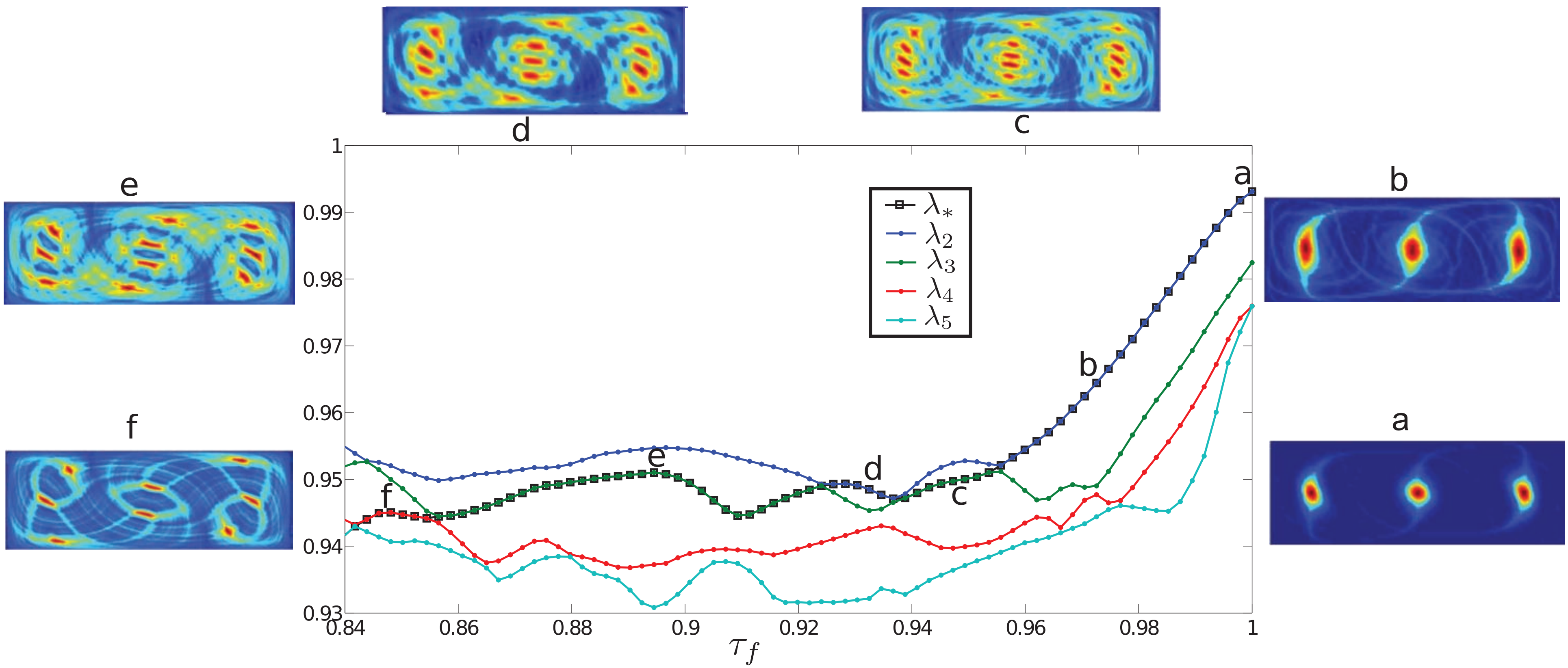
Topological entropy vs. bifurcation parameter



topological entropy as a function of τ_f

- h_{TN} shown for ACS braid on 3 strands

Eigenvalues/eigenvectors vs. bifurcation parameter



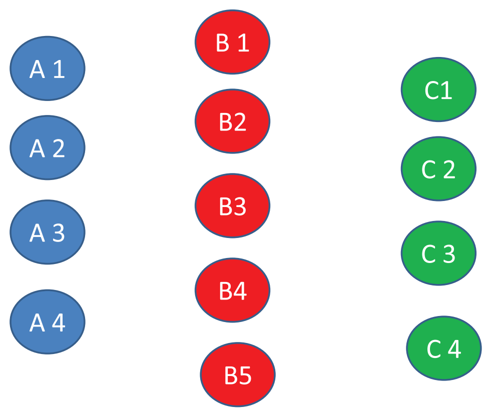
Bifurcation of ACSs

For example, braid on 13 strands for $\tau_f = 0.92$

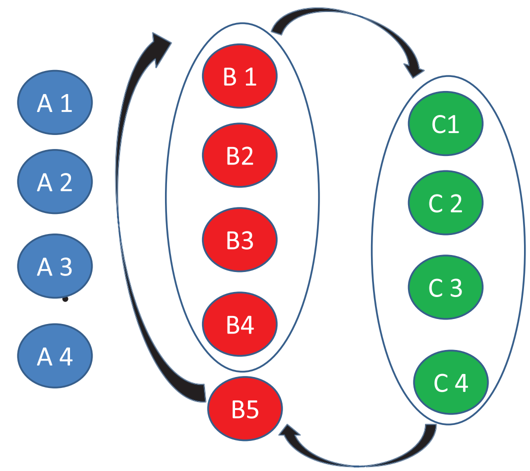
Movie shown is second eigenvector for $P_t^{t+\tau_f}$ for $t \in [0, \tau_f)$

Thurson-Nielsen for this braid provides lower bound on topological entropy

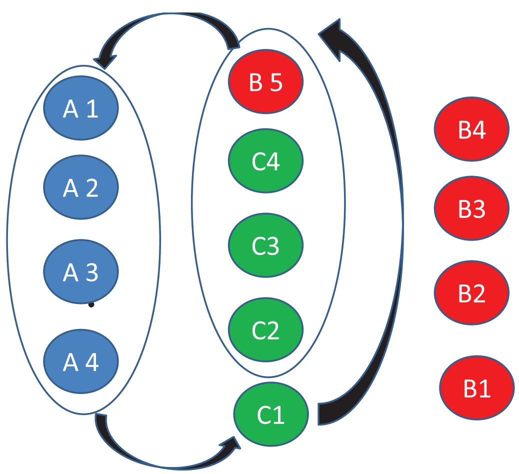
Bifurcation of ACSs



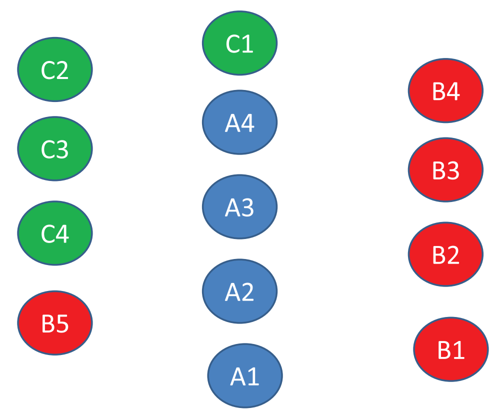
(a) Initial state



(b) First half-period

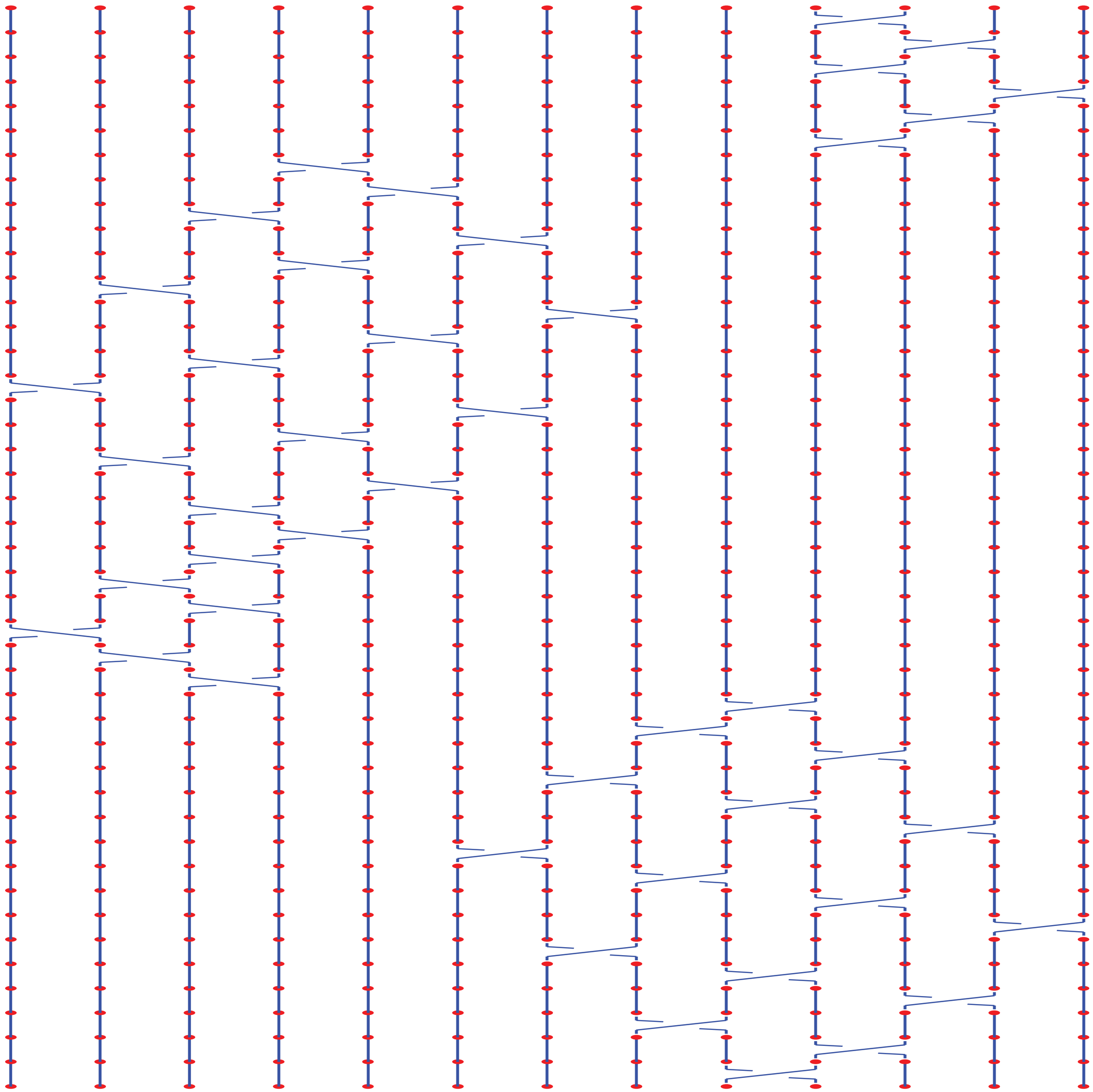


(c) Second half-period



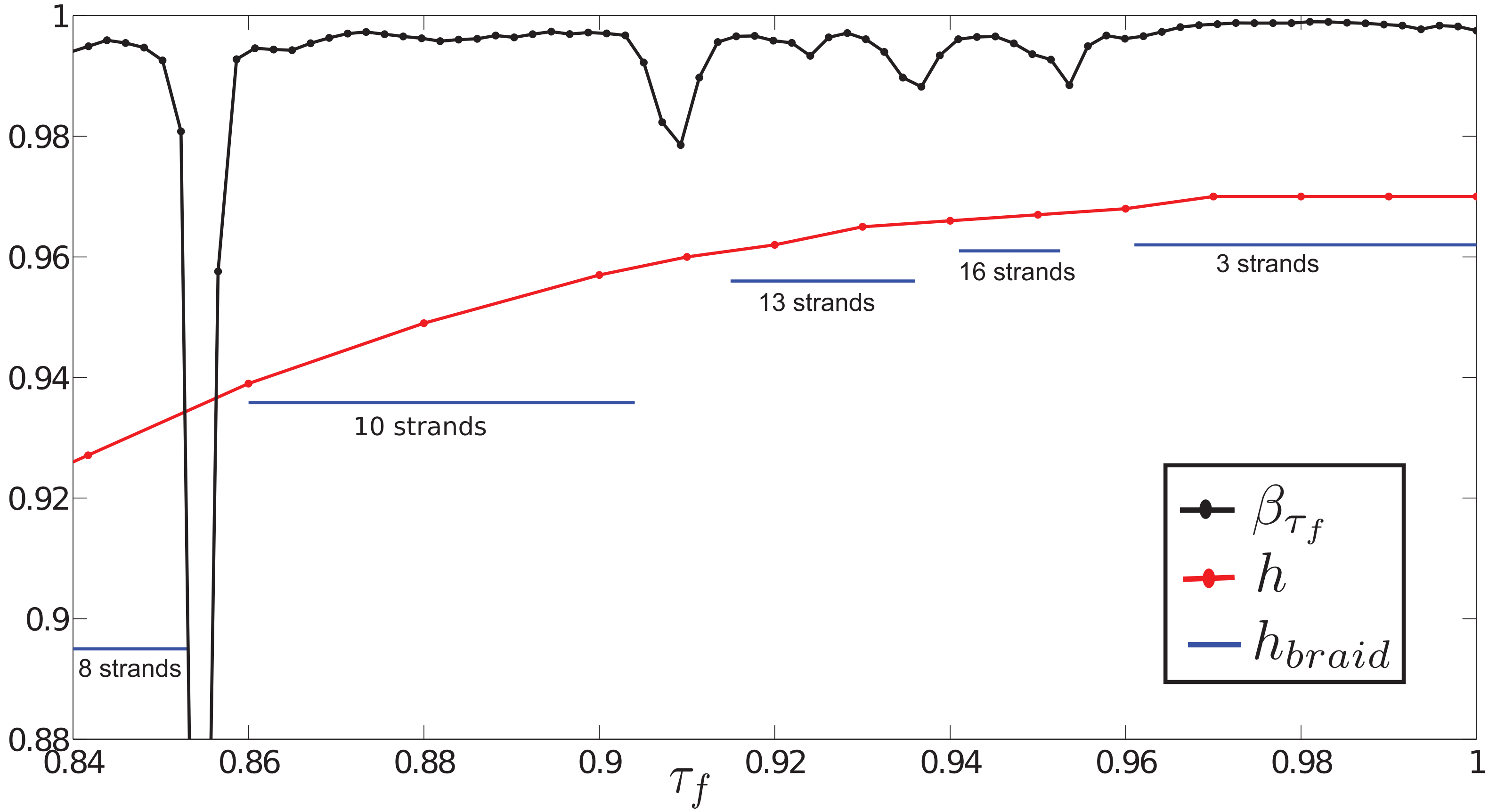
(d) State after 1 period

Bifurcation of ACSs



representation of braid

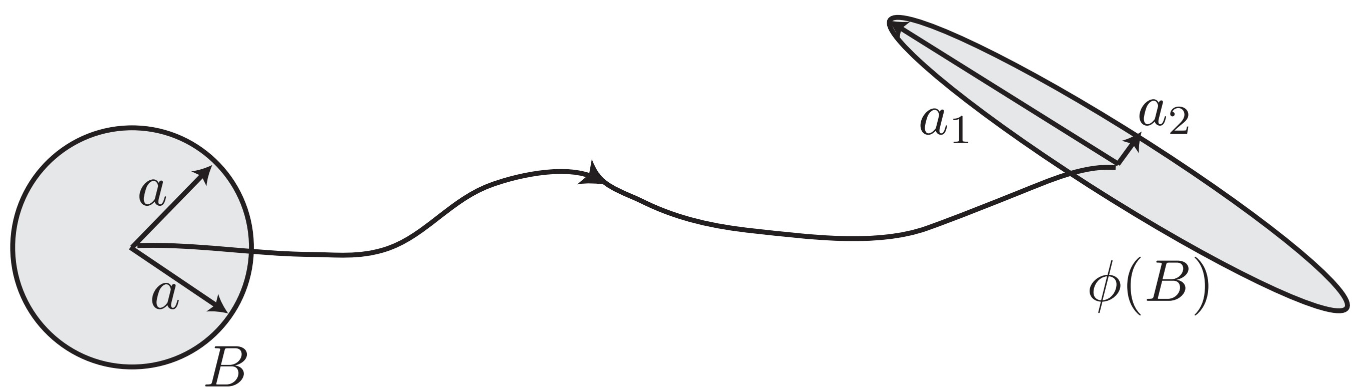
Sequence of ACS braids bounds entropy



For various braids of ACSs, the calculated entropy is given, bounding from below the true topological entropy over the range where the braid exists

Coherent sets and set-based definition of FTLE

- Consider, e.g., a flow ϕ_t^{t+T} in $(x_1, x_2) \in \mathbb{R}^2$.
- Treat the evolution of set $B \subset \mathbb{R}^2$ as evolution of two random variables X_1 and X_2 defined by probability density function $f(x_1, x_2)$, initially uniform on B, $f = \frac{1}{\mu(B)}\mathcal{X}_B$, with $\mathcal{X}_B$ the characteristic function of B.
- Under the action of the flow ϕ_t^{t+T}, f is mapped to Pf where P is the associated Perron-Frobenius operator.
- Let $I(f)$ be the covariance of f and $I(Pf)$ the covariance of Pf.

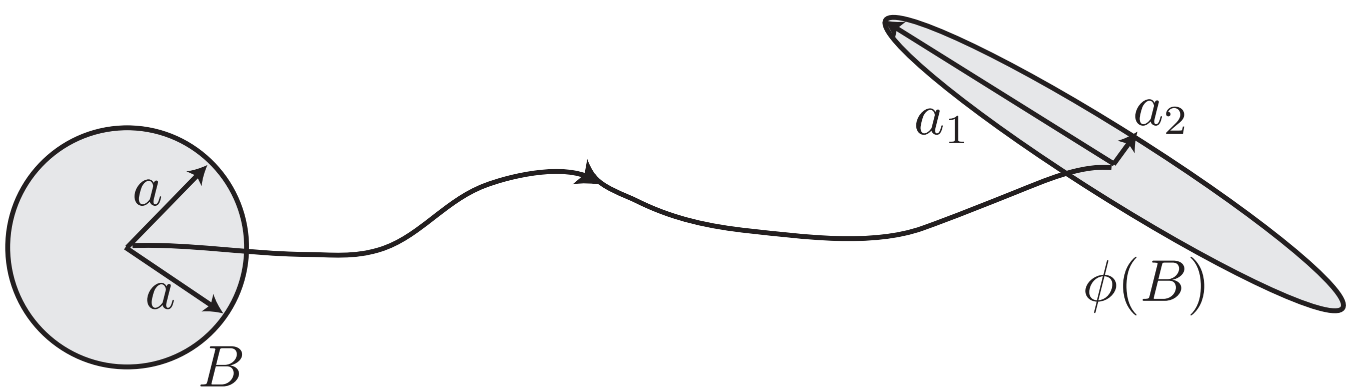


Deformation of a disk under the flow during $[t, t+T]$

Coherent sets and set-based definition of FTLE

- **Definition.** The **covariance-based FTLE** of B is
$$\sigma_I(B,t,T) = \frac{1}{|T|}\log\left(\frac{\sqrt{\lambda_{max}(I(Pf))}}{\sqrt{\lambda_{max}(I(f))}}\right).$$

- Reduces to usual definition of FTLE in the limit that the linearization approximation (i.e., line-stretching method) is valid



Deformation of a disk under the flow during $[t, t+T]$

Coherent sets and set-based definition of FTLE

- The **coherence** of a set B during $[t, t+T]$ is $\sigma_I(B, t, T)$.
- A set B is **almost-coherent** during $[t, t+T]$ if $\sigma_I(B, t, T) \approx 0$.

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- Captures the essential feature of a coherent set: it does not mix or spread significantly in the domain.
- This definition also can identify non-mixing **translating** sets.

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- **Values of $\sigma_I(B, t, T)$ determine the family of sets of various degrees of coherence**.
- Need to set a heuristic threshold on the value of $\sigma_I(B, t, T)$ to determine coherent sets.

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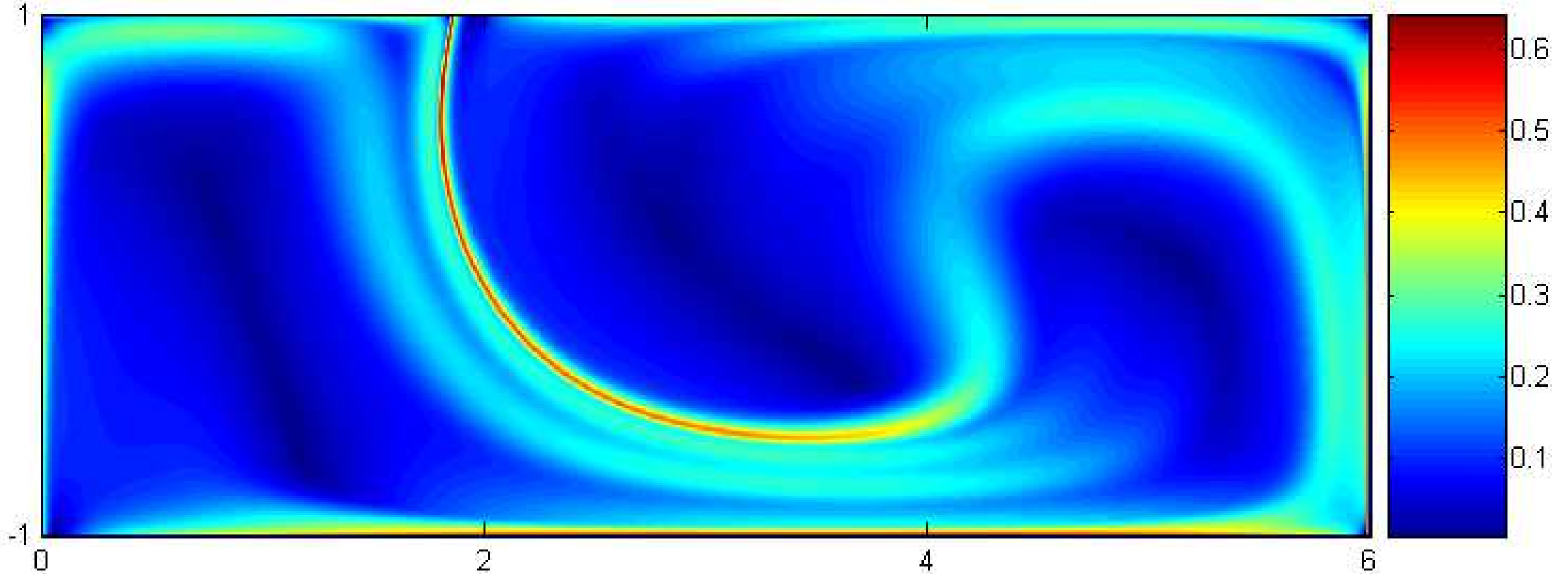
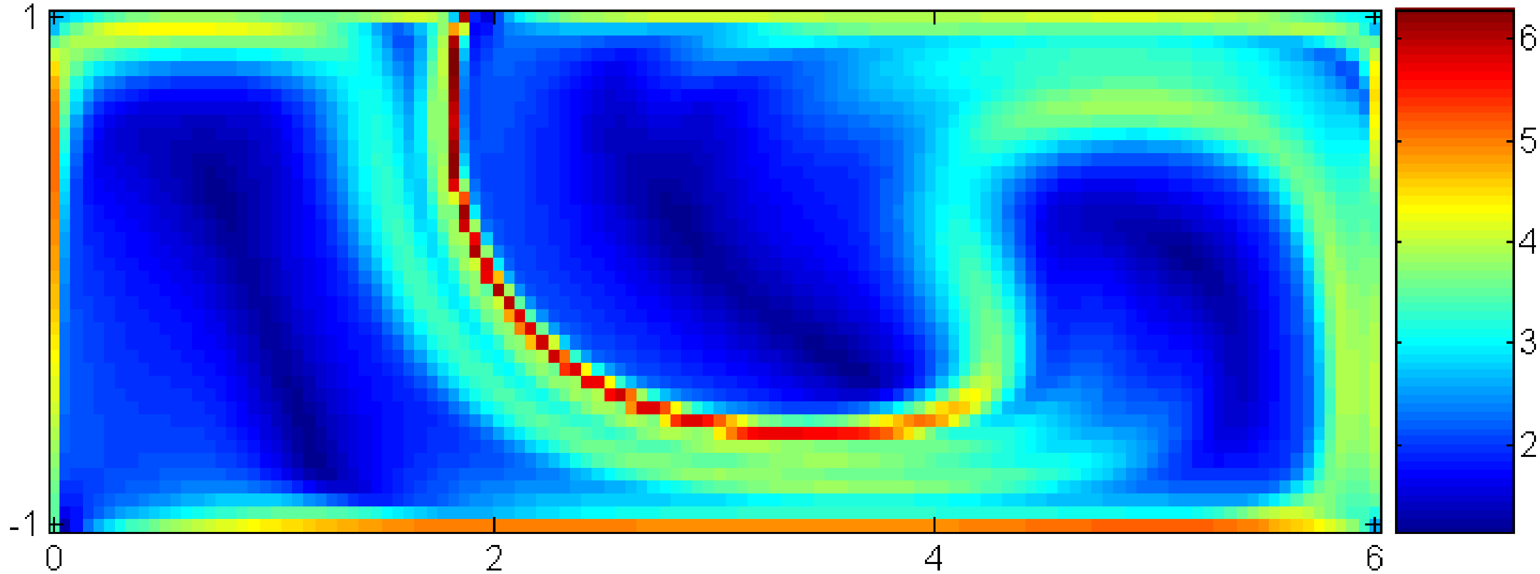
- Notice, coherent sets will be separated by ridges of high FTLE, i.e., LCS

Coherent sets in lid-driven cavity flow



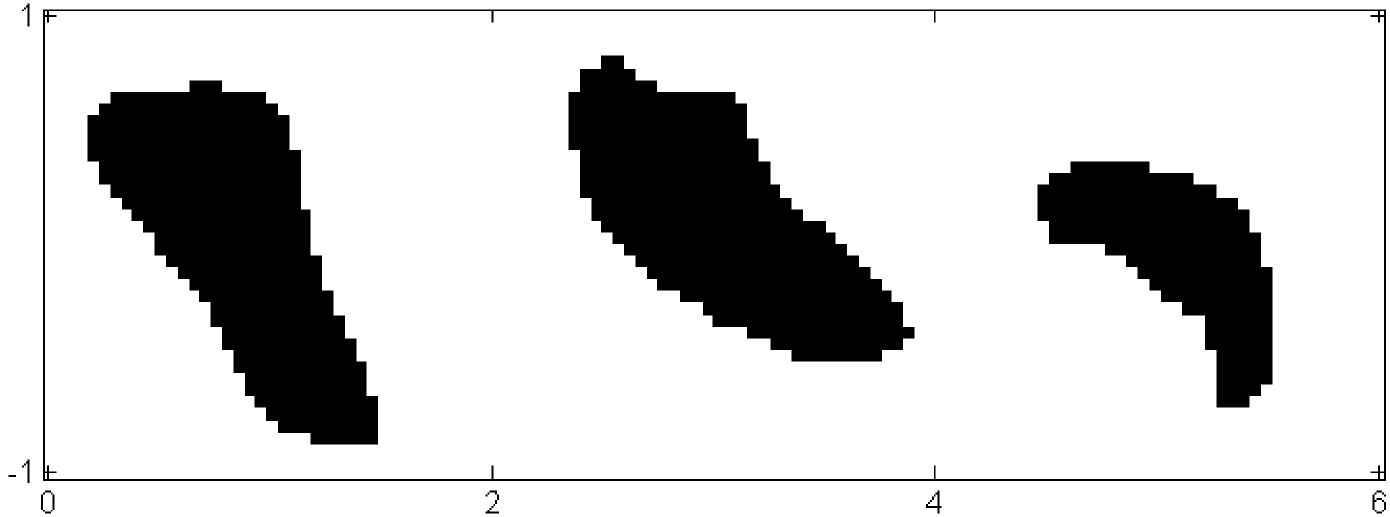
FTLE from line-stretching (conventional) during $[0, \tau_f]$

Coherent sets in lid-driven cavity flow



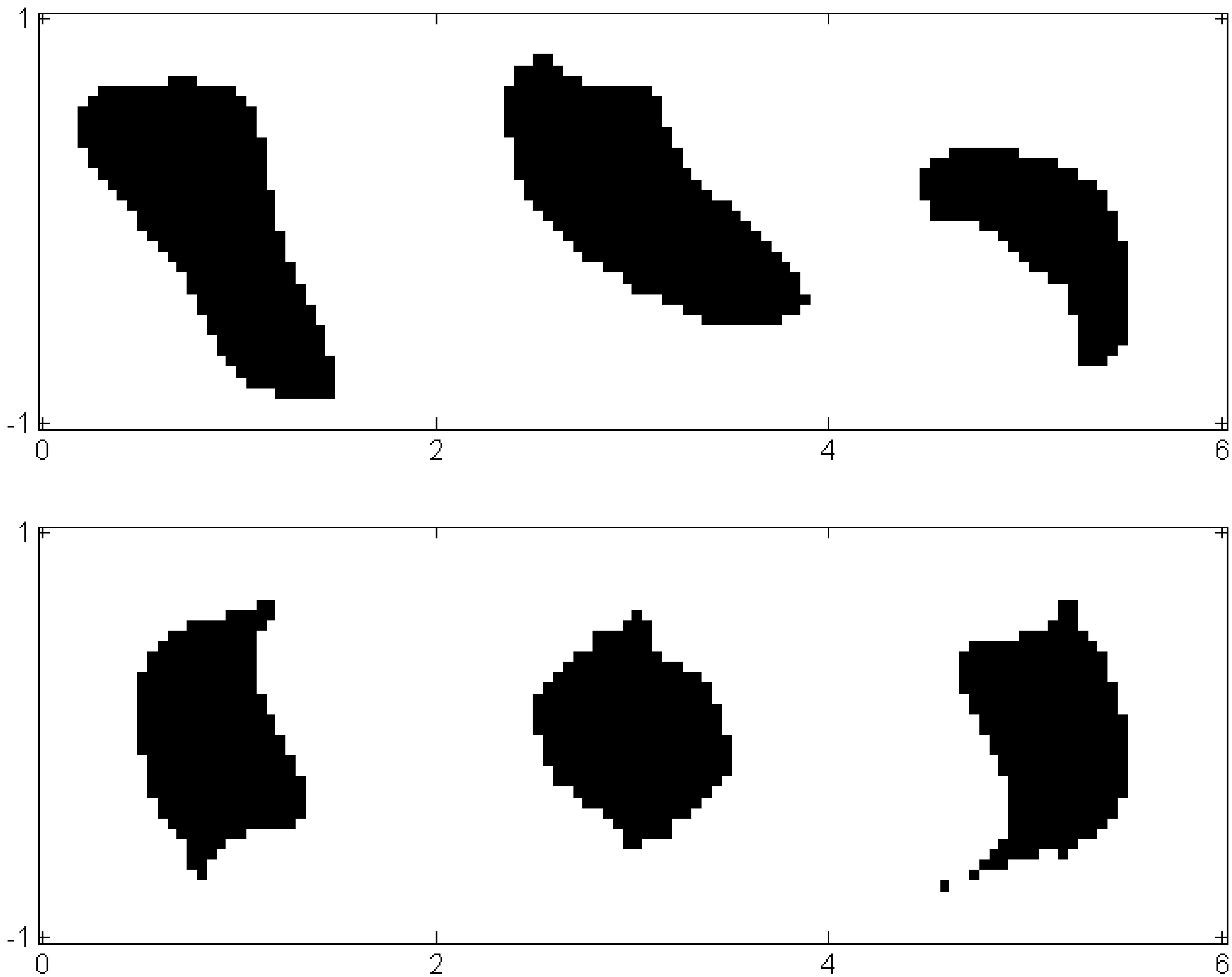
FTLE from covariance-based approach during $[0, \tau_f]$

Coherent sets in lid-driven cavity flow



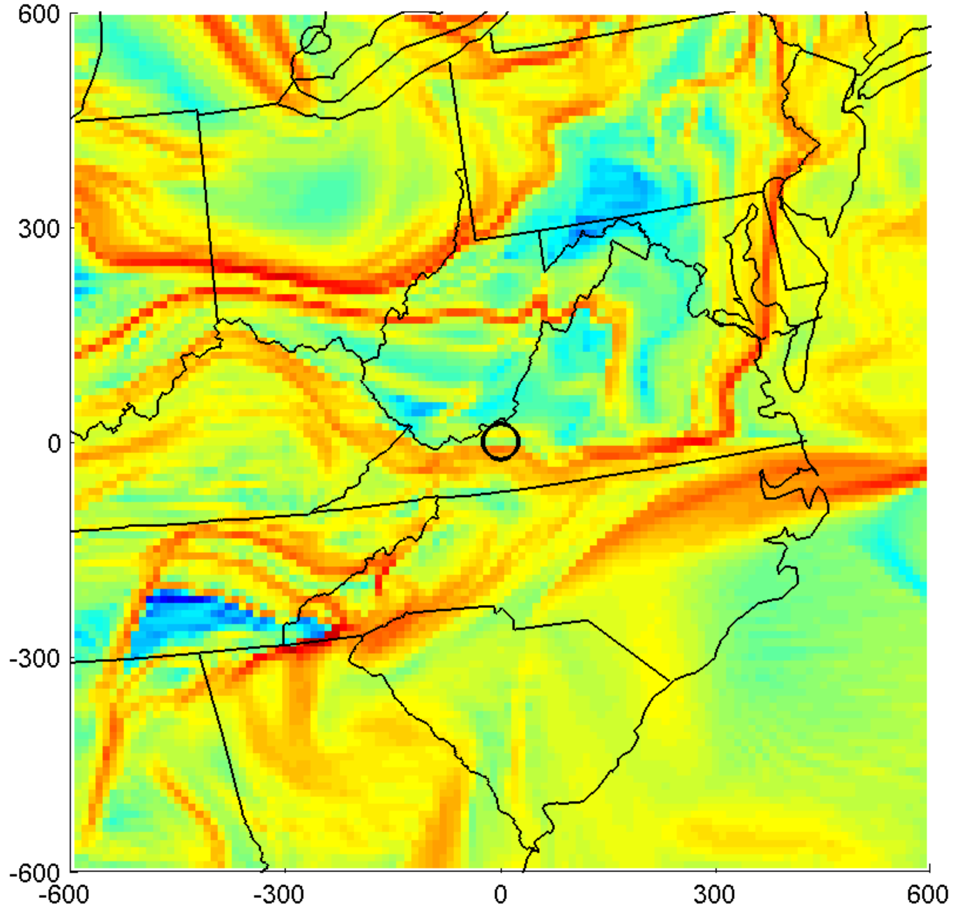
Sets of coherences $\sigma_I(0, \tau_f) < 1.6$

Coherent sets in lid-driven cavity flow



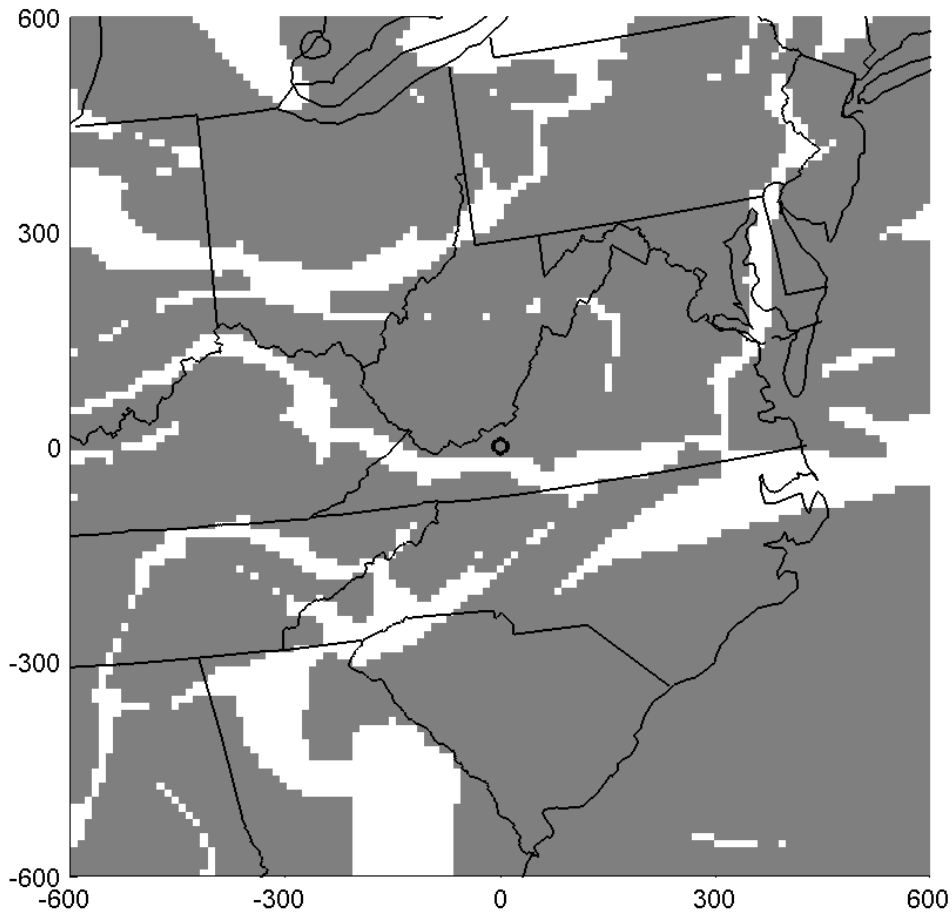
Compare with AIS from second eigenvector of P

Coherent sets in the atmosphere



- FTLE from covariance during 24 hours starting 09:00 1 May 2007

Coherent sets in the atmosphere



- Coherent sets during 24 hours starting 09:00 1 May 2007

Optimal navigation in an aperiodic setting?

- Selectively 'jumping' between coherent air masses using control
- Moving between mobile subregions of different finite-time itineraries

Final words on chaotic transport

- What are the robust descriptions of transport which work in data-driven aperiodic, finite-time settings?

  - Possibilities: finite-time lobe dynamics, finite-time symbolic dynamics may work

  - For these, use set-oriented approach

  - Many links between invariant manifolds, FTLE, LCS, AIS/coherent sets, and topological methods
    — e.g., boundaries between coherent sets are naturally LCS; follows from covariance-based definition of FTLE
    — fixed points $\Rightarrow$ AIS, so stable/unstable invariant manifolds $\Rightarrow$???

The End

For papers, movies, etc., visit:
www.shaneross.com

Main Papers:

- Stremler, Ross, Grover, Kumar [2011] Topological chaos and periodic braiding of almost-cyclic sets. *Physical Review Letters* 106, 114101.
- Senatore & Ross [2011] Detection and characterization of transport barriers in complex flows via ridge extraction of the finite time Lyapunov exponent field, *International Journal for Numerical Methods in Engineering* 86, 1163.
- Lekien & Ross [2010] The computation of finite-time Lyapunov exponents on unstructured meshes and for non-Euclidean manifolds. *Chaos* 20, 017505.
- Tallapragada & Ross [2011] A geometric and probabilistic description of coherent sets. Preprint.
- Grover, Ross, Stremler, Kumar [2011] Topological chaos, braiding and breakup of almost-invariant sets. Preprint.
- Tallapragada & Ross [2008] Particle segregation by Stokes number for small neutrally buoyant spheres in a fluid, *Physical Review E* 78, 036308.